# Electric Dipole Moments – probes of fundamental symmetries

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#### Outline

### **Electric Dipole Moments (EDMs)**

- What is it?
- Why is it interesting?
- What do we know about it?
- How to measure (charged particle) EDMs?



## What is it?



#### **Electric Dipoles**





	atomic physics	hadron physics
charges	е	
$ \vec{r}_1 - \vec{r}_2 $	1 Å= 10 <sup>-8</sup> cm	
EDM		
naive expectation	10 <sup>−8</sup> <i>e</i> · cm	
observed	water molecule	
	2 · 10 <sup>−8</sup> <i>e</i> · cm	



	atomic physics	hadron physics
charges	е	е
$ \vec{r}_1 - \vec{r}_2 $	1 Å= 10 <sup>-8</sup> cm	$1 \mathrm{fm} = 10^{-13} \mathrm{cm}$
EDM		
naive expectation	10 <sup>−8</sup> <i>e</i> · cm	$10^{-13} e \cdot cm$
observed	water molecule	neutron
	2 · 10 <sup>−8</sup> <i>e</i> · cm	$< 3 \cdot 10^{-26} e$ · cm





neutron EDM of  $d_n = 3 \cdot 10^{-26} e \cdot cm$  corresponds to separation of u- from d-quarks of  $\approx 5 \cdot 10^{-26} cm$ 

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### Operator $\vec{d} = q\vec{r}$

is odd under parity transformation  $(\vec{r} \rightarrow -\vec{r})$ :

 $\mathcal{P}^{-1}\vec{d}\mathcal{P}=-\vec{d}$ 

Consequences: In a state  $|a\rangle$  of given parity the expectation value is 0:

$$\langle a | \vec{d} | a \rangle = - \langle a | \vec{d} | a \rangle$$
  
If  $| a \rangle = \alpha | P = + \rangle + \beta | P = - \rangle$   
in general  $\langle a | \vec{d} | a \rangle \neq 0 \Rightarrow$  i.e. molecules



#### EDM of molecules



ground state: mixture of  $\Psi_s = \frac{1}{\sqrt{2}} (\Psi_1 + \Psi_2)$  P = + $\Psi_a = \frac{1}{\sqrt{2}} (\Psi_1 - \Psi_2)$  P = -(Cohen-Tannoudji, B. Diu, F. Laloë, Mécanique quantique)



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**Molecules** can have large EDM because of degenerated ground states with different parity



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**Elementary particles** (including hadrons) have a definite parity and cannot posess an EDM  $P|had >= \pm 1|had >$ 



**Molecules** can have large EDM because of degenerated ground states with different parity

**Elementary particles** (including hadrons) have a definite parity and cannot posess an EDM  $P|had >= \pm 1|had >$ 

unless

 ${\cal P}$  and time reversal  ${\cal T}$  invariance are violated!



#### ${\mathcal T}$ and ${\mathcal P}$ violation of EDM



 $\Rightarrow \text{EDM measurement tests violation of fundamental symmetries } \mathcal{P} \text{ and } \mathcal{T}(\stackrel{\mathcal{CPT}}{=} \mathcal{CP})$ 

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#### Symmetries in Standard Model

	electro-mag.	weak	strong
${\mathcal C}$	$\checkmark$	£	$\checkmark$
${\cal P}$	$\checkmark$	ź	$\checkmark$
$\mathcal{T} \stackrel{\textit{CPT}}{\rightarrow} \mathcal{CP}$	$\checkmark$	(ź)	(√)

- *C* and *P* are maximally violated in weak interactions (Lee, Yang, Wu)
- *CP* violation discoverd in kaon decays (Cronin,Fitch) described by CKM-matrix in Standard Model
- CP violation allowed in strong interaction but corresponding parameter  $\theta_{QCD} \lesssim 10^{-10}$  (strong CP-problem)

### Symmetries

- EDM requires violation of symmetries
- but particles may have large magnetic dipole moment (MDM),
- for **structureless** particles theory even predicts that  $e^{\hbar} |\vec{S}|$

$$\mu = g \frac{e n}{2m} \frac{|\mathbf{S}|}{\hbar}$$
 with  $g = 2$  in leading order

$$G = \frac{g-2}{2}$$
 for various particles:

	experiment	theory
electron	$1159652180.73(0.28)\cdot 10^{-12}$	$1159652181.13(0.86)\cdot 10^{-12}$
muon	$1165920.80(54)(33)\cdot 10^{-9}$	$1165918.28(49)\cdot 10^{-9}$
proton	1.792847356(23)	2*

\* ): static quark model, SU(6) wave function

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#### Nucleon Spin Puzzle





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## Why is it interesting?



#### $\mathcal{CP}$ violation

- We are surounded by matter (and not anti–matter)  $\eta = \frac{n_B - n_{\bar{B}}}{n_{\gamma}} = \mathbf{6} \times \mathbf{10}^{-10}$
- In 1967 Sakharov formulated three prerequisites for baryogenesis. One of these is the combined violation of the charge and parity, CP, symmetry.
- Starting from equal amount of matter and anti-matter at the Big Bang, from  $\mathcal{CP}$ -violation in Standard Model we expect only  $10^{-18}$
- New CP violating sources outside the realm of the SM are clearly needed to explain this discrepancy of eight orders of magnitude.
- They could manifest in EDMs of elementary particles







#### History of Neutron EDM



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no EDM observed yet, only limits





- no EDM observed yet, only limits
- no measurement for deuteron (or heavier nuclei),

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- Beyond SM values accessible by experiments





charged particle EDM measurements less precise





- charged particle EDM measurements less precise
- To measure EDMs one needs large electric fields. Charged particles are accelerated by electric fields





GOAL of JEDI (Jülich Electric Dipole Investigations)collaboration: Charged Hadron EDM measurements

- First measurement of deuteron, <sup>3</sup>He EDM,
- first direct measurement of proton EDM

ultimately with a precision of  $10^{-29}e$  cm

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# How to measure charged particle EDMs?



## Measurement of charged particle EDMs Generic Idea:

For **all** edm experiments (neutron, proton, atom, ...): Interaction of  $\vec{d}$  with electric field  $\vec{E}$ For charged particles: apply electric field in a storage ring:



Wait for build-up of vertical polarization  $s_\perp \propto |d|$ , then determine  $s_\perp$  using polarimeter

In general:

$$\frac{\mathrm{d}\vec{\boldsymbol{s}}}{\mathrm{d}t} = \vec{\Omega} \times \vec{\boldsymbol{s}}$$



Spin Motion is governed by Thomas-BMT equation (Bargmann, Michel, Telegdi)

$$rac{\mathrm{d}ec{s}}{\mathrm{d}t} = ec{\Omega} imes ec{s}$$
  
 $ec{\Omega} = rac{e\hbar}{mc} [Gec{B} + \left(G - rac{1}{\gamma^2 - 1}
ight) ec{v} imes ec{E} + rac{1}{2} \eta (ec{E} + ec{v} imes ec{B})]$ 

$$ec{d}=\etarac{e\hbar}{2mc}ec{S},\quadec{\mu}=2(G+1)rac{e\hbar}{2m}ec{S},\quad G=rac{g-2}{2},$$

- $\vec{d}$ : electric dipole moment  $\vec{\mu}$ : magnetic moment, g:g-factor, G: anomalous magnetic moment
- $\gamma$ : Lorentz factor

V. Bargmann, L. Michel and V. L. Telegdi, Phys. Rev. Lett. 2 (1959) 435.



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Several Options (try to get rid terms  $\propto$  G):



$$\vec{\Omega} = \frac{e\hbar}{mc} [G\vec{B} + \left(G - \frac{1}{\gamma^2 - 1}\right) \vec{v} \times \vec{E} + \frac{1}{2} \eta (\vec{E} + \vec{v} \times \vec{B})]$$

Several Options (try to get rid terms  $\propto$  G):

• Pure electric ring  
with 
$$\left(G - \frac{1}{\gamma^2 - 1}\right) = 0$$
, works only for  $G > 0$ 



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Several Options (try to get rid terms  $\propto$  G):

Pure magnetic ring

#### Required field strength

	$G=rac{g-2}{2}$	<i>p</i> /GeV/c	<i>E<sub>R</sub>/MV/m</i>	$B_V/T$
proton	1.79	0.701	10	0
deuteron	-0.14	1.0	-4	0.16
<sup>3</sup> He	-4.18	1.285	17	-0.05

Ring radius  $\approx$  40m Smaller ring size possible if  $B_V \neq 0$  for proton  $E = \frac{GBc\beta\gamma^2}{1 + G\beta^2\gamma^2}$ 





Figure 3: An all-electric storage ring lattice for measuring the electric dipole moment of the proton. Except for having longer straight sections and separated beam channels, the all-in-one lattice of Fig. 1 spatterned after this lattice. Quadrupole and sextupole families, and tunes and lattice functions of the allin-one lattice of Fig. 1 will be quite close to those given for this lattice in reference[3]. The match will be even closer with magnetic field set to zero for proton operation.

#### Brookhaven National Laboratory (BNL) Proposal



## 2. Combined $\vec{E}/\vec{B}$ ring



Figure 1: "All-In-One" lattice for measuring EDM's of protons, deuterons, and helions.

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Under discussion at Forschungszentrum Jülich (design: R. Talman)

Main advantage:

Experiment can be performed at the existing (upgraded) COSY (COoler SYnchrotron) in Jülich on a shorter time scale!



COSY provides (polarized ) protons and deuterons with  $p = 0.3 - 3.7 \text{GeV}/c \Rightarrow$  Ideal starting point



$$ec{\Omega} = rac{e\hbar}{mc} \left( G ec{B} + rac{1}{2} \eta ec{v} imes ec{B} 
ight)$$

Problem:

Due to precession caused by magnetic moment, 50% of time longitudinal polarization component is || to momentum, 50% of the time it is anti-||.





$$\vec{\Omega} = \frac{e\hbar}{mc} \left( G\vec{B} + \frac{1}{2} \frac{\eta \vec{v} \times \vec{B}}{\eta \vec{v} \times \vec{B}} \right)$$

Problem:

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 $E^*$  field in the particle rest frame tilts spin due to EDM up and down  $\Rightarrow$  **no net EDM effect** 



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Use resonant "magic Wien-Filter" in ring  $(\vec{E} + \vec{v} \times \vec{B} = 0)$ :

 $E^* = 0 \rightarrow \text{part.}$  trajectory is not affected but

 $B^* \neq 0 \rightarrow$  mag. mom. is influenced

 $\Rightarrow$  net EDM effect can be observed!



Horizontal spin motion  $\propto G$ 

vertical spin motion  $s_\perp \propto d$ 



#### Summary of different options

	$\bigcirc$	
1.) pure electric ring (BNL)	no $\vec{B}$ field needed	works only for p
2.) combined ring (Jülich)	works for $p, d, {}^{3}\text{He}, \dots$	both <i>Ē</i> and <i>B</i> required
<ol> <li>pure magnetic ring (Jülich)</li> </ol>	existing (upgraded) COSY ring can be used , shorter time scale	lower sensitivity



#### Statistical Sensitivity

 $\sigma \approx \frac{\hbar}{\sqrt{NfT\tau_{p}}PEA}$ 

Ε	electric field	10 MV/m
Ρ	beam polarization	0.8
Α	analyzing power	0.6
Ν	nb. of stored particles/cycle	$4  imes 10^7$
f	detection efficiency	0.005
$ au_p$	spin coherence time	1000 s
Т	running time per year	10 <sup>7</sup> s

 $\Rightarrow \sigma \approx 10^{-29} e \cdot cm/year \text{ (for magnetic ring } \approx 10^{-24} e \cdot cm/year\text{)}$ Expected signal  $\approx$  3nrad/s (for  $d = 10^{-29} e \cdot cm$ ) (BNL proposal)

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#### **Electrostatic Deflectors**



- Electrostatic deflectors from Fermilab ( $\pm$ 125kV at 5 cm  $\hat{=}$  5MV/m)
- large-grain Nb at plate separation of a few cm yields  $\approx$  20MV/m



#### Wien filter



Conventional design R. Gebel, S. Mey (FZ Jülich)



stripline design D. Hölscher, J. Slim (IHF RWTH Aachen)



#### Polarimeter

Principle: Particles hit a target: Left/Right asymmetry gives information on EDM Up/Down asymmetry gives information on MDM





#### Polarimeter



Cross Section & Analyzing Power for deuterons



#### Spin Coherence Time (SCT)

Short Spin Coherence Time



#### Spin Coherence Time (SCT)

Large Spin Coherence Time



#### Results on Spin Coherence Time (SCT)



Spins decohere during storage time results form Cosy run May 2012 using correction sextupole

 $\Rightarrow$  SCT increase from a few s to  $\approx$  200s already reached

(Ed. Stephenson)



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#### **Systematics**

One major source: Radial *B* field mimics an EDM effect:

- Difficulty: even small radial magnetic field, *B<sub>r</sub>* can mimic EDM effect if :μ*B<sub>r</sub>* ≈ *dE<sub>r</sub>*
- Suppose  $d = 10^{-29} e \cdot cm$  in a field of E = 10 MV/m

• This corresponds to a magnetic field:  

$$B_r = \frac{dE_r}{\mu_N} = \frac{10^{-22} eV}{3.1 \cdot 10^{-8} eV/T} \approx 3 \cdot 10^{-17} T$$
(Earth Magnetic field  $\approx 5 \cdot 10^{-5} T$ )

Solution: Use two beams running clockwise and counter clockwise, separation of the two beams is sensitive to  $B_r$ 



#### **JEDI** Collaboration

- JEDI = Jülich Electric Dipole Moment Investigations
- $\approx$  80 members

(Aachen, Dubna, Ferrara, Ithaca, Jülich, Krakow, Michigan, St. Petersburg, Minsk, Novosibirsk, Stockholm, Tbilisi, ...)

•  $\approx$  10 PhD students





#### Storage Ring EDM Efforts





#### JARA FAME

#### JARA=Jülich Aachen Research Alliance New section founded: FAME (=Forces and Matter Experiments)





## Summary



#### Summary

- EDM of charged particles can be measured in storage rings
- EDMs of elementary particles are of high interest to disentangle various sources of CP violation searched for to explain matter - antimatter asymmetry in the Universe

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- Experimentally very challenging because effect is tiny
- Efforts at Brookhaven and Jülich to perform such measurements





#### Sources of $\mathcal{CP}$ violation



#### Sources of $\mathcal{CP}$ violation



 $\Rightarrow$  It is mandatory to measure EDM of many different particles to disentangle various sources of CP violation.