

Theory of Electric Dipole Moments of complex systems

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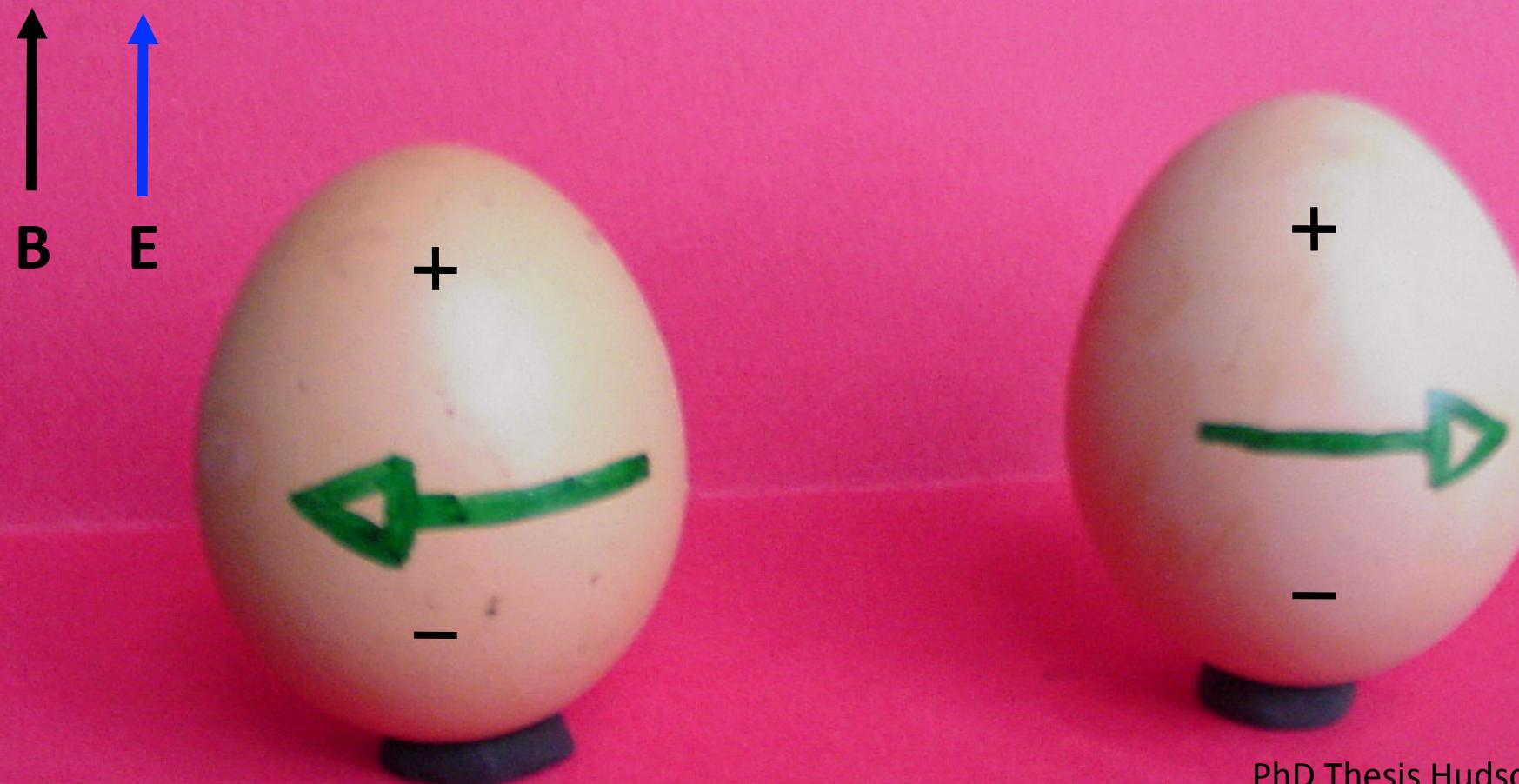
With: E. Mereghetti, W. Dekens, U. van Kolck, R. Timmermans
J. Bsaisou, A. Wirzba, Ulf-G. Meißner, C. Hanhart.....



Outline of this talk

- **Part I:** What are EDMs and why are they interesting in the first place ?
- **Part II:** Effective field theory framework
- **Part III:** EDMs of nucleons and nuclei

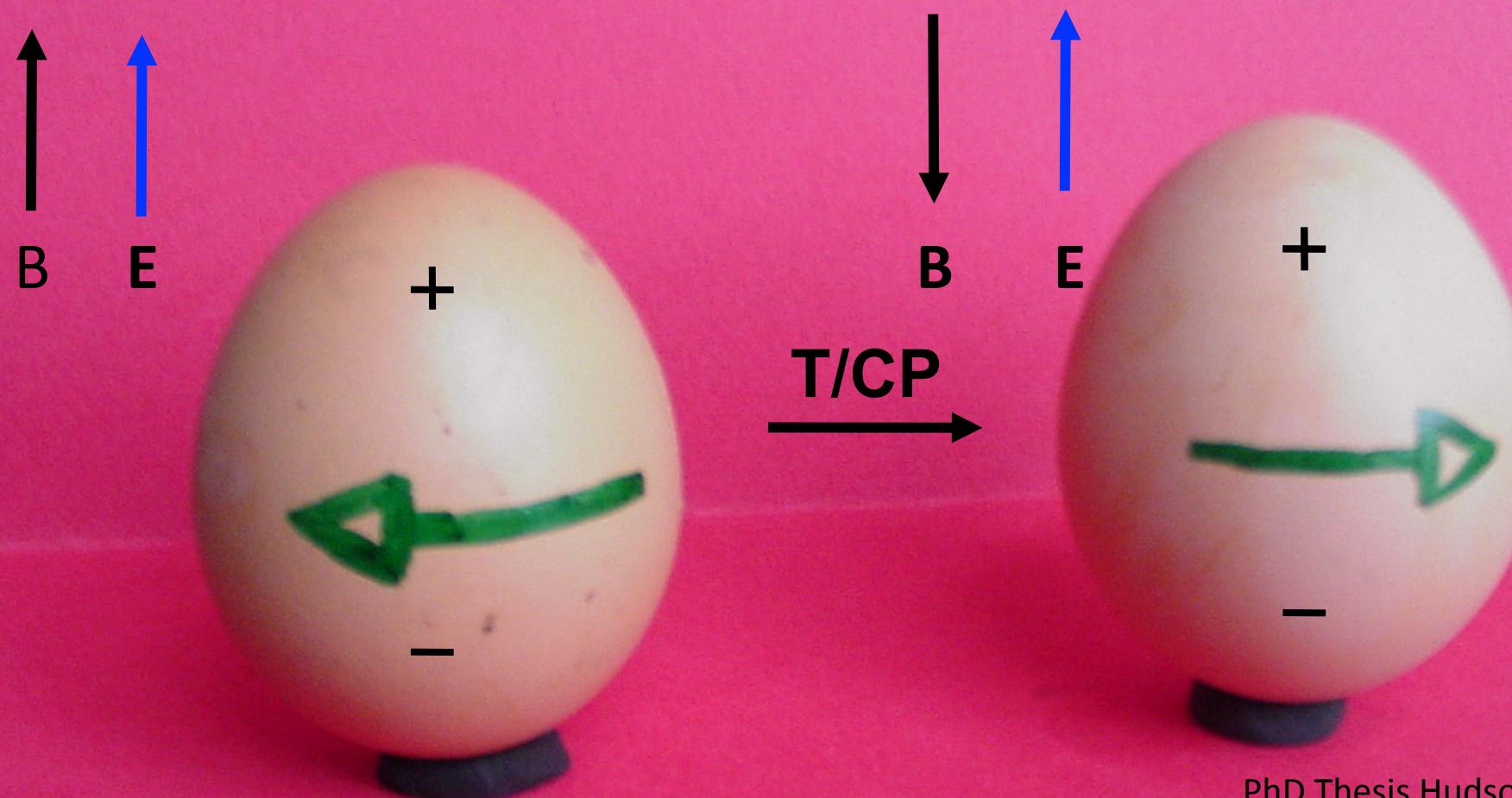
Electric Dipole Moments



PhD Thesis Hudson

$$H = -\mu(\vec{\sigma} \cdot \vec{B}) - d(\vec{\sigma} \cdot \vec{E})$$

Electric Dipole Moments



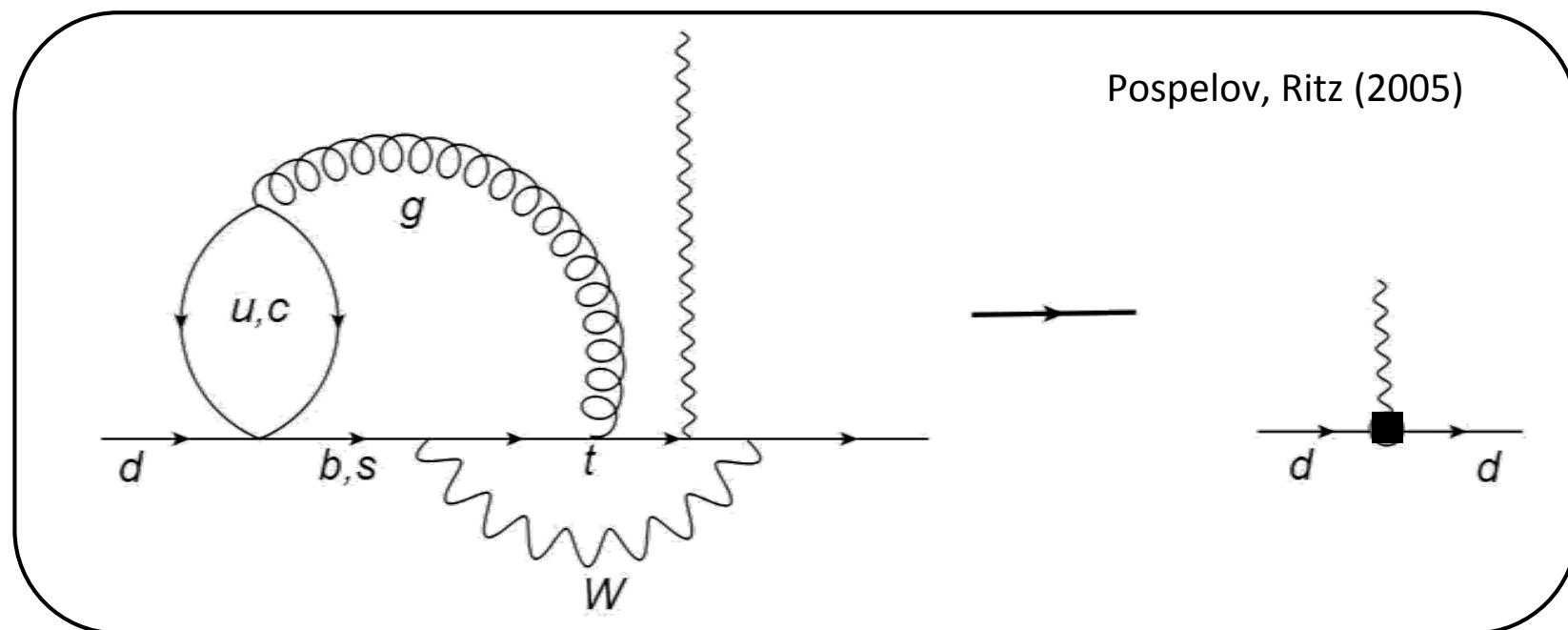
PhD Thesis Hudson

$$H = -\mu(\vec{\sigma} \cdot \vec{B}) + d(\vec{\sigma} \cdot \vec{E})$$

EDM's in the Standard Model

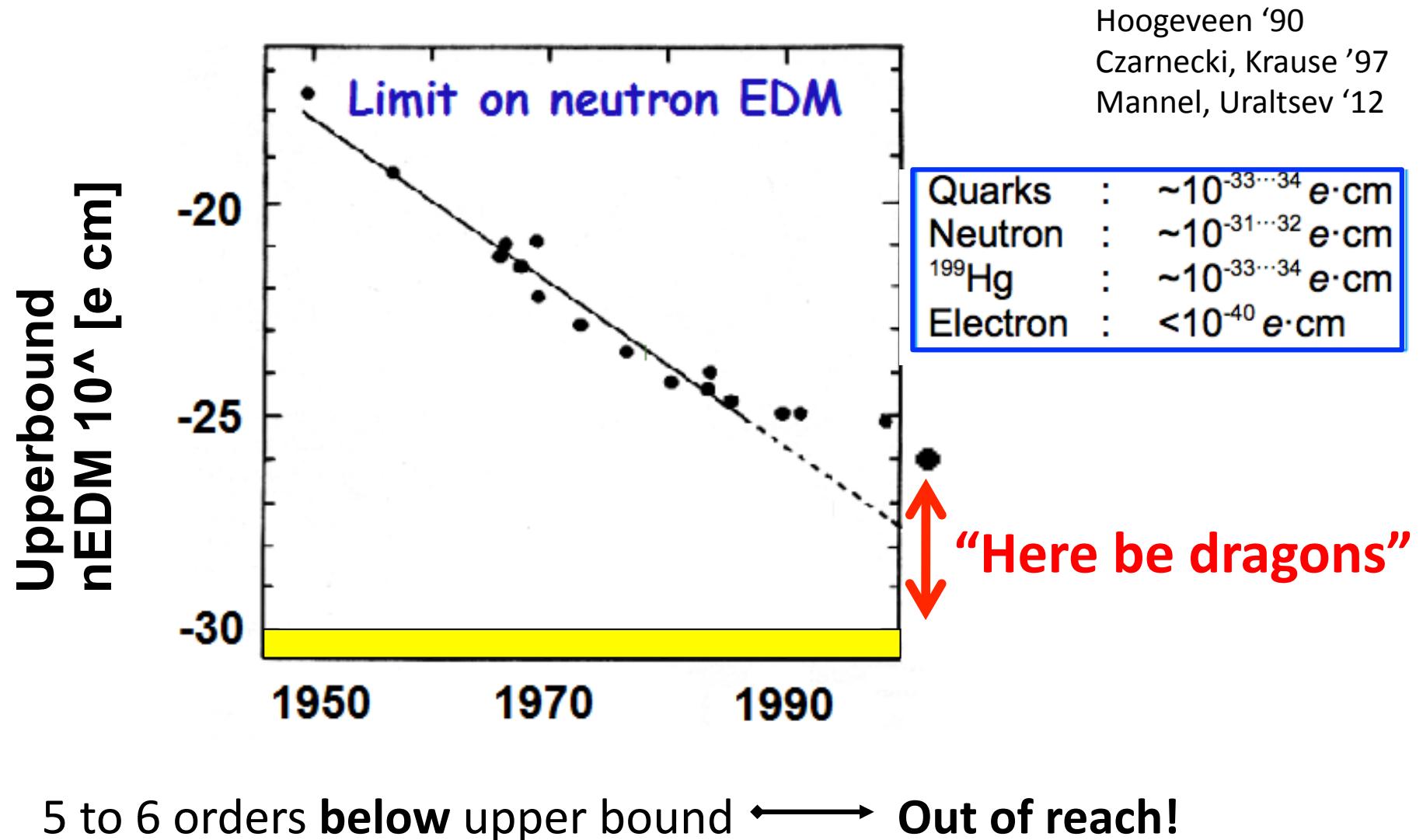
- Electroweak CP-violation
- Nobel prize for predicting **third** generation

$$L = \bar{d} \overline{\Psi} \sigma^{\mu\nu} i\gamma^5 \Psi F_{\mu\nu}$$



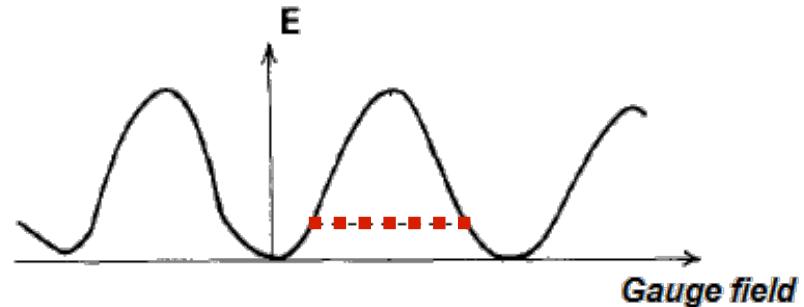
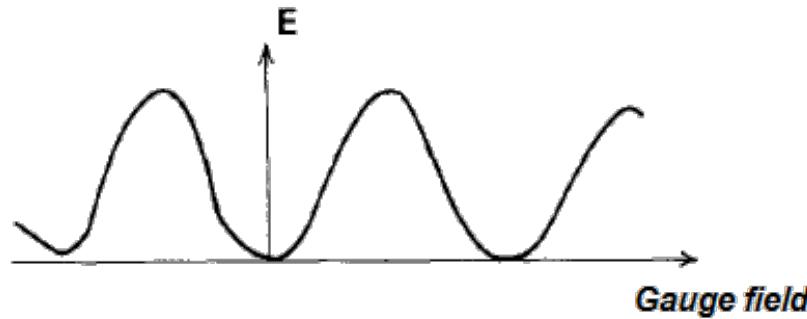
Highly Suppressed

Electroweak CP-violation



EDM's in the Standard Model

- Second source: QCD **theta-term**
- Due to complicated vacuum structure of QCD

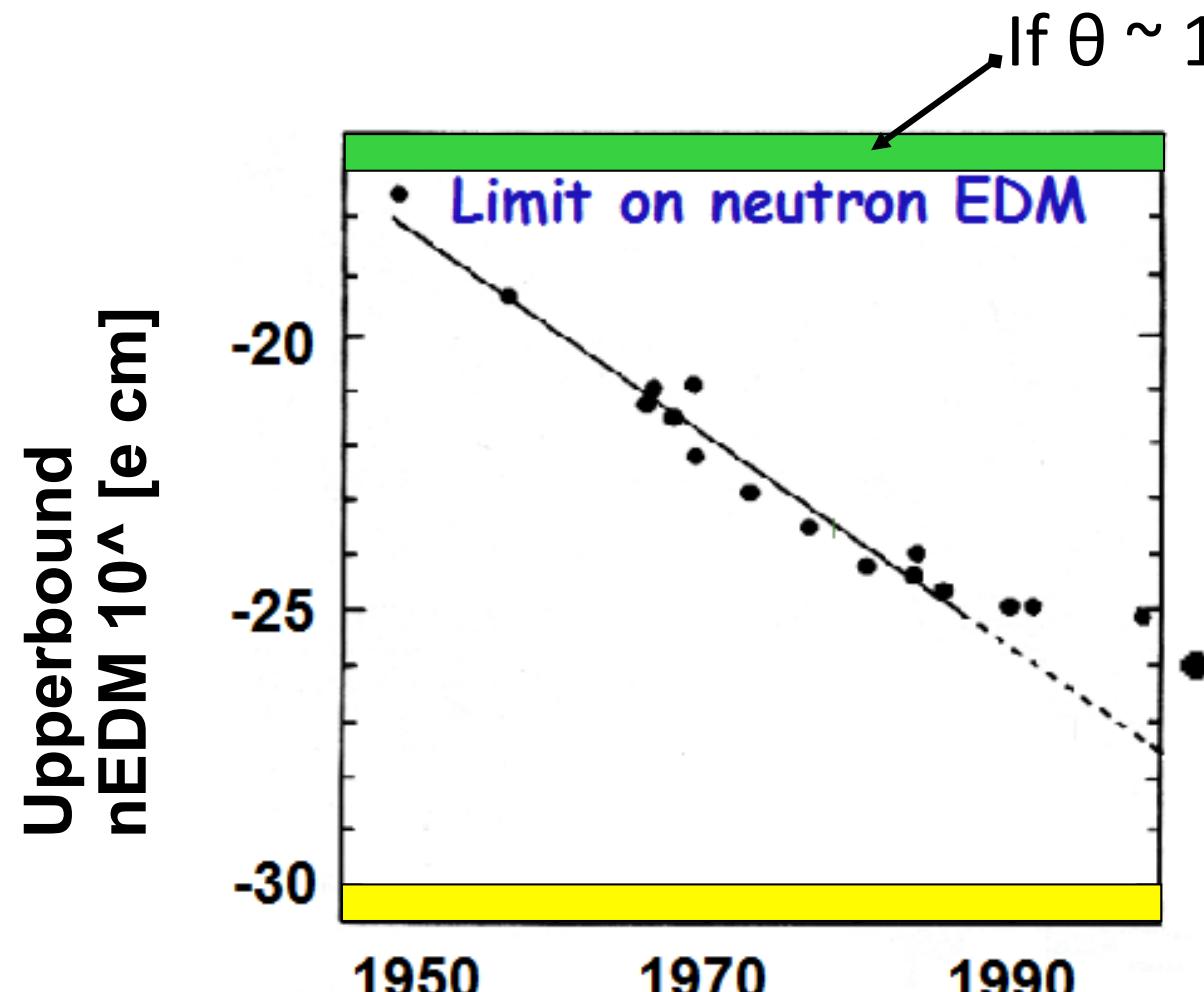


- Causes a ‘new’ CP-violating interaction with **coupling constant θ**

$$\theta \epsilon^{\mu\nu\alpha\beta} G_{\mu\nu} G_{\alpha\beta} \quad (\text{in QED} \sim \vec{E} \cdot \vec{B})$$

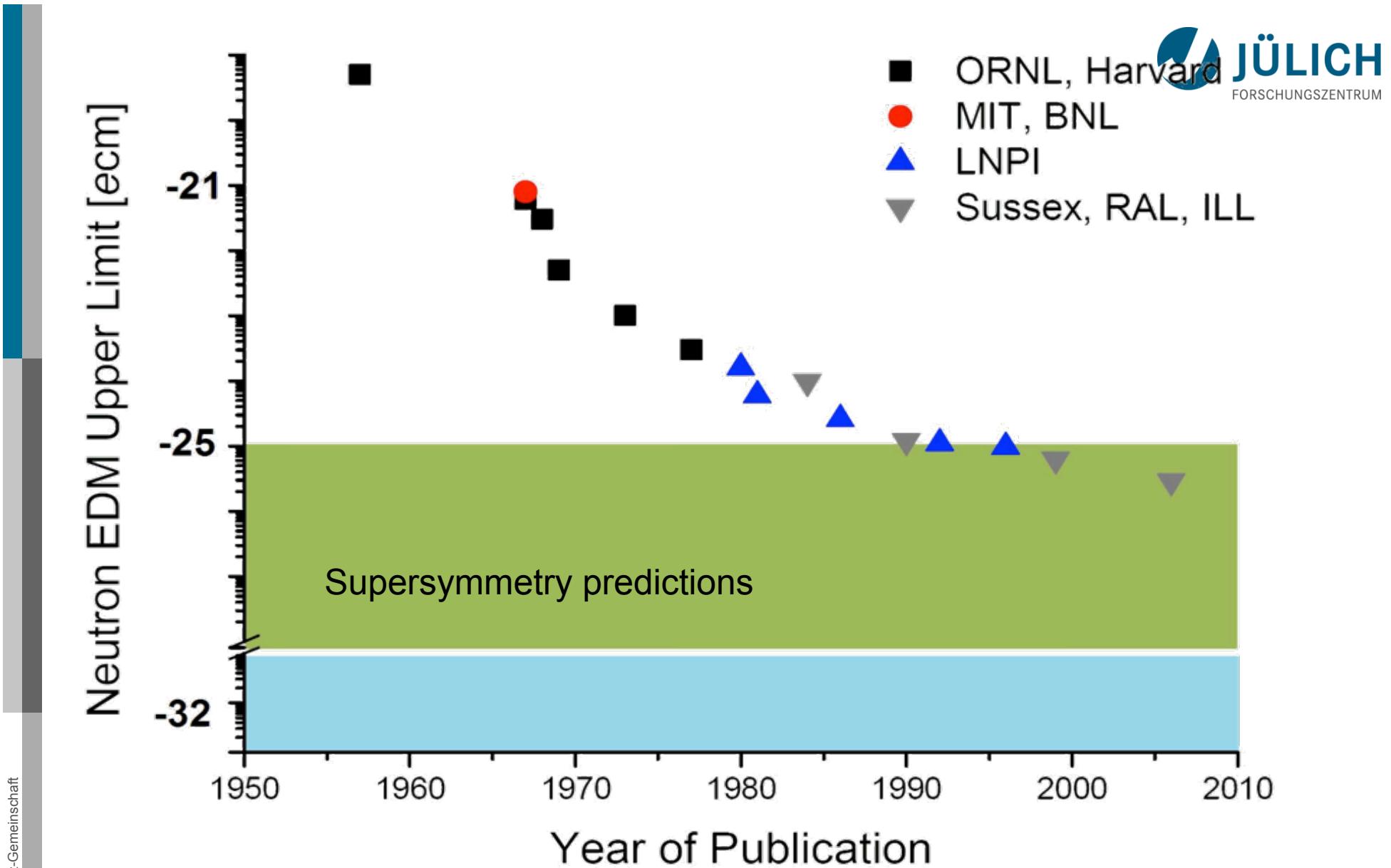
- Size of θ is **unknown, one of SM parameters**

Theta Term Predictions



Crewther et al. (1979)

Sets θ upper bound: $\theta < 10^{-10}$



Electric Dipole Moments =

“the poor man’s high-energy physics” (S. Lamoreaux)

Active experimental field

System	Group	Limit	C.L.	Value	Year
^{205}TI	Berkeley	1.6×10^{-27}	90%	$6.9(7.4) \times 10^{-28}$	2002
YbF	Imperial	10.5×10^{-28}	90	$-2.4(5.7)(1.5) \times 10^{-28}$	2011
$\text{Eu}_{0.5}\text{Ba}_{0.5}\text{TiO}_3$	Yale	6.05×10^{-25}	90	$-1.07(3.06)(1.74) \times 10^{-25}$	2012
PbO	Yale	1.7×10^{-26}	90	$-4.4(9.5)(1.8) \times 10^{-27}$	2013
ThO	ACME	8.7×10^{-29}	90	$-2.1(3.7)(2.5) \times 10^{-29}$	2014
n	Sussex-RAL-ILL	2.9×10^{-26}	90	$0.2(1.5)(0.7) \times 10^{-26}$	2006
^{129}Xe	UMich	6.6×10^{-27}	95	$0.7(3.3)(0.1) \times 10^{-27}$	2001
^{199}Hg	UWash	3.1×10^{-29}	95	$0.49(1.29)(0.76) \times 10^{-29}$	2009
muon	E821 BNL $g-2$	1.8×10^{-19}	95	$0.0(0.2)(0.9) \times 10^{-19}$	2009

Current EDM null results → **probe few TeV scale** or $\phi_{\text{CP}} \leq O(10^{-2,-3})$

Next generation sensitive to **10 TeV (beyond LHC)** or $\phi_{\text{CP}} \leq O(10^{-4,-5})$

Storage rings experiments

$$\frac{d\vec{S}}{dt} = \vec{S} \times \vec{\Omega}$$

Anomalous magnetic moment

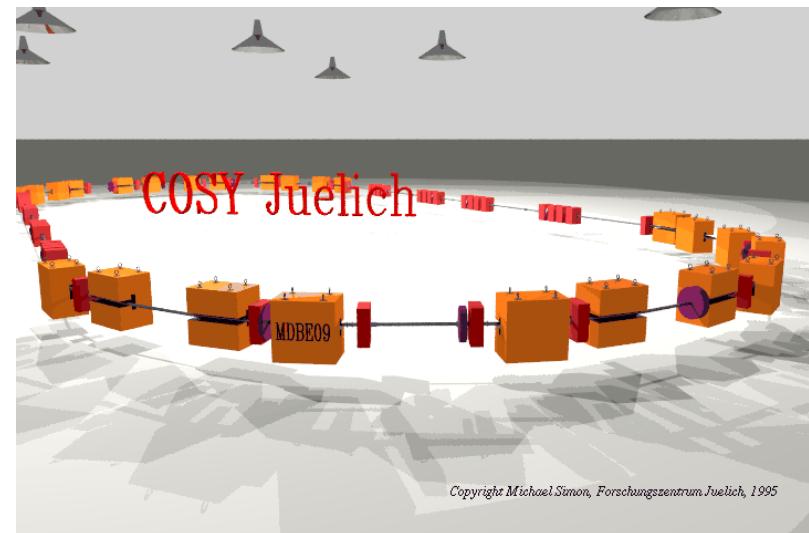
$$\vec{\Omega} = \frac{q}{m} \left[a\vec{B} + \left(\frac{1}{v^2} - a \right) \vec{v} \times \vec{E} \right] + 2d \left(\vec{E} + \vec{v} \times \vec{B} \right)$$

Electric dipole moment

Proton at
Brookhaven/
Fermilab?



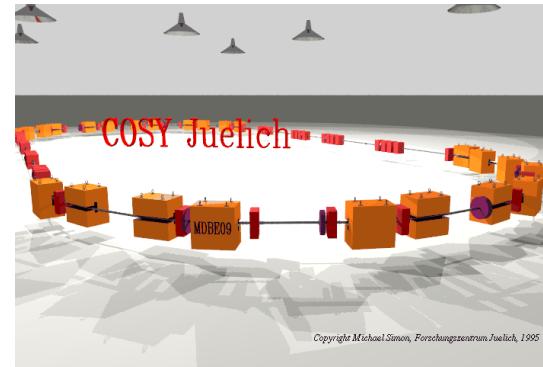
All-purpose ring (proton, deuteron, helion) at COSY?



Experiments on hadronic EDMs

Farley *et al* PRL '04

- New kid on the block: **Charged particle in storage ring**



Bennett *et al* (BNL g-2) PRL '09

- Limit on muon EDM

$$d_\mu \leq 1.8 \cdot 10^{-19} \text{ e cm} \quad (95\% \text{ C.L.})$$

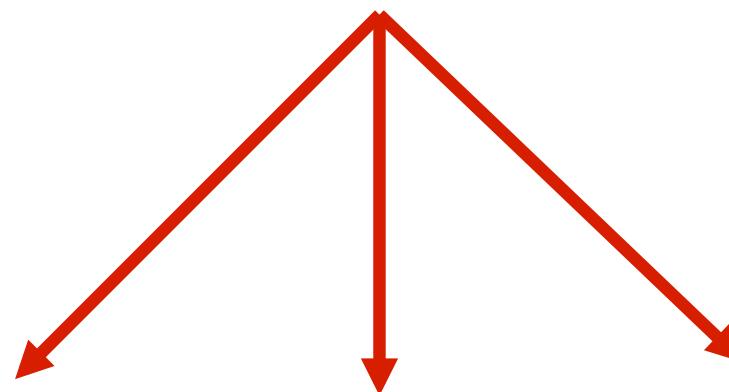
- **Proposals to measure EDMs of proton, deuteron, ^3He at level**
- Other light nuclei?

$$\sim 10^{-29} \text{ e cm}$$

COSY @ Jülich
Brookhaven/Fermilab

Why bother with them ?

Why are EDMs interesting to measure?

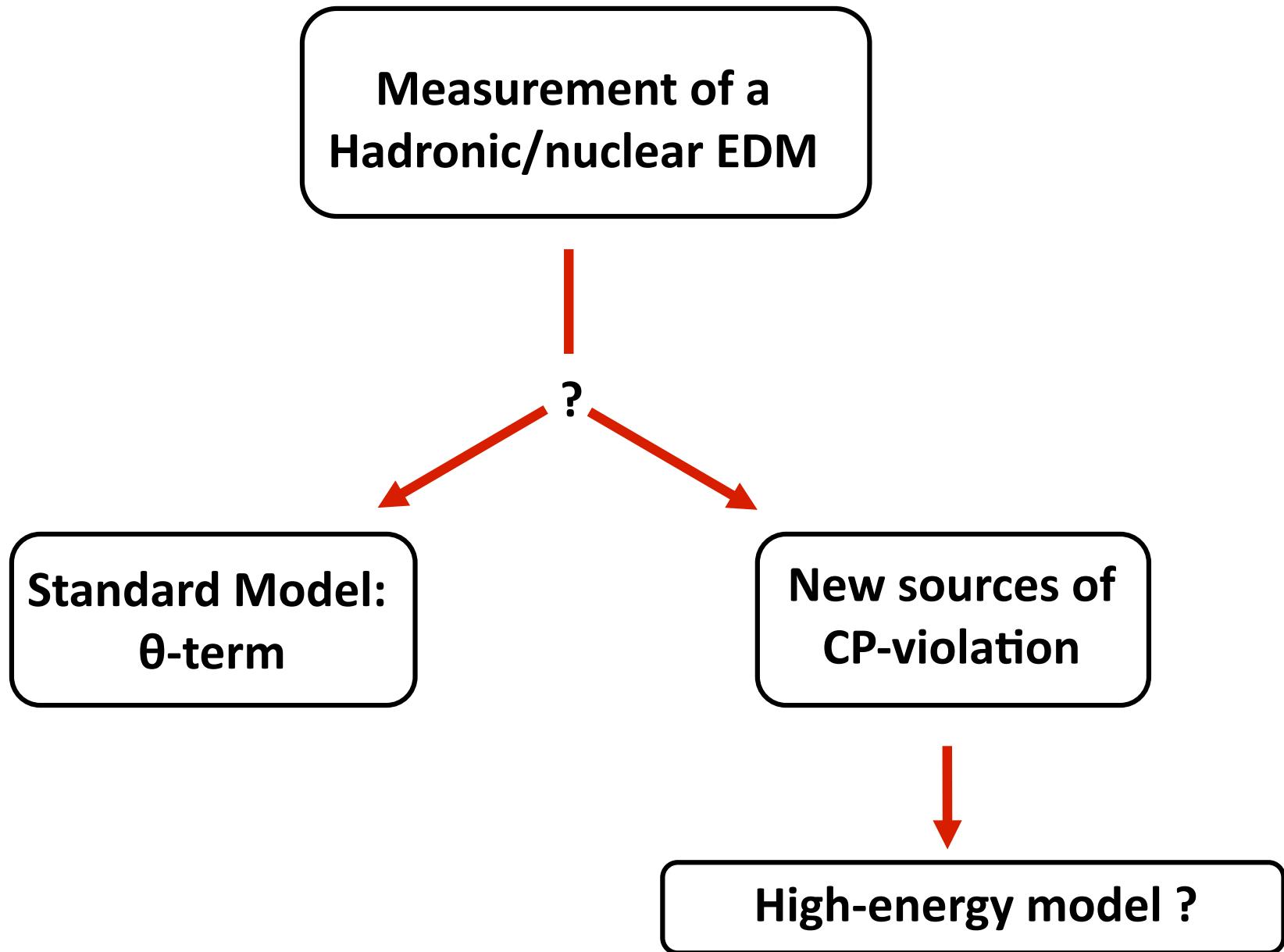


A search for
new physics which is
'background free'

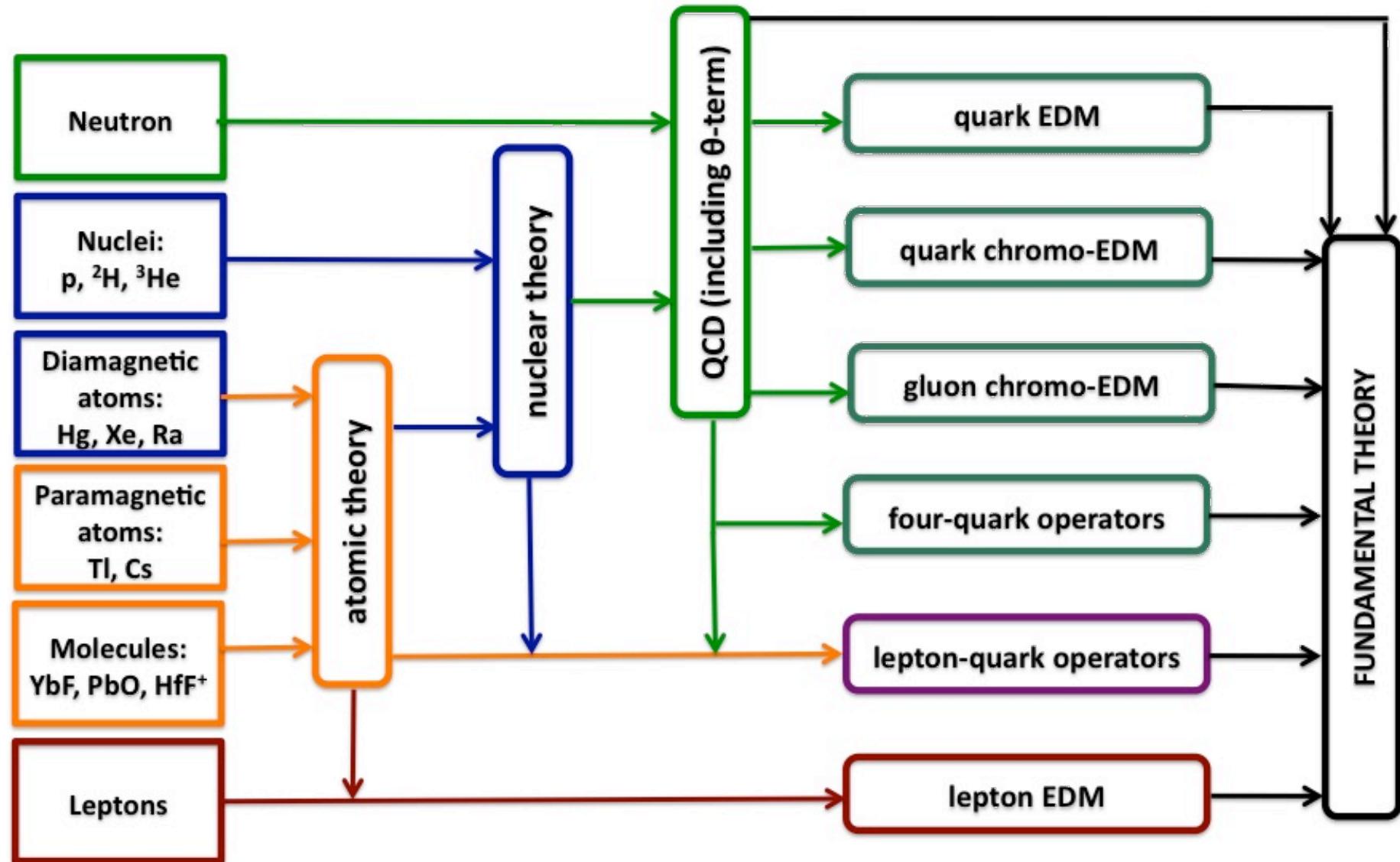
Many beyond-the-
SM models predict
large EDMs:
*Complementary to
LHC search*

Matter/Antimatter
asymmetry
requires more CPV:
*EDMs are excellent
probes*

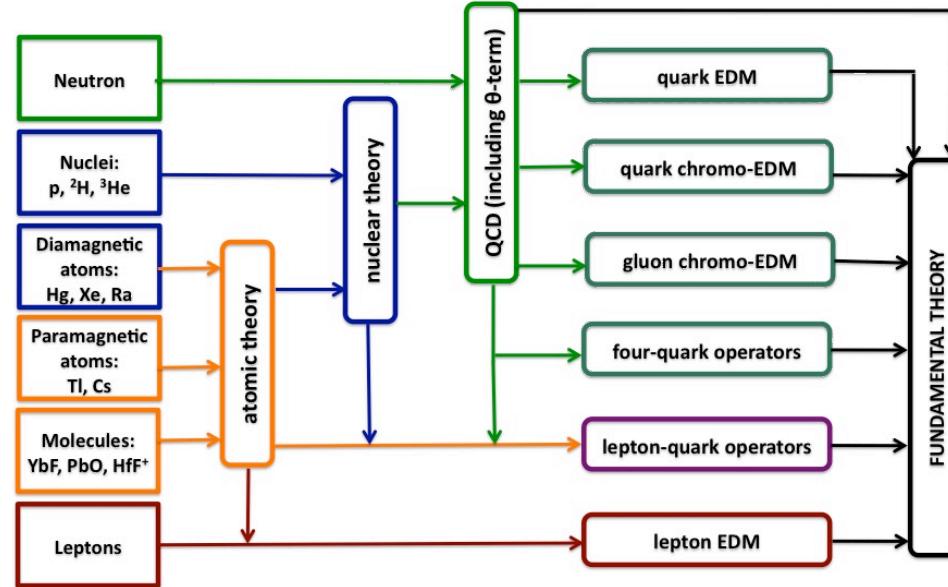
Not really background free.....



Finding the Source



Footprint ?



Questions :

Do **different** models of CP-violation leave
behind a **different** ‘EDM-footprint’



Can we **pinpoint** the microscopic source of CP-violation
from **EDM measurements**?

Outline of this talk

- **Part I:** What are EDMs and why are they interesting in the first place ?
- **Part II: Effective field theory framework**
- **Part III:** EDMs of nucleons, light nuclei, and heavier systems

Separation of scales

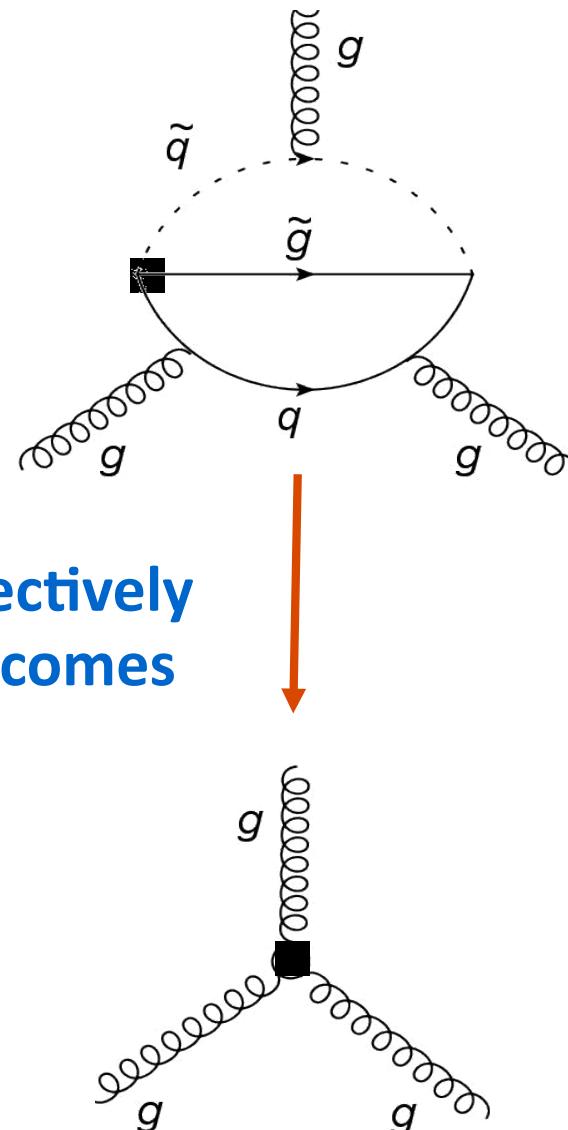
1 TeV ?

SUSY?

100 GeV

Standard
Model

Energy



$$\propto \frac{1}{M_{CP}^2}$$

A systematic approach

- Add all possible CP-odd operators with:

Buchmuller & Wyler NPB '86
 Gradkowski et al JHEP '10

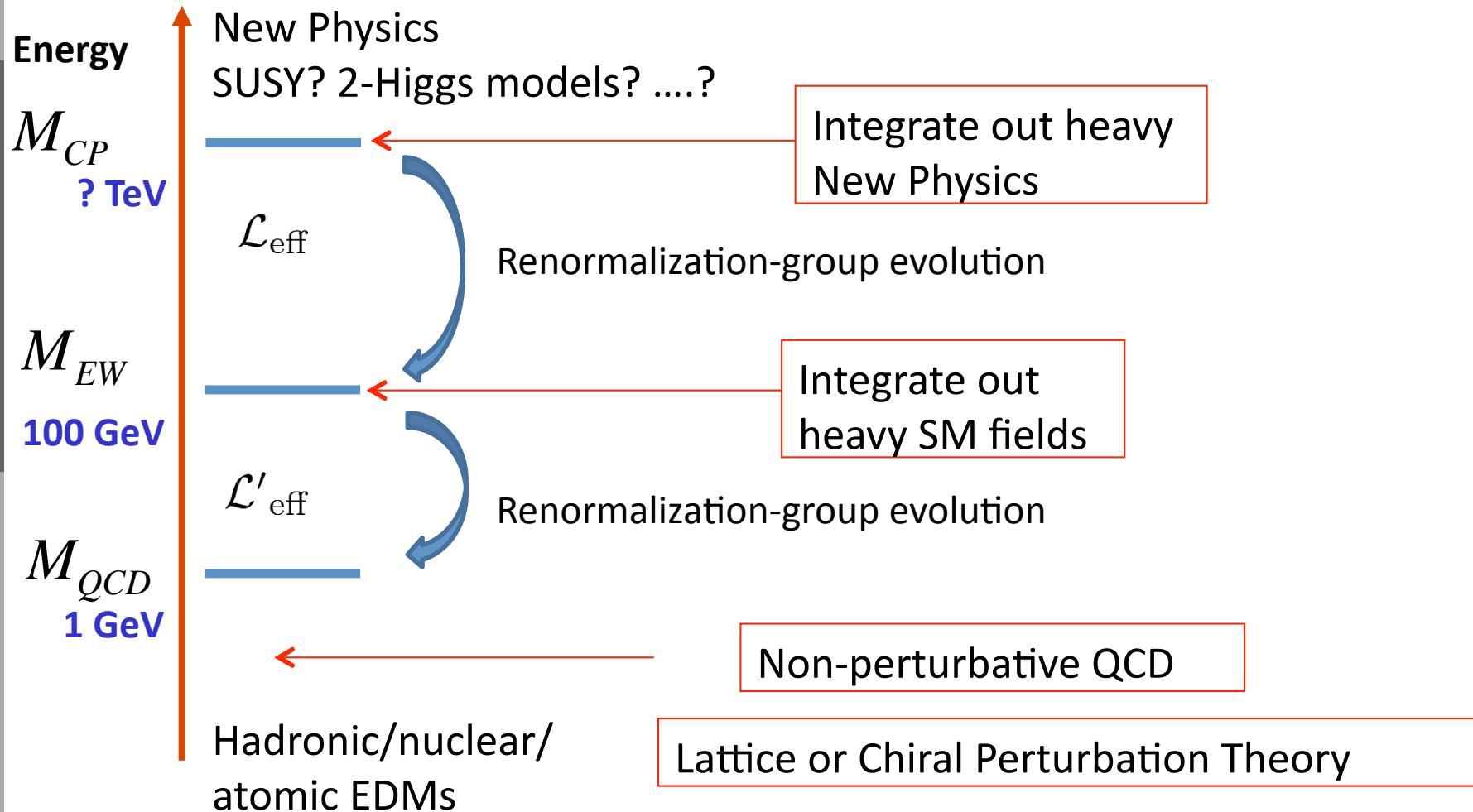
- 1) Degrees of freedom: Full SM field content
- 2) Symmetries: Lorentz, SU(3)xSU(2)xU(1)

$$L_{new} = \cancel{\frac{1}{M_{CP}} L_5} + \frac{1}{M_{CP}^2} L_6 + \dots$$

- Model independent, ***but can be matched to particular models***

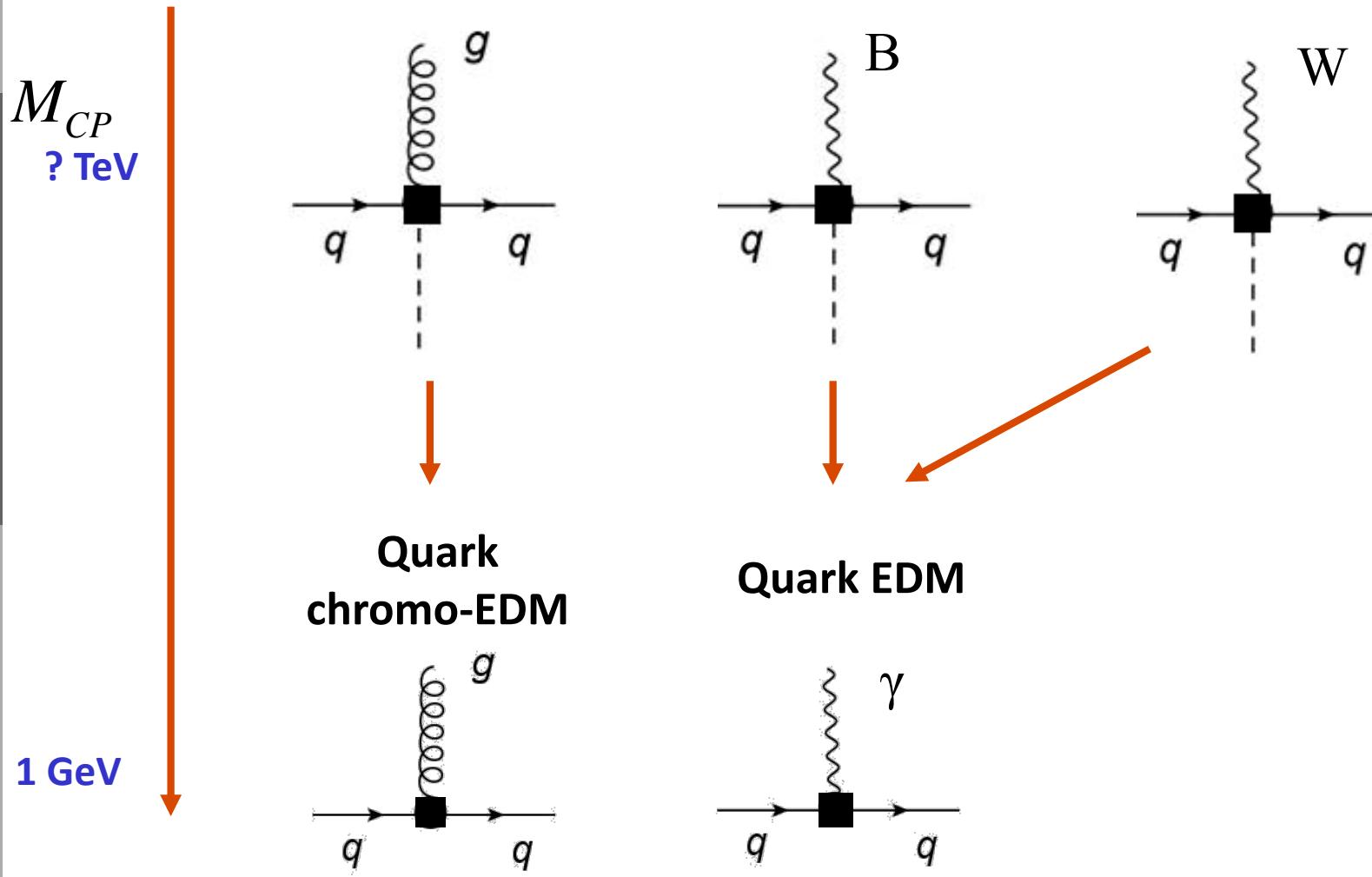
Separation of scales

Dekens & JdV, JHEP, '13



Examples of dim6 operators

Quark-Higgs-Gauge couplings

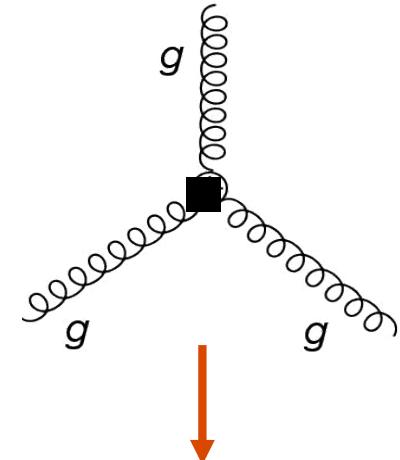


Examples of dim6 operators

Pure gauge operators

Weinberg PRL '89

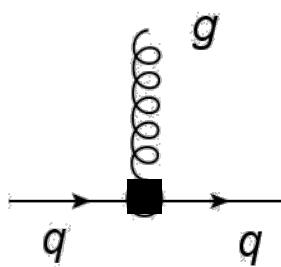
M_{CP}
? TeV



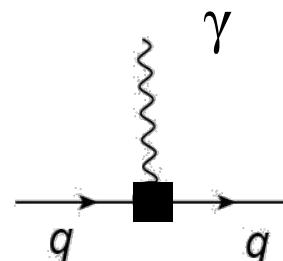
Gluon
chromo-EDM

$$d_w f^{abc} \epsilon^{\mu\nu\alpha\beta} G_{\alpha\beta}^a G_{\mu\lambda}^b G_{\nu}^{c\lambda}$$

Quark
chromo-EDM



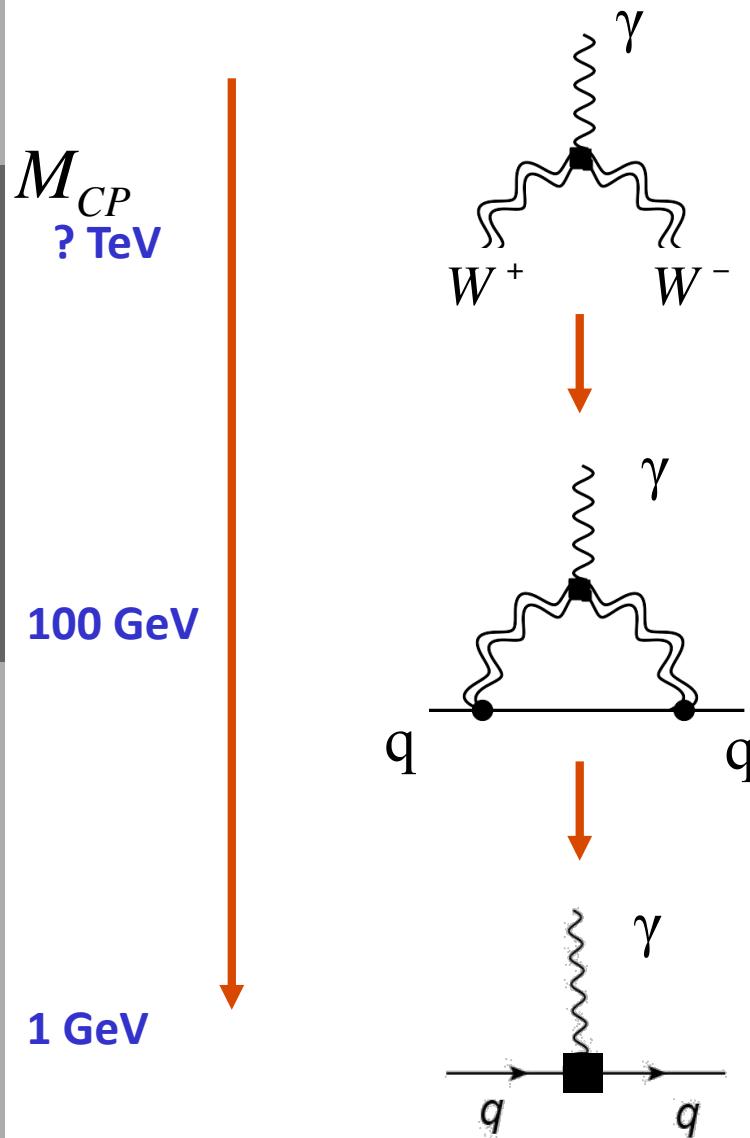
Quark EDM



1 GeV

Complements LHC

Anomalous triple gauge
boson vertices

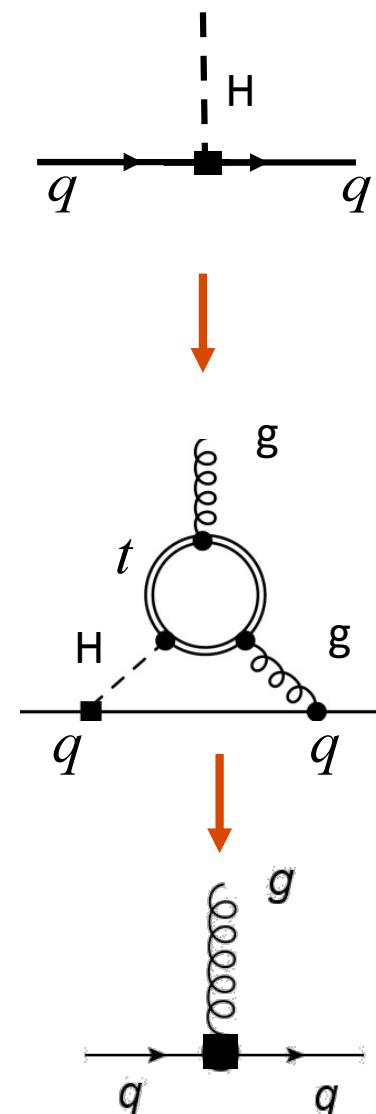


M_{CP}
? TeV

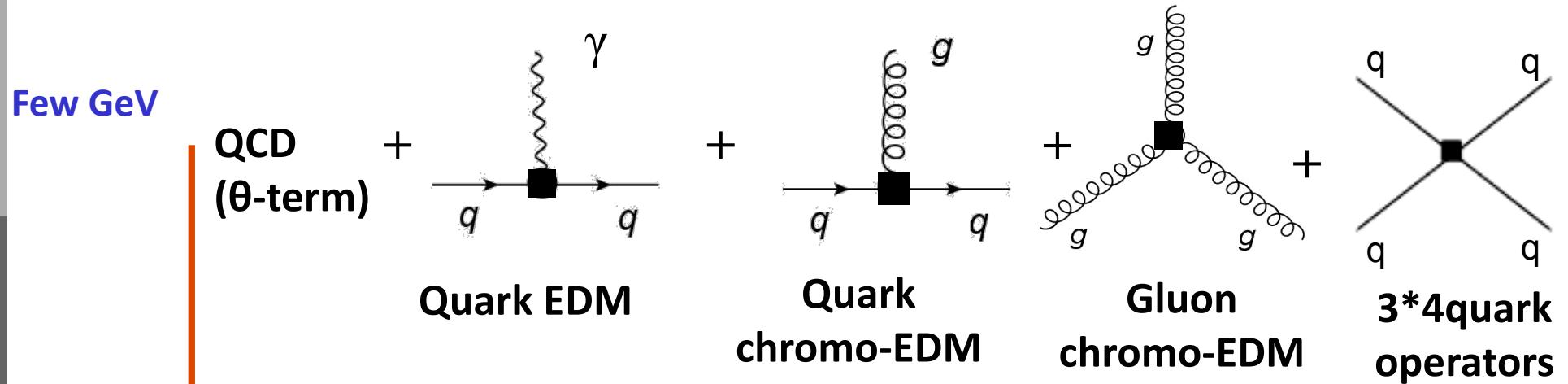
100 GeV

1 GeV

CP-odd quark-Higgs
vertices



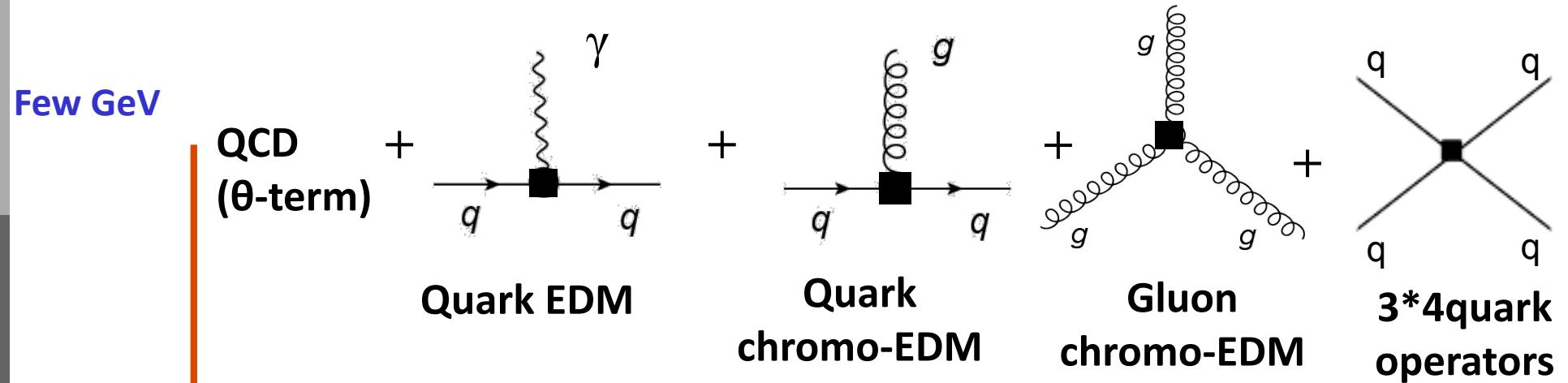
When the dust settles....



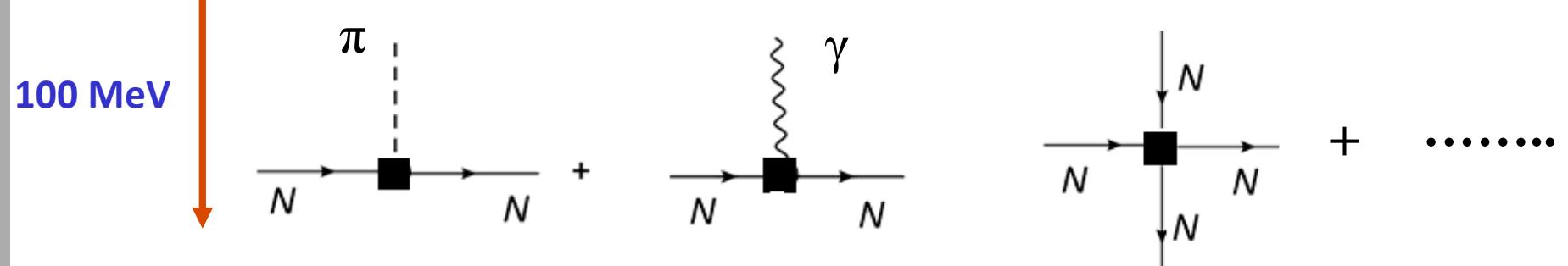
**Different beyond-the-SM models predict
different *dominant operator(s)***

$$\begin{aligned}
 \mathcal{L}_{PT} = & -\bar{\theta} \frac{g^2}{64\pi^2} \epsilon^{\mu\nu\alpha\beta} G_{\mu\nu}^a G_{\alpha\beta}^a - \frac{1}{2} \sum_{q=u,d} \left(d_q \bar{q} i\sigma^{\mu\nu} \gamma_5 q F_{\mu\nu} + \tilde{d}_q \bar{q} i\sigma^{\mu\nu} \gamma_5 t_a q G_{\mu\nu}^a \right) \\
 & + \frac{d_W}{6} f_{abc} \epsilon^{\mu\nu\alpha\beta} G_{\alpha\beta}^a G_{\mu\rho}^b G_{\nu}^{c\rho} + \sum_{i,j,k,l=u,d} C_{ijkl} \bar{q}_i \Gamma q_j \bar{q}_k \Gamma' q_l ,
 \end{aligned}$$

Crossing the barrier



Chiral Perturbation Theory



Chiral effective field theory

- Use the symmetries of QCD to obtain chiral Lagrangians

$$L_{QCD} \rightarrow L_{chiPT} = L_{\pi\pi} + L_{\pi N} + L_{NN} + \dots$$

- Massless QCD Lagrangian has a global $SU(2) \times SU(2)$ symmetry spontaneously broken to $SU(2)$

$$q \rightarrow \left(e^{i \vec{\theta}_V \cdot \vec{\tau} + i \vec{\theta}_A \cdot \vec{\tau} \gamma^5} \right) q \quad q = (u \ d)^T$$

- Pions are the corresponding **Goldstone** bosons (small pion mass due to small quark mass)

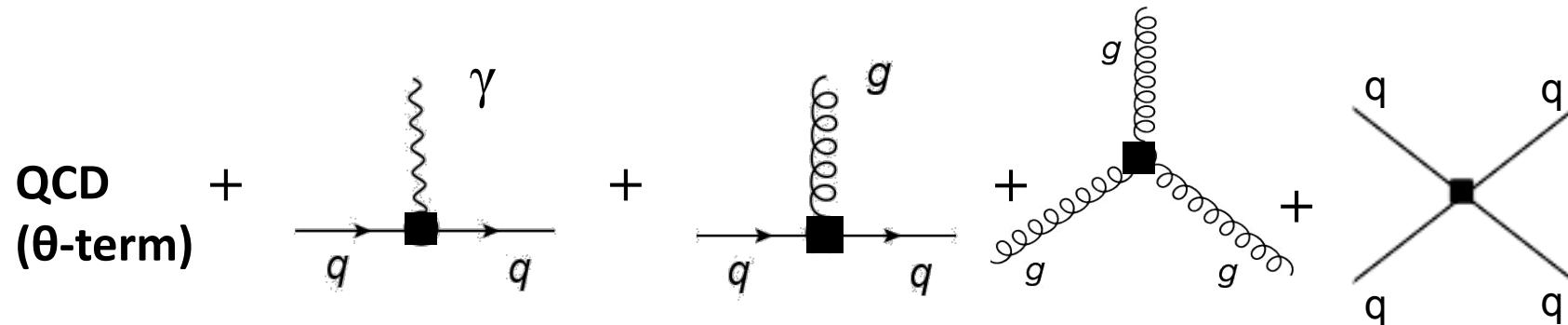
Chiral effective field theory

$$L_{QCD} \rightarrow L_{chiPT} = L_{\pi\pi} + L_{\pi N} + L_{NN} + \dots$$

- The chiral Lagrangian can be constructed **order by order**
- Form of interactions **fixed by symmetry** considerations
- Each interaction associated with **unknown** LEC (needs to be fitted or from **lattice QCD**)
- Provides a **perturbative** expansion in $\frac{Q}{\Lambda_\chi}$ $\Lambda_\chi \cong 1 \text{ GeV}$
- Extremely successful in CP-even case (See for instance talk by Evgeny Epelbaum last monday)

ChiPT with CP violation

- After integrating out new physics (***whatever it is***) at low energies:



- They all break CP
- But **transform differently** under chiral symmetry



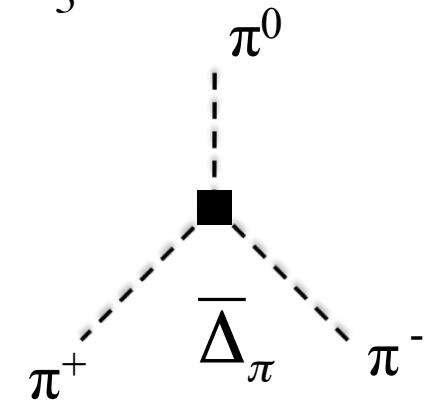
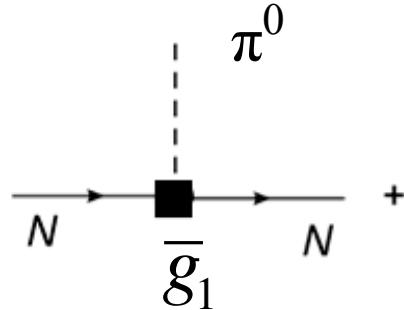
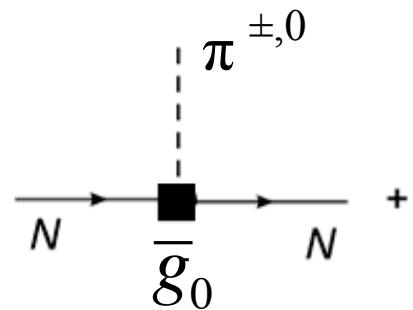
Different CP-odd chiral Lagrangians!

This is the key to unravel them !

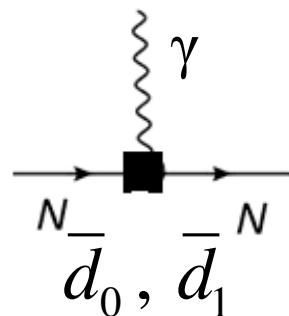
Seven interactions up to N2LO

- We extend chiPT to include **CP-odd interactions**

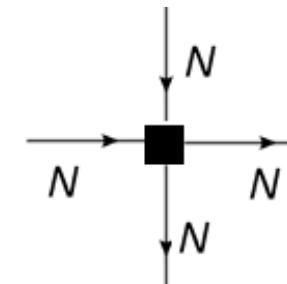
$$L = \bar{g}_0 \bar{N} (\vec{\pi} \cdot \vec{\tau}) N + \bar{g}_1 \bar{N} \pi_3 N + \bar{\Delta}_\pi \pi^2 \pi_3$$



$$L = \bar{d}_0 \bar{N} (\vec{\sigma} \cdot \vec{E}) N + \bar{d}_1 \bar{N} (\vec{\sigma} \cdot \vec{E}) \tau^3 N$$



$$L = \bar{C}_1 (\bar{N} \vec{\sigma} N) \cdot \vec{\partial} (\bar{N} N) + \bar{C}_2 (\bar{N} \vec{\sigma} \tau N) \cdot \vec{\partial} (\bar{N} \tau N)$$



Hierarchy among the sources

- Example: CP-odd pion-nucleon interactions

$$L = \bar{g}_0 \bar{N} (\vec{\pi} \cdot \vec{\tau}) N + \bar{g}_1 \bar{N} \pi_3 N$$

- ❖ θ-term **breaks chiral symmetry** but **conserves isospin symmetry**

- $\bar{g}_0 \gg \bar{g}_1$ because \bar{g}_1 is isospin-breaking

$$\bar{g}_0 = \frac{(m_n - m_p)^{\text{strong}}}{4F_\pi \varepsilon} \bar{\theta} = -0.018(7) \bar{\theta}$$

$$\bar{g}_1 = \frac{8c_1(\delta m_\pi^2)^{\text{strong}}}{F_\pi} \frac{1 - \varepsilon^2}{2\varepsilon} \bar{\theta} = 0.003(2) \bar{\theta}$$

$$\left. \begin{array}{l} \\ \end{array} \right\} \quad \frac{\bar{g}_1}{\bar{g}_0} = -(0.2 \pm 0.1)$$

$$\varepsilon = (m_u - m_d)/2\bar{m} \cong 1/3$$

Hierarchy among the sources

- Example: CP-odd pion-nucleon interactions

$$L = \bar{g}_0 \bar{N} (\vec{\pi} \cdot \vec{\tau}) N + \bar{g}_1 \bar{N} \pi_3 N$$

- ❖ Quark chromo-EDM: **both** isospin-breaking and –conserving part

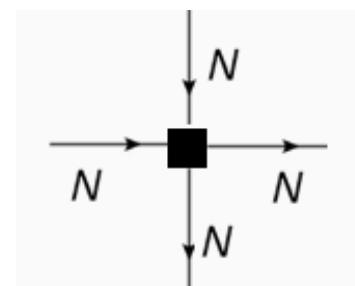
$$|\bar{g}_0| \approx |\bar{g}_1|$$

Relatively **large** uncertainty in LECs
e.g. from QCD sum rules

Pospelov, Ritz '02 '05

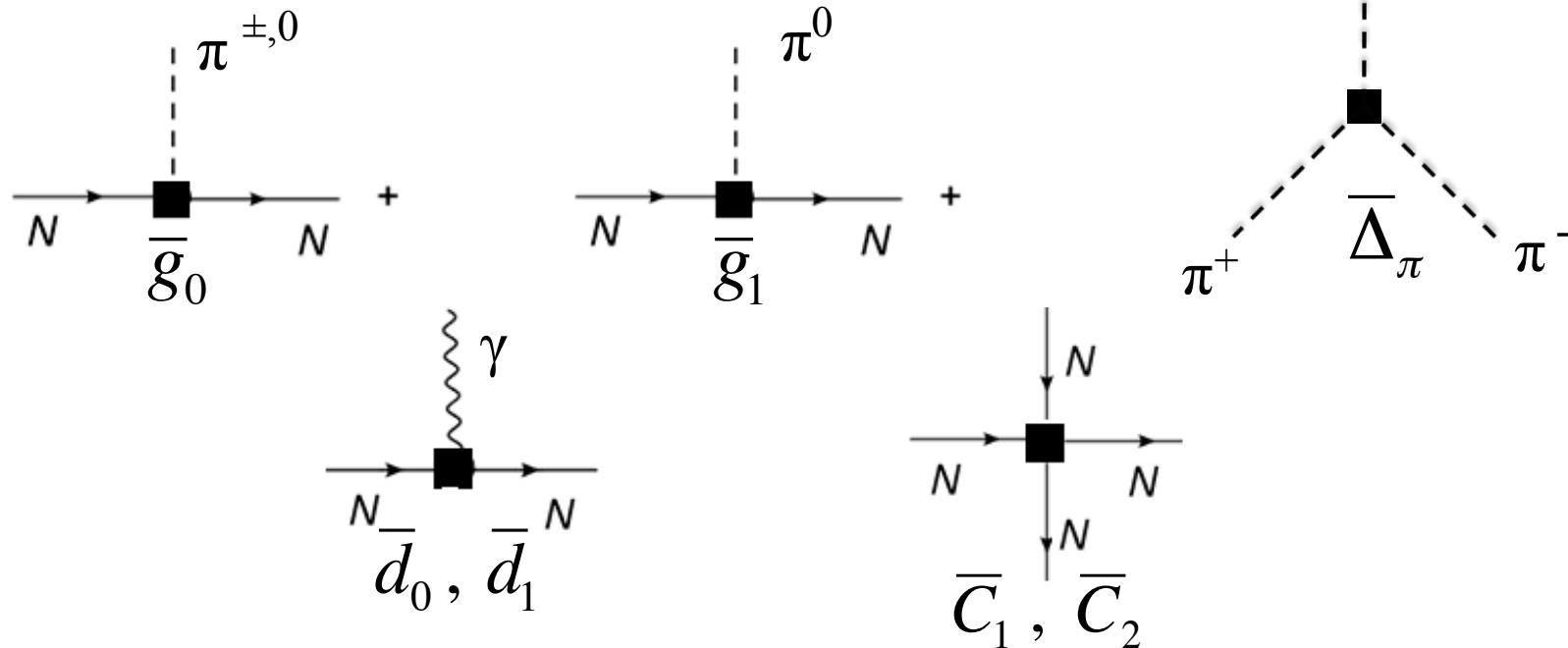
- ❖ Gluon chromo-EDM, LECs suppressed due to **chiral symmetry**.
Contributions from short-range NN interactions.

$$L = \bar{C} (\bar{N} \vec{\sigma} N) \cdot \vec{\partial} (\bar{N} N)$$



JdV, Mereghetti et al '12

The magnificent seven



Different sources of CP-violation (theta, quark EDM, gluon CEDM....)

Induce different hierarchies between these interactions



Probe these hierarchies experimentally

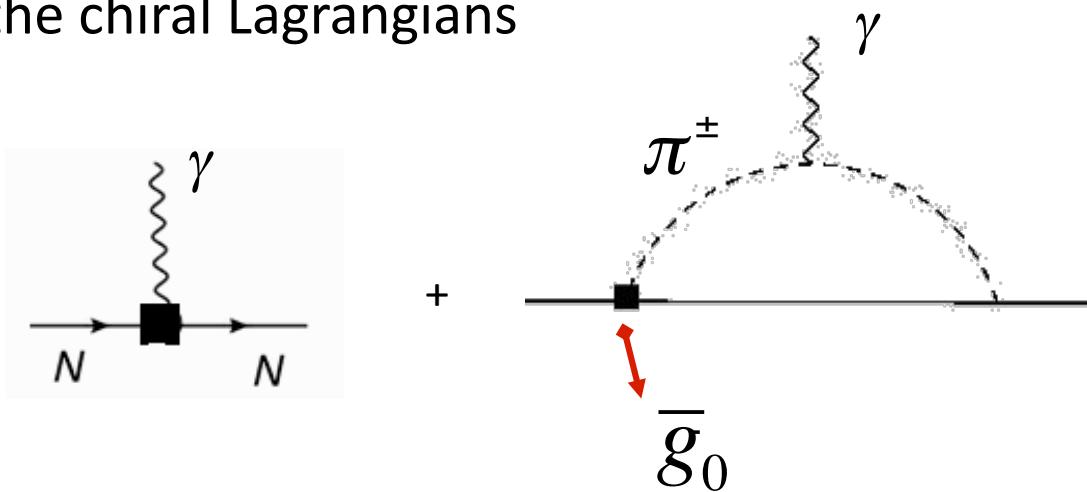
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The Nucleon EDM

- Calculate directly from the chiral Lagrangians

Nucleon EDM



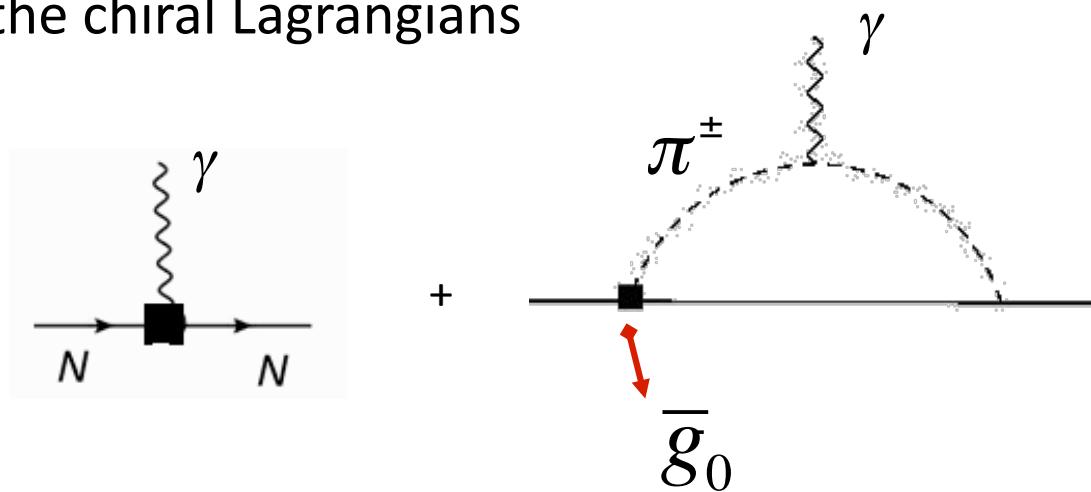
$$d_n = \bar{d}_0 - \bar{d}_1 - \frac{e g_A \bar{g}_0}{4\pi^2 F_\pi} \left(\ln \frac{m_\pi^2}{M_N^2} - \frac{\pi}{2} \frac{m_\pi}{M_N} \right)$$

$$d_p = \bar{d}_0 + \bar{d}_1 + \frac{e g_A}{4\pi^2 F_\pi} \left[\bar{g}_0 \left(\ln \frac{m_\pi^2}{M_N^2} - 2\pi \frac{m_\pi}{M_N} \right) - \bar{g}_1 \frac{\pi}{2} \frac{m_\pi}{M_N} \right]$$

The Nucleon EDM

- Calculate directly from the chiral Lagrangians

Nucleon EDM



$$d_n = \bar{d}_0 - \bar{d}_1 + (0.13 + 0.01) \bar{g}_0 e \text{ fm}$$

$$d_p = \bar{d}_0 + \bar{d}_1 - (0.13 + 0.03) \bar{g}_0 e \text{ fm}$$

**Very similar and
3 LECs...**

Hard to probe the hierarchy with only neutron and proton EDMs

Lattice QCD to the rescue

Test for the QCD theta term. With lattice input for LECs

$$d_n = (2.7 \pm 1.2) \cdot 10^{-16} \overline{\theta} e \text{ cm}$$

Shintani '12 '13
Guo, Meißner '13 '14

$$d_p = -(2.1 \pm 1.2) \cdot 10^{-16} \overline{\theta} e \text{ cm}$$

Could provide **strong hint** for theta term!

New lattice approach for theta

Talk by Andrea Shindler

Lattice QCD to the rescue

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Shintani '12 '13
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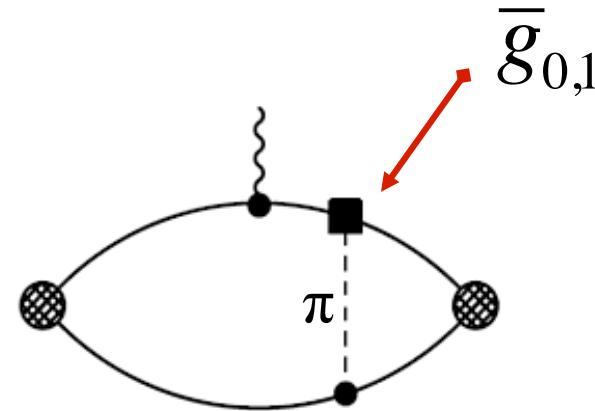
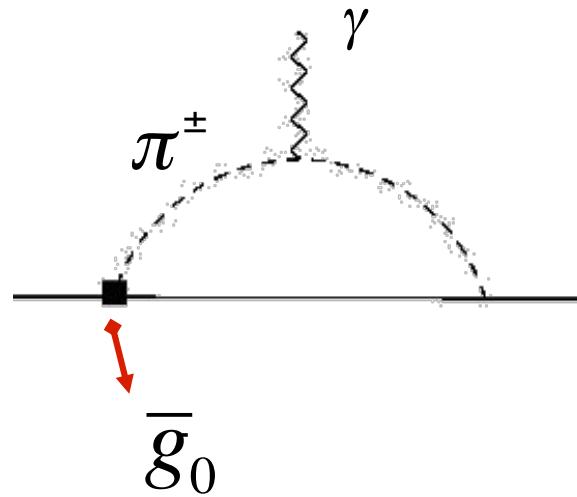
But..... No lattice data yet for other CPV scenarios

Work in progress (e.g. Cirigliano et al , Shindler et al)

Generally: Need more observables to unravel sources !

Nuclear EDMs

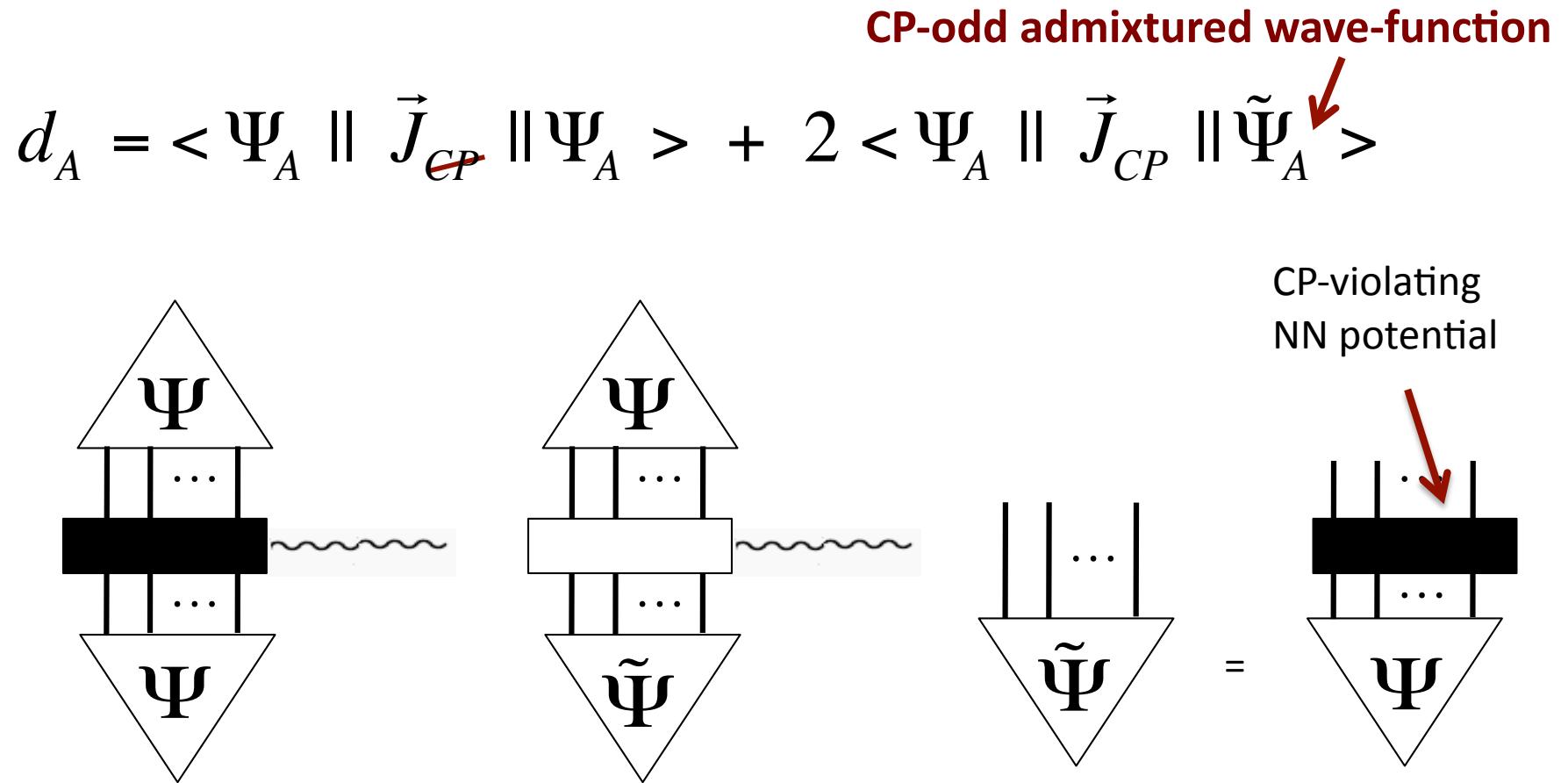
- Nuclear EDMs: **tree-level dependence !**



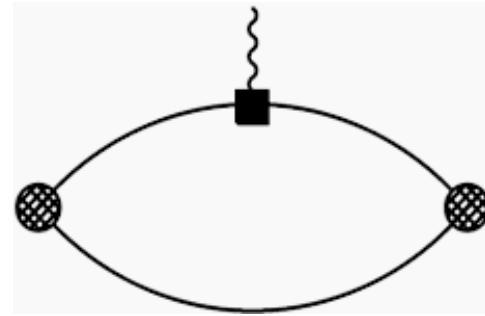
- **No loop suppression and no counter terms!**
- Possible to calculate light-nuclear EDMs with **high accuracy!**

EDM of a general light nucleus

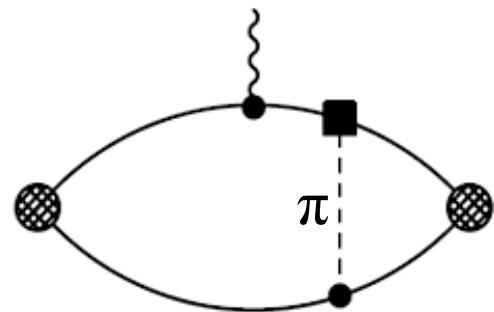
- EDM of a nucleus with A nucleons can be separated in 2 contributions



The deuteron EDM



One-body: $d_D = d_n + d_p$

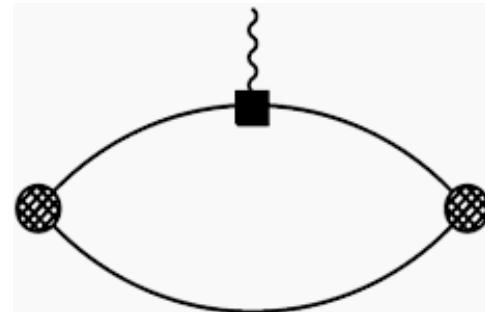


T-violating pion-exchange

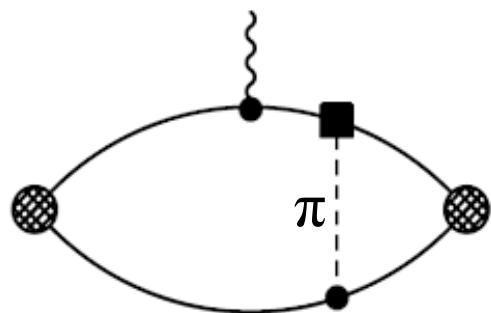
$$L = \bar{g}_0 \overline{N} (\vec{\pi} \cdot \vec{\tau}) N + \bar{g}_1 \overline{N} \pi_3 N$$

Khriplovich+Korkin NPA '00
Liu+Timmermans PRC '04
JdV et al PRL '11
Bsaisou et al '12

The deuteron EDM



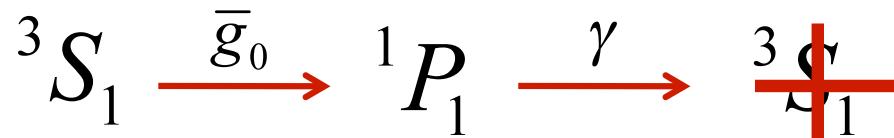
One-body: $d_D = d_n + d_p$



T-violating pion-exchange

$$L = \bar{g}_0 \overline{N} (\vec{\pi} \cdot \vec{\tau}) N + \bar{g}_1 \overline{N} \pi_3 N$$

- Deuteron is a special case due to $N=Z$



The deuteron EDM

Obtain deuteron wave function from **chiral EFT potential**

$$(E - H_{PT}) |\Psi_A\rangle = 0$$

$$(E - H_{PT}) |\tilde{\Psi}_A\rangle = V_{CP} |\Psi_A\rangle$$



Both consistently derived
in chiral EFT

The deuteron EDM

Obtain deuteron wave function from **chiral EFT potential**

$$(E - H_{PT}) |\Psi_A\rangle = 0$$

$$(E - H_{PT}) |\tilde{\Psi}_A\rangle = V_{CP} |\Psi_A\rangle$$

Both consistently derived
in chiral EFT

Strong isospin filter

$$d_D = d_n + d_p + [(0.18 \pm 0.02) \bar{g}_1 + (0.0028 \pm 0.0003) \bar{g}_0] e \text{ fm}$$

Theoretical accuracy is excellent !

Do the same for ${}^3\text{He}$ and ${}^3\text{H}$

No isospin filter in these nuclei, both g_0 and g_1 dependence

$$d_{^3\text{He}} = 0.9 d_n - 0.05 d_p + [(0.14 \pm 0.03) \bar{g}_1 + (0.10 \pm 0.03) \bar{g}_0] e \text{ fm}$$

$$d_{^3\text{H}} = -0.05 d_n + 0.9 d_p + [(0.14 \pm 0.03) \bar{g}_1 - (0.10 \pm 0.03) \bar{g}_0] e \text{ fm}$$

Still good accuracy, can most likely be improved by using N3LO

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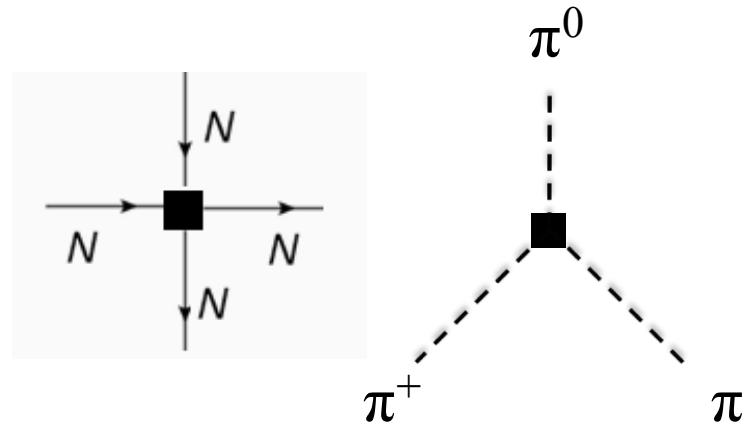
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Still good accuracy, can most likely be improved by using N3LO

^3H probably not a candidate for storage rings

And dependence on:

Not discussed now



What would this tell us ?

- From measurements of dn, dp, dD, and/or d3He we can:
 - 1) Extract the (relative) sizes of the couplings \bar{g}_1 \bar{g}_0 ($\bar{\Delta}_\pi$)
 - 2) The size of the two- and three-body contributions, e.g.

$$\left| \frac{d_D}{d_p} \right| \quad \left| \frac{d_D - d_n - d_p}{d_n + d_p} \right| \quad \left| \frac{d_{^3He}}{d_p} \right| \quad \left| \frac{d_{^3He} - d_n}{d_n} \right|$$

- These quantities point towards underlying source
(theta, quark chromo-EDM, etc)

Heavier EDMs

- **Can't we get this info from EDMs of Hg, Ra, Rn.... ??**

Strong bound on atomic EDM: $d_{^{199}Hg} < 3.1 \cdot 10^{-29} e \text{ cm}$

Griffiths et al, PRL '09

- The **atomic** part of the calculation is well under control

$$d_{^{199}Hg} = (2.8 \pm 0.6) \cdot 10^{-4} S_{Hg} e \text{ fm}^2$$

Dzuba et al, PRA '02, '09



Nuclear Schiff moment

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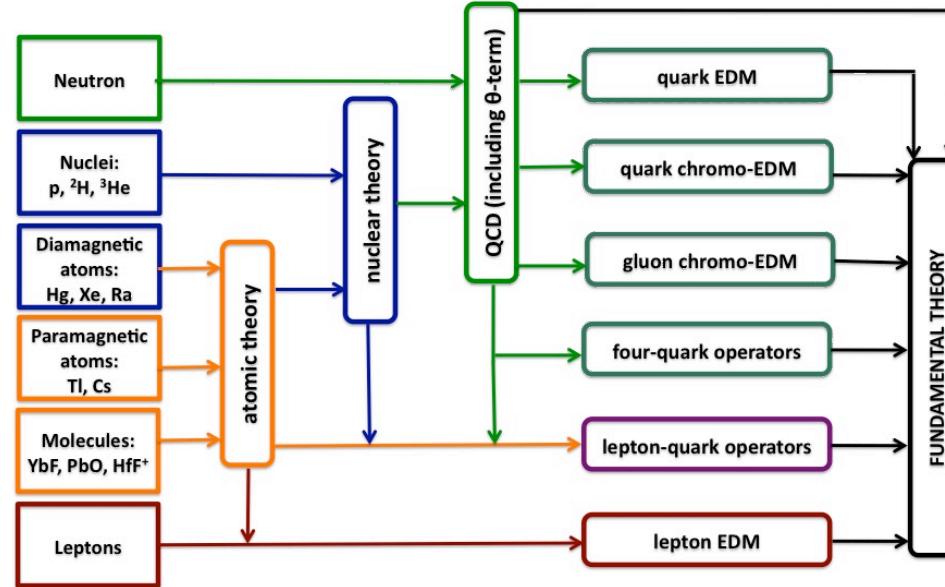
- But the **nuclear** part isn't.....

$$S_{^{199}Hg} = [(0.3 \pm 0.4)\bar{g}_0 + (0.4 \pm 0.8)\bar{g}_1] e \text{ fm}^3$$

Engel et al, PPNP '13

- There is no **power counting** for nuclei with so many nucleons

Footprint ?



Big questions :

Do **different** models of CP-violation leave
behind a **different** ‘EDM-footprint’



Can we **pinpoint** the microscopic source of CP-violation
from **EDM measurements**?

Unraveling models of CPV

- In recent work we studied 4 popular scenarios of CP-violation
 - 1) **Standard Model including QCD theta term**
 - 2) **The minimal left-right symmetric model** Mohapatra et al '08
 - 3) The aligned two-Higgs-Doublet model Pich & Jung '13
 - 4) The MSSM (well, specific versions of it)
- Can we **unravel** these models with EDM experiments? Dekens, JdV, Bsaisou, et al, JHEP '14

Two scenarios of CPV

1) We first look at the QCD theta term

Crewther et al. (1979)

$$L = L_{QCD} + \frac{\theta}{32\pi^2} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

$$\bar{m} = (m_u + m_d)/2$$



Axial U(1) transformation

$$L = L_{QCD} + i \frac{\bar{m}}{2} \theta \bar{q} i\gamma^5 q$$

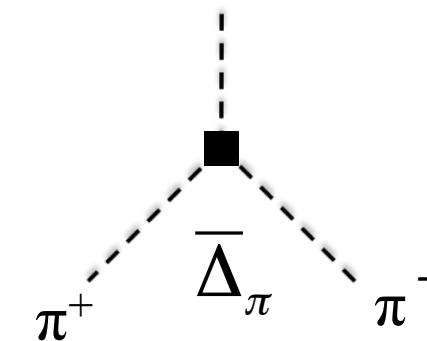
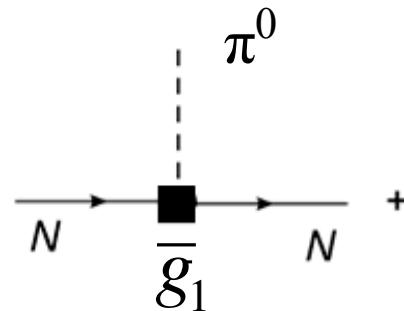
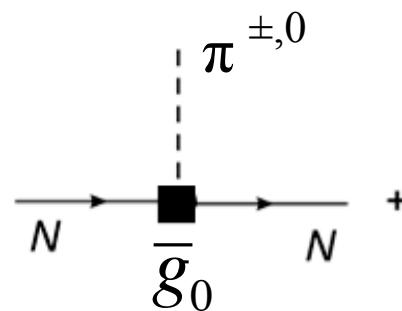


a CP-odd quark mass
Isospin-symmetric!!

Leads to a very **specific** hadronic CP-odd Lagrangian

Two scenarios of CPV

1) Chiral Lagrangian for theta term



$$\bar{g}_0 = \frac{(m_n - m_p)^{\text{strong}}}{4F_\pi \varepsilon} \bar{\theta} = -0.018(7) \bar{\theta}$$

$$\bar{g}_1 = \frac{8c_1(\delta m_\pi^2)^{\text{strong}}}{F_\pi} \frac{1 - \varepsilon^2}{2\varepsilon} \bar{\theta} = 0.003(2) \bar{\theta}$$

$$\varepsilon = (m_u - m_d)/2\bar{m} \cong 1/3$$

Accuracy will improve!

$$\frac{\bar{g}_1}{\bar{g}_0} = -(0.2 \pm 0.1)$$

$\bar{\Delta}_\pi$ **Is N2LO effect**

Two scenarios of CPV

2) Now the minimal left-right symmetric model

Mohapatra et al '08

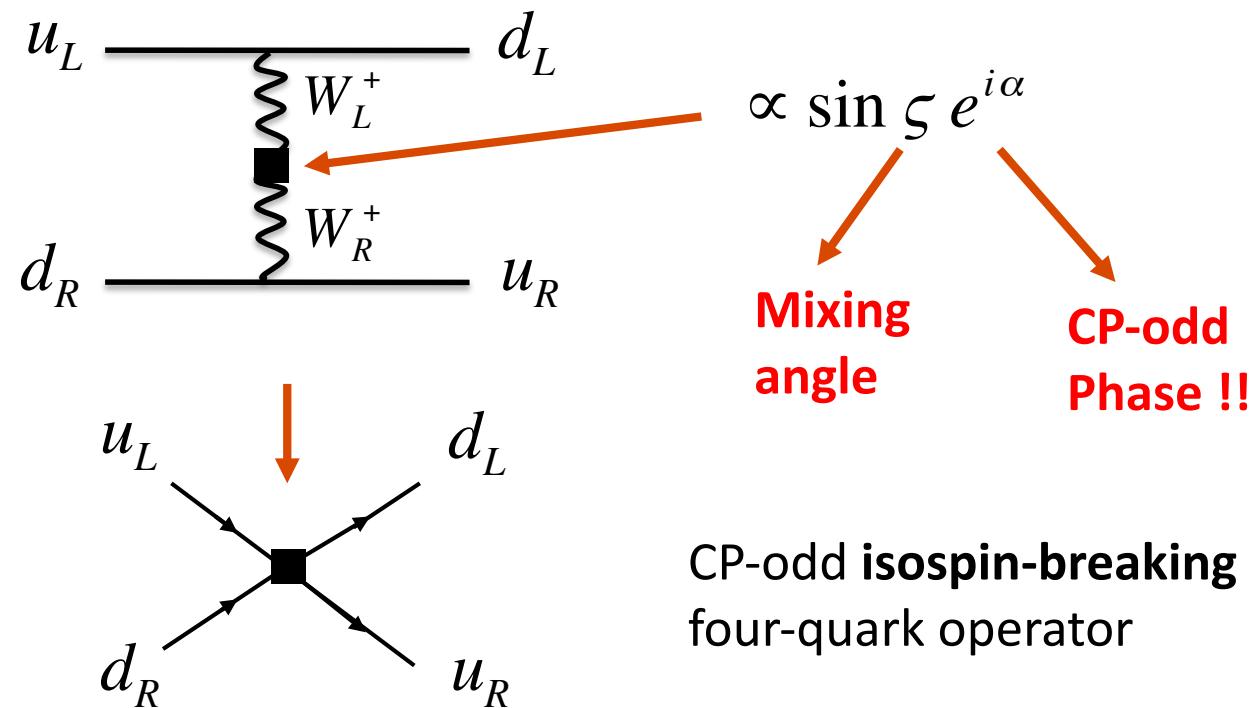
- Based on unbroken Parity symmetry at high energies
- Gauge group: $SU_R(2) \times SU_L(2) \times SU_c(3) \times U(1)$
- Additional Higgs fields to break $SU_R(2)$
- Additional **heavy** right-handed gauge bosons W_R^{+-} and Z_R

Two scenarios of CPV

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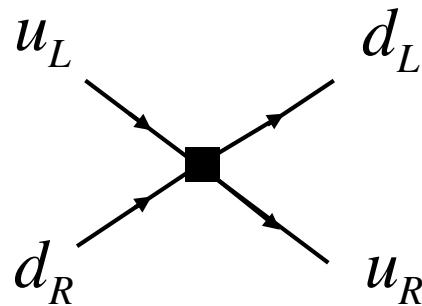
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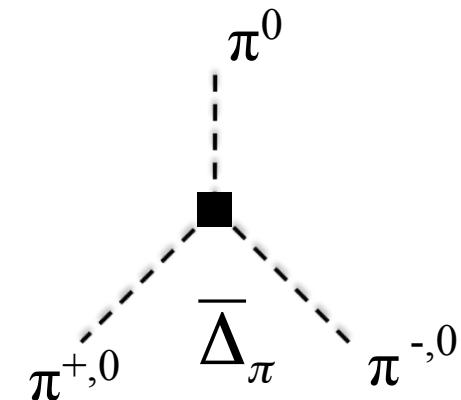
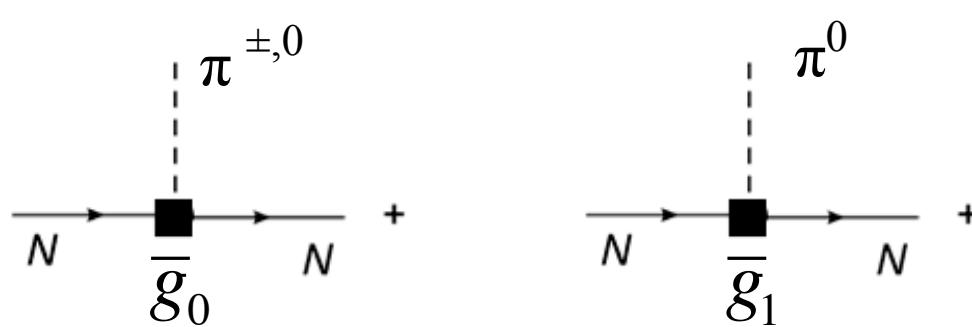


Two scenarios of CPV

2) minimal left-right symmetric model



CP-odd **isospin-breaking**
four-quark operator



Completely **opposite** behavior with respect to theta term

$$\frac{\bar{g}_0}{\bar{g}_1} = \frac{(m_n - m_p)^{str}}{8c_1 m_\pi^2} - (0.02 \pm 0.01)$$

Δ_{π} **Leading order interaction!**

“EDM sum rules”

1) Theta term

$$\frac{\bar{g}_1}{\bar{g}_0} = -(0.2 \pm 0.1)$$

2) mLRSM

$$\frac{\bar{g}_0}{\bar{g}_1} = -(0.02 \pm 0.01)$$

Deuteron EDM, sensitive mainly to \bar{g}_1

$$d_D = d_n + d_p + [(0.18 \pm 0.02) \bar{g}_1 + (0.0028 \pm 0.0003) \bar{g}_0] e \text{ fm}$$

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$$\left| \frac{d_D - d_n - d_p}{d_n} \right| < 1$$

$$d_D - d_n - d_p \cong \frac{d_n}{6}$$

$$\left| \frac{d_D - d_n - d_p}{d_n} \right| > 1$$

Rather big
Uncertainty

$$d_D - d_n - d_p \cong (3 - 10) d_n$$

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$$d_{^3\text{He}} - 0.9d_n \cong \frac{d_n}{2}$$

$$\cong 3(d_D - d_n - d_p)$$

$$d_{^3\text{He}} - 0.9d_n \cong 0.7 d_D$$

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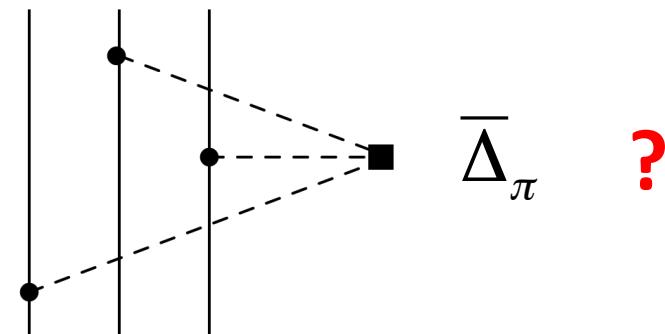
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$$\begin{aligned} d_{^3\text{He}} - 0.9d_n &\cong \frac{d_n}{2} \\ &\cong 3(d_D - d_n - d_p) \end{aligned}$$

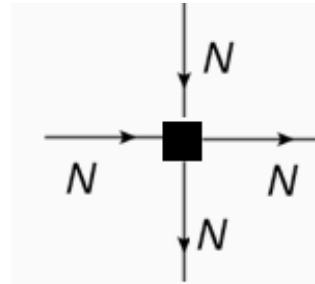
$$d_{^3\text{He}} - 0.9d_n \cong 0.7 d_D$$



Unraveling models of CPV

- Multi-Higgs and the MSSM leave different hierarchies behind

$$\left| \frac{g_1}{g_0} \right| \approx 1$$



- And give rise do different **electron/nucleon EDM ratio**

- **EDM experiments can tell us a lot about new physics**

(or its absence....)

To-do list

Light nuclei

Wrap up some things (i.e. 3-body interactions)

Reduce hadronic uncertainties

Lattice QCD determinations of CP-odd LECs from theta and BSM.
In particular: nucleon EDMs and pion-nucleons

Increase the number of nucleons

Can we extend the framework to heavier systems ?

$$d_A(d_n, d_p, \bar{g}_0, \bar{g}_1, \bar{\Delta}_\pi, \bar{C}_{1,2})$$

Perhaps with nuclear lattice EFT ?

Conclusion/Summary

- 1) EDMs are very good probes of new CP-odd physics
- 2) Probe **similar energy scales** as LHC, strong bounds on new physics

EFT approach

- a) Framework exists for CP-violation (EDMs) from 1st principles
- b) Keep track of **symmetries** from multi-Tev to few MeV

The chiral filter

- a) Chiral symmetry determines form of hadronic interactions
- b) n, p, D, 3He can be used to **unravel** the CP-odd source (if existing..)

Quantified uncertainties

- a) Hadronic uncertainties dominate for nucleons & light nuclei → LQCD
- b) Nuclear uncertainties dominate for heavy systems