STORI'14: 9th International Conference on Nuclear Physics at Storage Rings

September 28th – October 3rd, 2014 St. Goar (Germany)

SPIN COHERENCE TIME studies

of a POLARIZED DEUTERON beam @ COSY

On behalf of the JEDI Collaboration

G. Guidoboni University of Ferrara and INFN

See "Electrice Dipole Moment measurements at storage rings" by J. Pretz

Electric Dipole Moment of fundamental particles Origin Now matter Matter Anti-matter Matter BARYOGENESIS Sakharov's conditions (1967): $B = n_b - n_{\overline{b}}$ violation C and *CP* violation Far from thermal equilibrium

Standard Model (SM):

- Not enough to explain *Baryon Asymmetry*
- Too small CP violation

Electric Dipole Moment of fundamental particles



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Def: permanent charge displacement within the particle volume



Assuming CPT symmetry T violation = **CP violation**

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Assuming CPT symmetry

T violation = **CP violation**

 EDM_{SM} too small to be observed EDM_{bSM} within exp. limits

* C.A. Baker et al., Phys. Rev. Lett. 97, 131801 (2006) * The ACME collaboration Science 343, p. 269-272 (2014)



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EDM EXPERIMENTS

Neutral systems

1951 Purcell and Ramsey:

parity violation in neutron scattering from nuclei

heavy atoms, molecules and neutrons

<u>Method</u>: apply electric field \vec{E} measure energy shift $\vec{d} \cdot \vec{E}$

EDM EXPERIMENTS

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Neutral systems

Charged particles

New method:

STORAGE RING as <u>trap</u> for POLARIZED charged particles beams

Storage ring

All rings have inward electric field in particle frame

$$\frac{d\vec{s}}{dt} = \vec{d} \times \vec{E}$$

EDM signal = spin precession in the vertical plane

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 $G = \frac{g-2}{2}$

EDM signal = spin precession in the vertical plane

TO DO: FREEZE the spin along the velocity

From the Thomas-BMT equation:

$$\vec{\omega_a} = \vec{\omega_s} - \vec{\omega_r} = \frac{-q}{m} \left\{ G \vec{B} + \left[G - \left(\frac{m}{p} \right)^2 \right] \frac{\vec{\beta} \times \vec{E}}{c} \right\}$$

Anomalous magnetic moment

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 $G = \frac{g-2}{2}$ Anomalous magnetic moment

Find E (radial) and B fields combination!

DEUTERON CASE

• Minimal detectable precession $\theta \approx 10^{-6} rad$

• Assuming $d \approx 10^{-29} e \cdot cm$

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• Possible ring: $E_{lab} = 17 MV/m$ $B_{lab} = 0,42 T$ p = 1,5 GeV/c

$$\theta_{EDM} = \frac{d}{\hbar} (E_{lab} + v B_{lab}) \sim (10^{-9} rad/s) t$$

$$\theta_{EDM} \sim \frac{10^{-15} rad}{turn}$$

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 Is this an issue? YES!

Horizontal polarization vanishes!

SCT STUDIES @ COSY

AIM

Demonstrate sextupole fields can counteract the spread of spin tunes

associated with **emittance** and **(Δp/p)**² of a **deuteron** beam. Second order effects!

In combination with beam preparation based on

- eCooling to shrink transverse and longitudinal beam size
- Bunching to remove first order Δp/p contribution

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Beam preparation

Polarized deuterons at p=0,97 GeV/c Bunched beam Horizontal polarization with RF solenoid

Experimental setup

Beam extraction

beam

continuous extraction with **white noise** applied to electric field plates (vertical direction, y axis)

17 mm

Carbon target

Elastic scattering events

(high spin sensitivity)

Polarimeter: EDDA

*Z. Bagdasarian et al., Phys. Rev. ST Accel. Beams 17, 052803 (2014)

- Timing \rightarrow Count turn number **n** (bunched beam)
- Compute total spin precession angle

 $\theta_s = 2 \pi G \gamma n$

- Bin by phase $\boldsymbol{\phi}$ the spin precession angle circle
- Compute asymmetry in each bin

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 $\epsilon_h = \frac{U - D}{U + D}$

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projection of the polarization.

Sextupole effect

Decoherence sources

Spin Tune spread: $\Delta v_s = G \Delta \gamma$

Betatron oscillations increase particle **path length**

$$\frac{\Delta L}{L} \propto \frac{\theta_x^2 + \theta_y^2}{4}$$

Bunching freezes the revolution frequency $\Delta \frac{y}{y} \propto \frac{\Delta L}{L}$ x 2 θ_x 4 Δv_s ldeal particle θ_x Δv_s z

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Sextupole effects

SCT dependence on sextupoles

$$\frac{1}{SCT} = |A + a_1 S + a_2 L + a_3 G| \theta_x^2$$
$$+ |B + b_1 S + b_2 L + b_3 G| \theta_y^2$$
$$+ |C + c_1 S + c_2 L + c_3 G| \left(\frac{\Delta p}{p}\right)^2$$
$$Drivers:$$

- beam widths

- 2nd order mom.

spread

Sextupole fields MXS, MXL, MXG

Sextupole effects

Sextupole effect

SCT dependence on sextupoles

 $\frac{1}{SCT} = |A + a_1 S| \theta_x^2$

- Flip 1/SCT sign above zero crossing
- Different slopes
- **SCT** does not go to infinity. Point near zero may be above or below the line due to other contributions.
- The same zero crossing, independent of horizontal width

2012 DATA

Chromaticity = ZERO

<u>Cromaticity</u> ξ

$$\frac{\Delta Q_{x,y}}{Q_{x,y}} = \xi_{x,y} \frac{\Delta p}{p}$$

A uniq MXG,

Q = tune, number of betatron socillation per turn

ue combination of
MXS and MXL
$$\xi = 0$$
 $\Rightarrow \frac{\Delta L}{L} = 0 \Rightarrow \frac{\Delta \gamma}{\gamma} = 0$ $\Rightarrow \frac{\Delta v_s}{v_s} = 0$

20

Chromaticity in MXG x MXS plane. MXL = -2.0 %.

Note the overlap of the two dotted lines that represent the places where the chromaticities vanish.

Picture from Ed. J. Stephenson

Chromaticity = ZERO

Make scans in 2D MXS x MXG space with MXL = -1.45%

Horizontal heating (large X emittance)
 Cool, then bunch (large synchrotron orbits)

Both transverse (X) and longitudinal spreads of the beam produce decoherence; both are canceled at places of zero chromaticity. Errors less than the size of the symbols.

The longest polarization lifetimes are found near the middle of this range.

lines of zero chromaticity (X or Y) in this plane – errors ~ 1 %

Scales are in percent of power supply full range.

Y Picture from Ed. J. Stephenson

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MXG

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A non-zero EDM within the actual experimental limits would be a clear probe of new physics

Fundamental requirement for the EDM experiment on charged particles is 1000 s SCT

Conclusions

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It has been demonstrated that the **lifetime of a horizontally polarized deuteron beam** may be substantially extended (**up to ~ 1000 s**) through a combination of:

- beam bunching on the first harmonic,
- electron cooling
- combination of SEXTUPOLE families
- both X and Y chromaticities are zero.

This meets the requirement for a storage ring to search for an EDM.

Conclusions

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Fundamental requirement for the EDM experiment on charged particles is 1000 s SCT

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FUTURE

- systematic error studies
- spin tracking simulations
- polarimeter
- development of a new ring

THANKS FOR YOUR ATTENTION

JEDI collaboration @ Juelich

CP violating sources

STANDARD MODEL

- Weak interaction: complex phase $\boldsymbol{\delta}$ in CKM quark mixing matrix
- Strong interaction: $\boldsymbol{\theta}_{ocn}$

$$|d_{n}| = |d_{p}| \simeq 4,5 \cdot 10^{-15} \theta_{QCD} \longrightarrow |d_{n}^{exp}| \le 10^{-26} e \cdot cm \Rightarrow \theta_{QCD} \le 10^{-11}$$

Axion search

SUSY

- quark-EDM $\Delta = d_{down} d_{up}/4$
- Chromo-EDM: EDM generated by a loop with SS-particle

$$\Delta^+ = d_{up}^c + d_{down}^c \qquad \Delta^- = d_{up}^c - d_{down}^c$$

$$d_{n} = 1,4 \Delta + 0,83 \Delta^{+} - 0,27 \Delta^{-}$$
$$d_{p} = 1,4 \Delta + 0,83 \Delta^{+} + 0,27 \Delta^{-}$$
$$d_{d} = d_{up} + d_{down} - 0,2 \Delta^{+} - 6 \Delta^{-}$$

If a **non-zero deuteron EDM** is measured, it would have a special sensitivity to the chromo-EDM due to the large coefficient of Δ^- .

The EDM measurement of **several particles** is needed to determine the CP violating sources scenario.