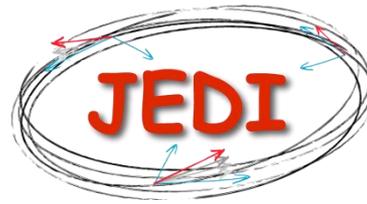




**RWTHAACHEN**  
UNIVERSITY



# Investigation of Beam and Spin Dynamics for EDM Measurements at COSY

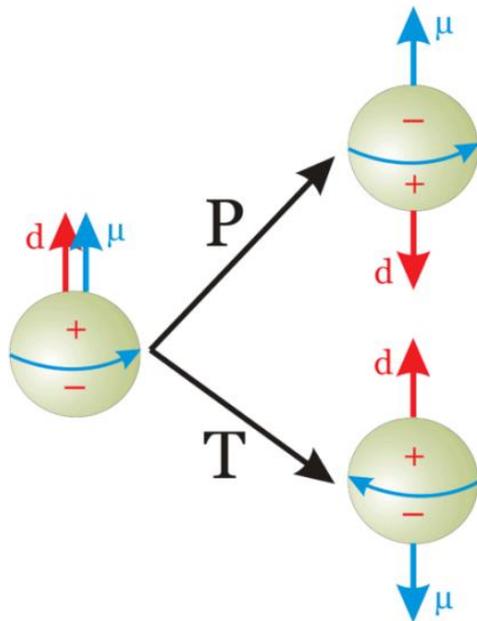
2014-09-04 | Marcel Rosenthal on behalf of the JEDI Collaboration

## ➤ **Part 1: Introduction**

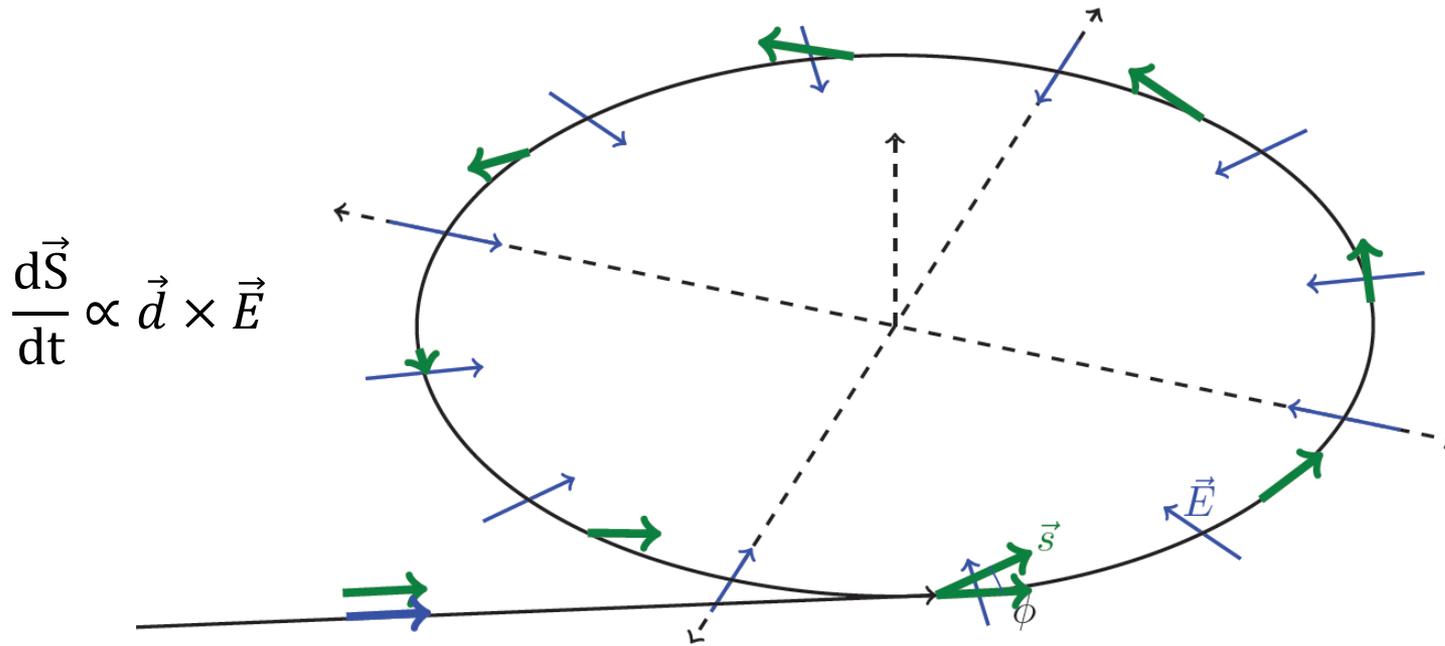
- What are Electric Dipole Moments?
- General idea for EDM measurements in storage rings
- The Cooler Synchrotron COSY, Jülich
- Thomas-BMT-equation

## ➤ **Part 2: Simulations**

- Simulation framework
- Measurement principle at COSY, Jülich
- Spin Coherence Time studies
- False signal due to magnet imperfections



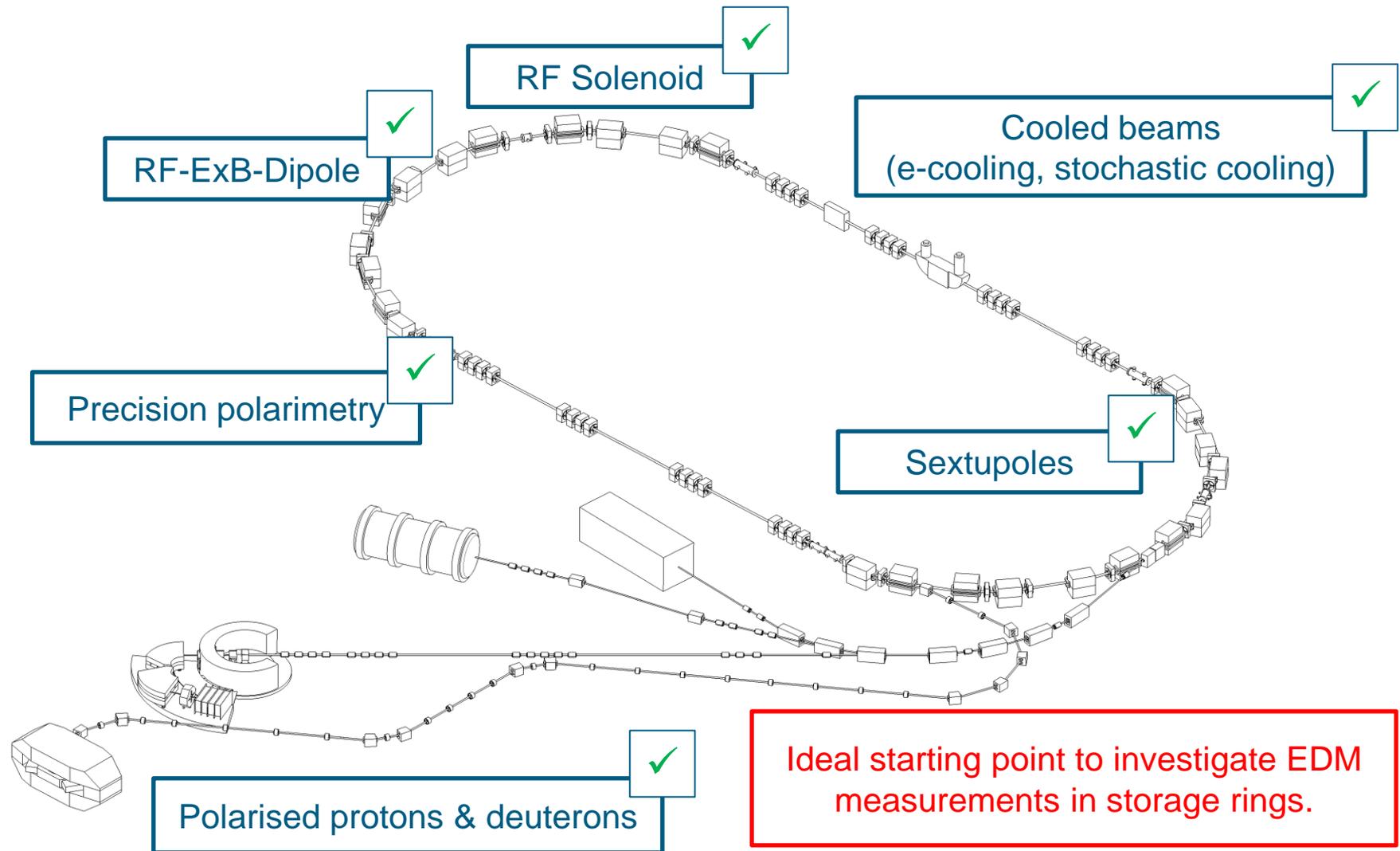
- Electric Dipole Moments:
  - Charge separation
  - Fundamental property
- Permanent EDMs are P- and T-violating
  - CPT-Theorem: CP-Violation
- Known CP-Violation not sufficient to explain Matter-Antimatter-Asymmetry in universe
- Search for new sources of CP-Violation ( $\Theta$ -term, BSM) by measuring Electric Dipole Moments of charged hadrons in storage rings



➤ General idea:

- Inject polarised particles with spin pointing towards momentum direction
- „Frozen Spin“-Technique: without EDM spin stays aligned to momentum
- EDM couples to electric bending fields
- Slow buildup of EDM related vertical polarisation

# The Cooler Synchrotron COSY



- Equation of spin motion for relativistic particles in electromagnetic fields:

$$\frac{d\vec{S}}{dt} = \vec{S} \times \vec{\Omega}_{\text{MDM}} + \vec{S} \times \vec{\Omega}_{\text{EDM}}$$

$$\vec{\Omega}_{\text{MDM}} = \frac{e}{\gamma m} \left[ G\gamma \vec{B} - \left( G - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{E} \times \vec{\beta}}{c} - \frac{G\gamma^2}{\gamma + 1} \vec{\beta}(\vec{\beta} \cdot \vec{B}) \right]$$

$$\vec{\Omega}_{\text{EDM}} = \frac{e}{m} \frac{\eta}{2} \left[ \frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} - \frac{\gamma}{\gamma + 1} \vec{\beta} \left( \vec{\beta} \cdot \frac{\vec{E}}{c} \right) \right]$$

$$\vec{\mu} = 2(G + 1) \cdot \frac{e}{2m} \vec{S}$$

	G
Proton	1.792847357
Deuteron	-0.142561769

$$\vec{d} = \frac{\eta}{2} \cdot \frac{e}{2mc} \vec{S}$$

d	η
$10^{-24}$ e cm	$\sim 10^{-9}$
$10^{-29}$ e cm	$\sim 10^{-14}$

- Equation of spin motion for relativistic particles in electromagnetic fields:

$$\frac{d\vec{S}}{dt} = \vec{S} \times \vec{\Omega}_{\text{MDM}} + \vec{S} \times \vec{\Omega}_{\text{EDM}}$$

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$$|\vec{\Omega}_{\text{EDM}}| = \frac{e}{m} \frac{\eta}{2} \left[ \frac{|\vec{E}|}{c} + \vec{\beta} \times \vec{B} - \frac{\gamma}{\gamma + 1} \vec{\beta} \left( \vec{\beta} \cdot \frac{\vec{E}}{c} \right) \right] \ll |\vec{\Omega}_{\text{MDM}}|$$

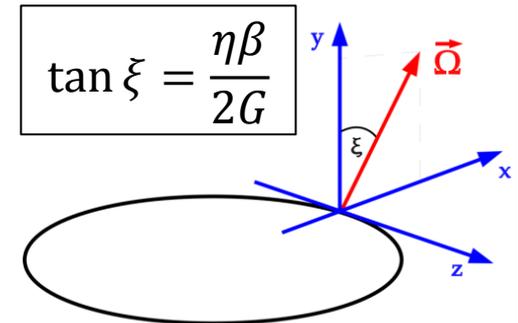
- Cooler Synchrotron Jülich is conventional pure magnetic ring:
  - Spin precesses around vertical guiding field.
  - Number of spin precessions per revolution (with respect to the momentum vector) is given by the spin tune  $\nu_s = G\gamma$
  - „Frozen Spin“-Technique requires  $\vec{\Omega}_{\text{MDM}} = 0$ 
    - Not applicable

- Static storage ring:
  - Tilt of  $\vec{\Omega}$  in main dipoles due to EDM contribution

$$\frac{d\vec{S}}{dt} = \vec{S} \times \vec{\Omega}_{\text{MDM}} + \vec{S} \times \vec{\Omega}_{\text{EDM}}$$

$$\vec{\Omega}_{\text{MDM}} = \frac{e}{\gamma m} \left[ G\gamma \vec{B} - \left( G - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{E} \times \vec{\beta}}{c} - \frac{G\gamma^2}{\gamma + 1} \vec{\beta} (\vec{\beta} \cdot \vec{B}) \right]$$

$$\vec{\Omega}_{\text{EDM}} = \frac{e \eta}{m 2} \left[ \frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} - \frac{\gamma}{\gamma + 1} \vec{\beta} \left( \vec{\beta} \cdot \frac{\vec{E}}{c} \right) \right]$$

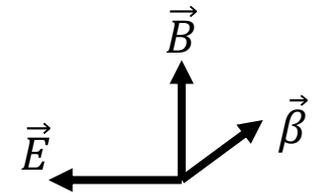


- $\vec{\Omega}_{\text{MDM}}$  vertical,  $\vec{\Omega}_{\text{EDM}}$  radial
- Small tilt of spin precession axis.
- Vertical oscillation is too small to measure, if  $|\vec{\Omega}_{\text{EDM}}| \ll |\vec{\Omega}_{\text{MDM}}|$

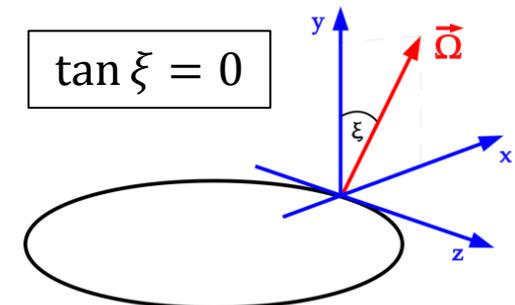
- Idea:
  - Radiofrequent field oscillating with spin precession frequency
  - Pure electric field: coherent betatron oscillations
  - Minimization using Wien filter configuration
    - RF-E×B-Dipole

$$\vec{\Omega}_{\text{MDM}} = \frac{e}{\gamma m} \left[ G\gamma \vec{B} - \left( G - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{E} \times \vec{\beta}}{c} - \frac{G\gamma^2}{\gamma + 1} \vec{\beta}(\vec{\beta} \cdot \vec{B}) \right]$$

$$\vec{\Omega}_{\text{EDM}} = \frac{e \eta}{m 2} \left[ \frac{\vec{E}}{c} + \vec{\beta} \times \vec{B} - \frac{\gamma}{\gamma + 1} \vec{\beta} \left( \vec{\beta} \cdot \frac{\vec{E}}{c} \right) \right] = \vec{0}$$



- Device is EDM transparent, no tilt  $\xi$  like in guiding dipoles
  - Slowly oscillating signal



## ➤ Part 1: Introduction

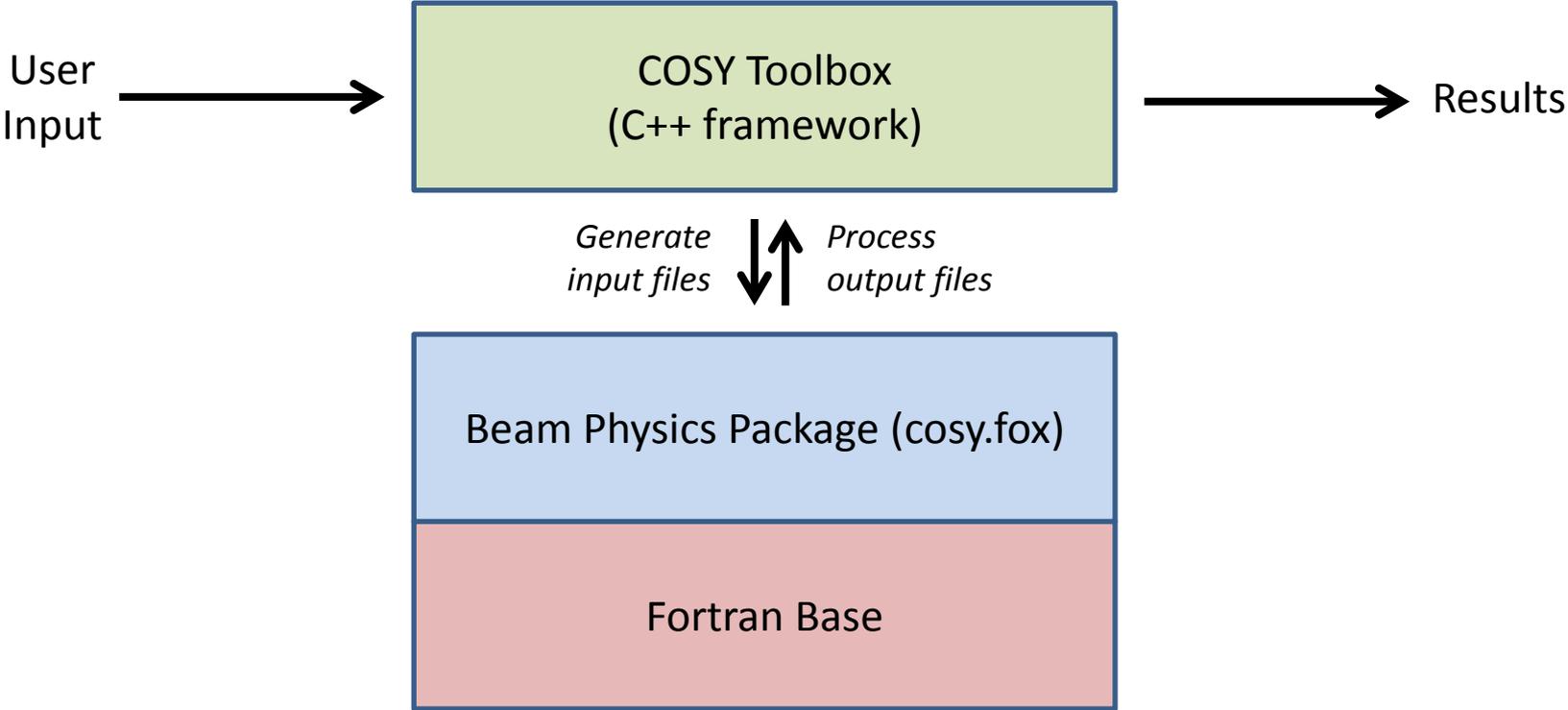
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## ➤ Part 2: Simulations

- Simulation framework
- Measurement principle at COSY, Jülich
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# Simulation framework

- New package on top of COSY Infinity



# COSY Toolbox

COSY INFINITY

COSY Lattice

Armadillo

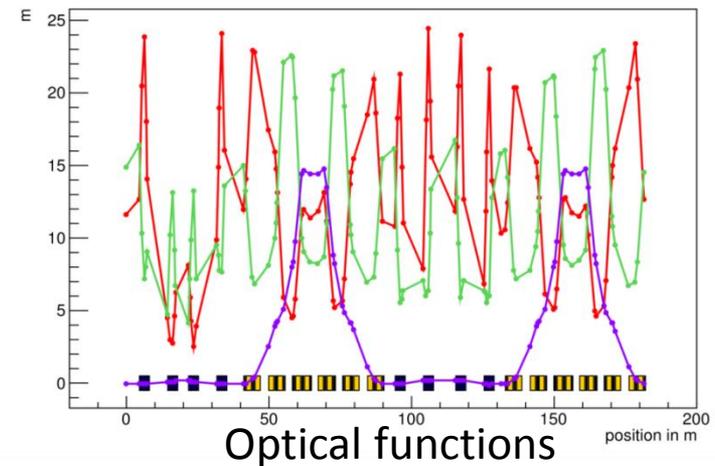
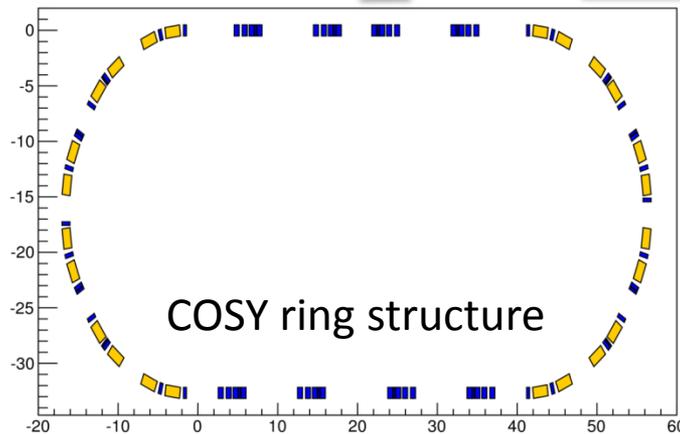
COSY Toolbox

ROOT

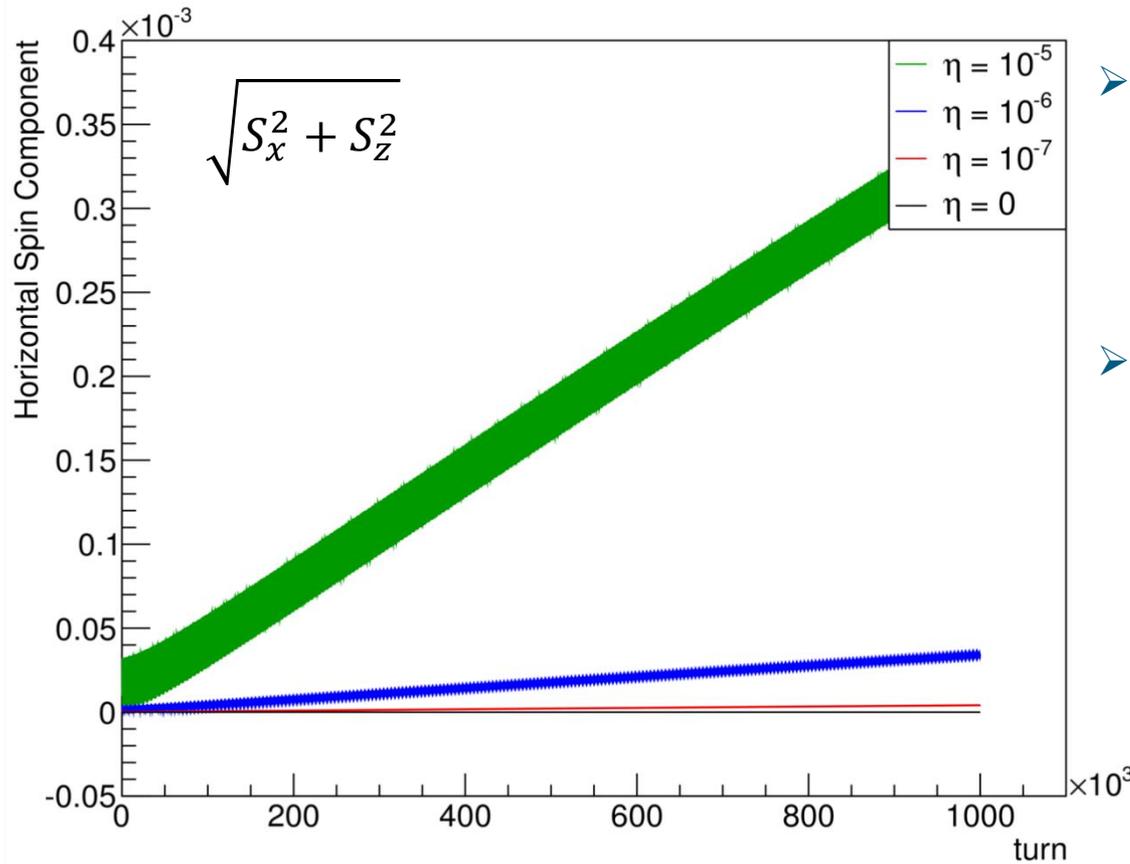
- Calculator:
  - Optical functions
  - Closed orbit
  - Spin tune ...
- Tracker:
  - Static maps
  - RF maps (\*)

- linear algebra operations
- SVD, pseudo-inverse

- plotting
- storage (ROOT files/trees)



- RF-Simulation for ideal ring shows effect
  - Polarized deuteron,  $p = 970 \text{ MeV}/c$ , RF-E×B-Dipole strength: 0.3 mT



- Initial vertical polarization will precess into horizontal plane, if  $\eta \neq 0$

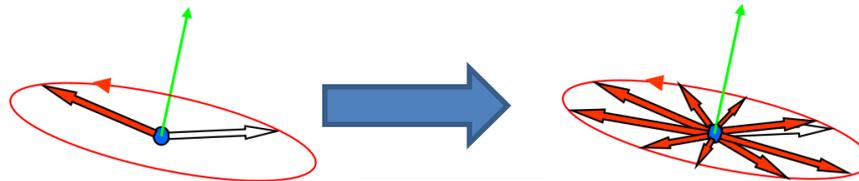
- Oscillation of horizontal spin component:

$$\sin(\theta \cdot n) \approx \theta \cdot n$$

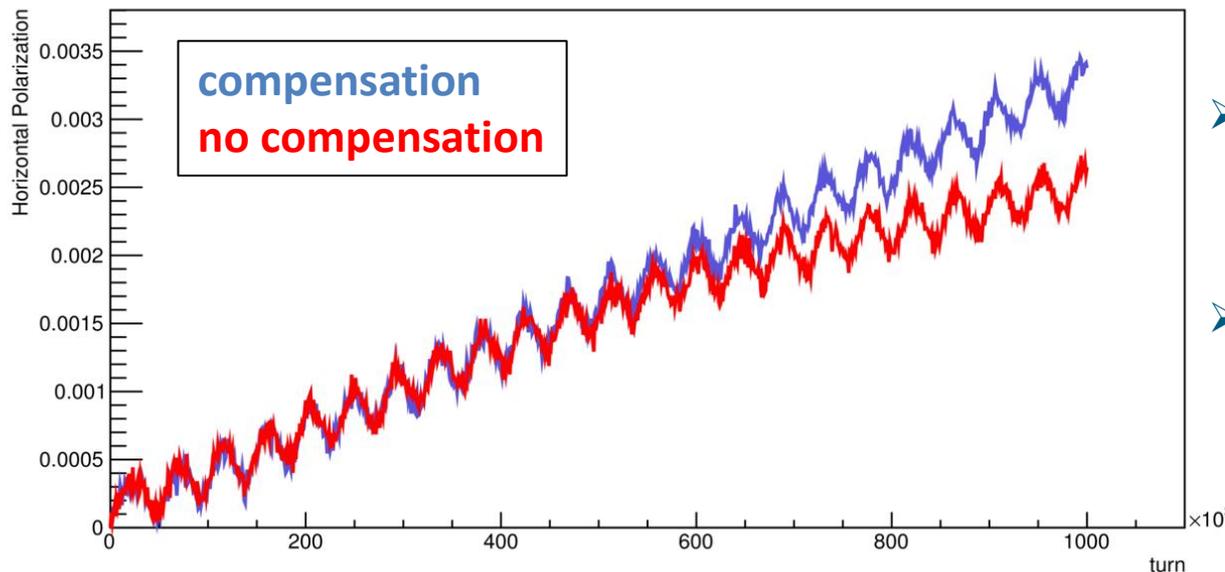
$$\theta \approx \frac{\psi \xi}{2}$$

$\psi$ : max. spin rotation in RF field  
 $\xi$ : EDM tilt

- Spin precession in ideal magnetic ring around vertical axis:
  - Spin tune:  $\nu_s = G\gamma$
  - Energy deviations lead to different precession speed

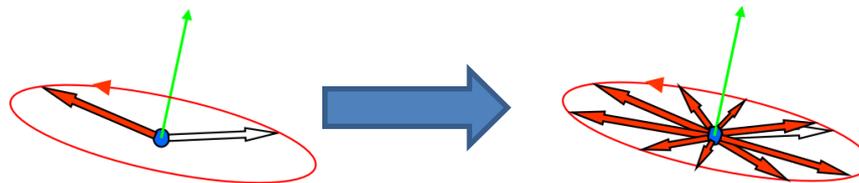


**Horizontal polarization vanishes!**



- Buildup limited by Spin Coherence Time
- Decoherence needs to be minimized

- Spin precession in ideal magnetic ring around vertical axis:
  - Spin tune:  $\nu_s = G\gamma$
  - Energy deviations lead to different precession speed



**Horizontal polarization vanishes!**

- Consider relative change of revolution time of single particle:

$$\text{➤ } \frac{\Delta T}{T_0} = \frac{\Delta L}{L_0} - \frac{\Delta\beta}{\beta_0} - \frac{\Delta L}{L_0} \frac{\Delta\beta}{\beta_0} + \left(\frac{\Delta\beta}{\beta_0}\right)^2 \quad \text{with} \quad T_0 = \frac{L_0}{\beta_0 c}$$

- No coupling:

$$\text{➤ } \frac{\Delta L}{L_0} = \left(\frac{\Delta L}{L_0}\right)_x + \left(\frac{\Delta L}{L_0}\right)_y + \left(\frac{\Delta L}{L_0}\right) \frac{\Delta p}{p} \quad \text{➤ } \left(\frac{\Delta L}{L_0}\right) \frac{\Delta p}{p} = \alpha_0 \cdot \frac{\Delta p}{p} + \alpha_1 \cdot \left(\frac{\Delta p}{p}\right)^2$$

$$\left\langle \frac{\Delta T}{T_0} \right\rangle = \left( \alpha_0 - \frac{1}{\gamma_0^2} \right) \left\langle \frac{\Delta p}{p} \right\rangle + \left( \alpha_1 + \frac{3\beta_0^2}{2\gamma_0^2} - \frac{\alpha_0}{\gamma_0^2} + \frac{1}{\gamma_0^4} \right) \left\langle \left( \frac{\Delta p}{p} \right)^2 \right\rangle + \left\langle \left( \frac{\Delta L}{L_0} \right)_x \right\rangle + \left\langle \left( \frac{\Delta L}{L_0} \right)_y \right\rangle = 0$$

↑  
bunched

➤ Canceling energy deviations ( $v_s = G\gamma$ ):  $\left\langle \frac{\Delta \gamma}{\gamma_0} \right\rangle = 0$

➤  $\frac{\Delta p}{p} \approx \frac{1}{\beta_0^2} \frac{\Delta \gamma}{\gamma_0} - \frac{1}{2\beta_0^4 \gamma_0^2} \left( \frac{\Delta \gamma}{\gamma} \right)^2$

➤ Three conditions for  $\left\langle \frac{\Delta \gamma}{\gamma_0} \right\rangle = 0$ :

➤  $\left\langle \left( \frac{\Delta L}{L_0} \right)_u \right\rangle = -\frac{\pi}{L_0} \cdot \epsilon_u \cdot \xi_u = 0, \quad u \in \{x, y\}$

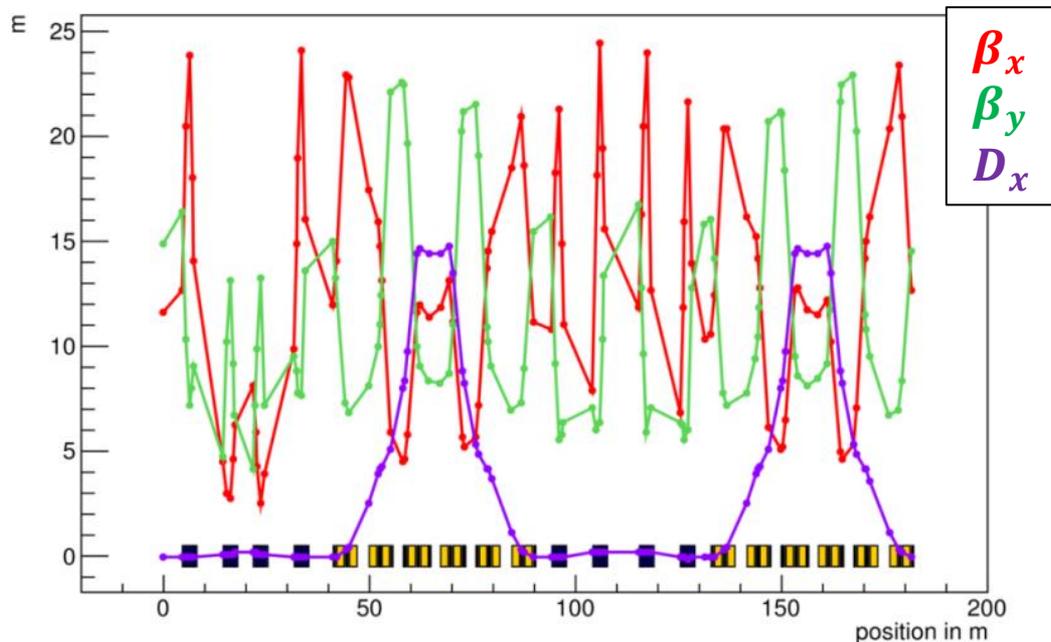
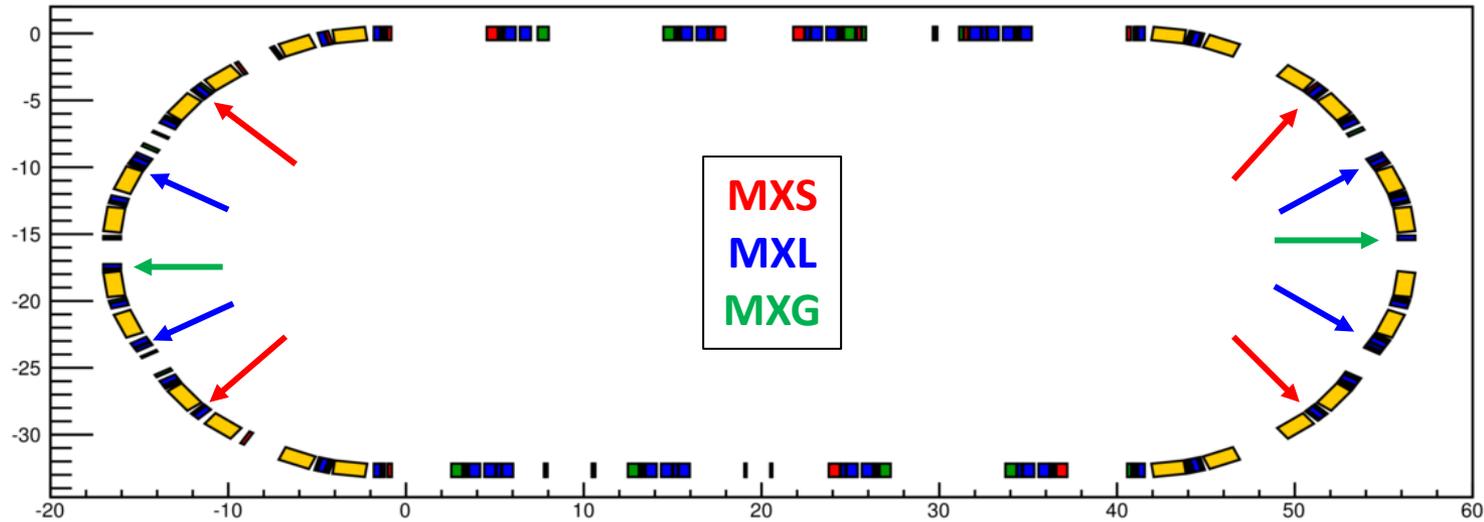
$$\epsilon_u = \frac{u_{max}^2}{\beta_u}$$

$$\xi_u = \frac{\Delta Q_u / Q_u}{\Delta p / p}$$

➤  $\Delta \equiv \left[ \alpha_1 + \frac{3}{2\gamma_0^2} \left( \beta_0^2 - \left( \alpha_0 - \frac{1}{\gamma_0^2} \right) \right) \right] = 0$

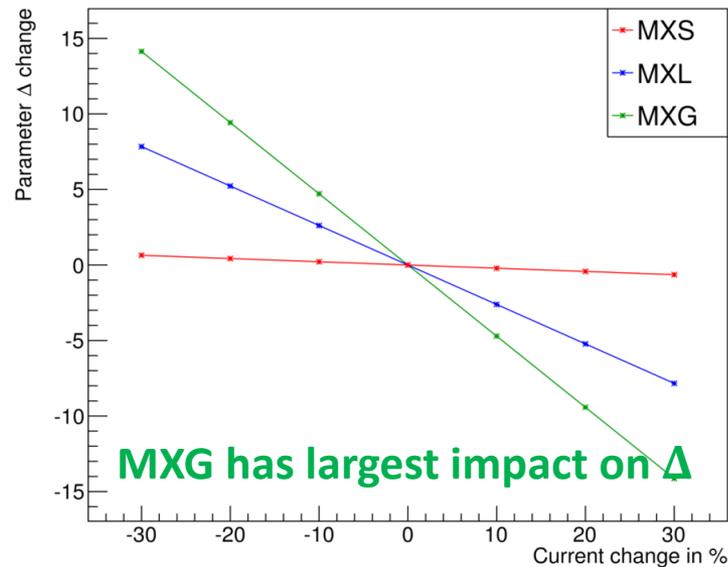
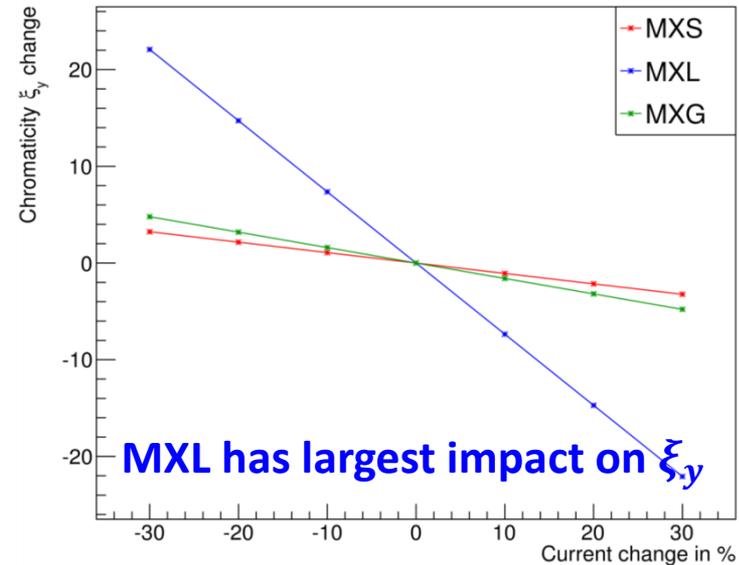
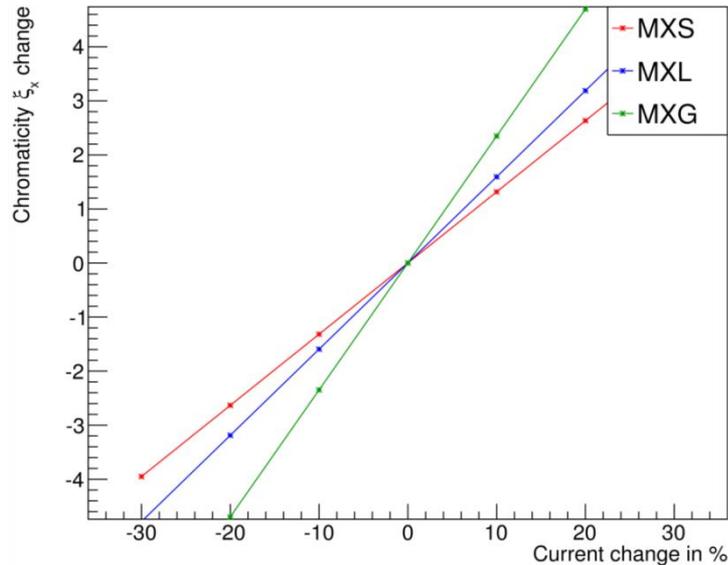
➤ Magnetic sextupoles are an effective tool to maintain these conditions.

# Sextupoles at COSY



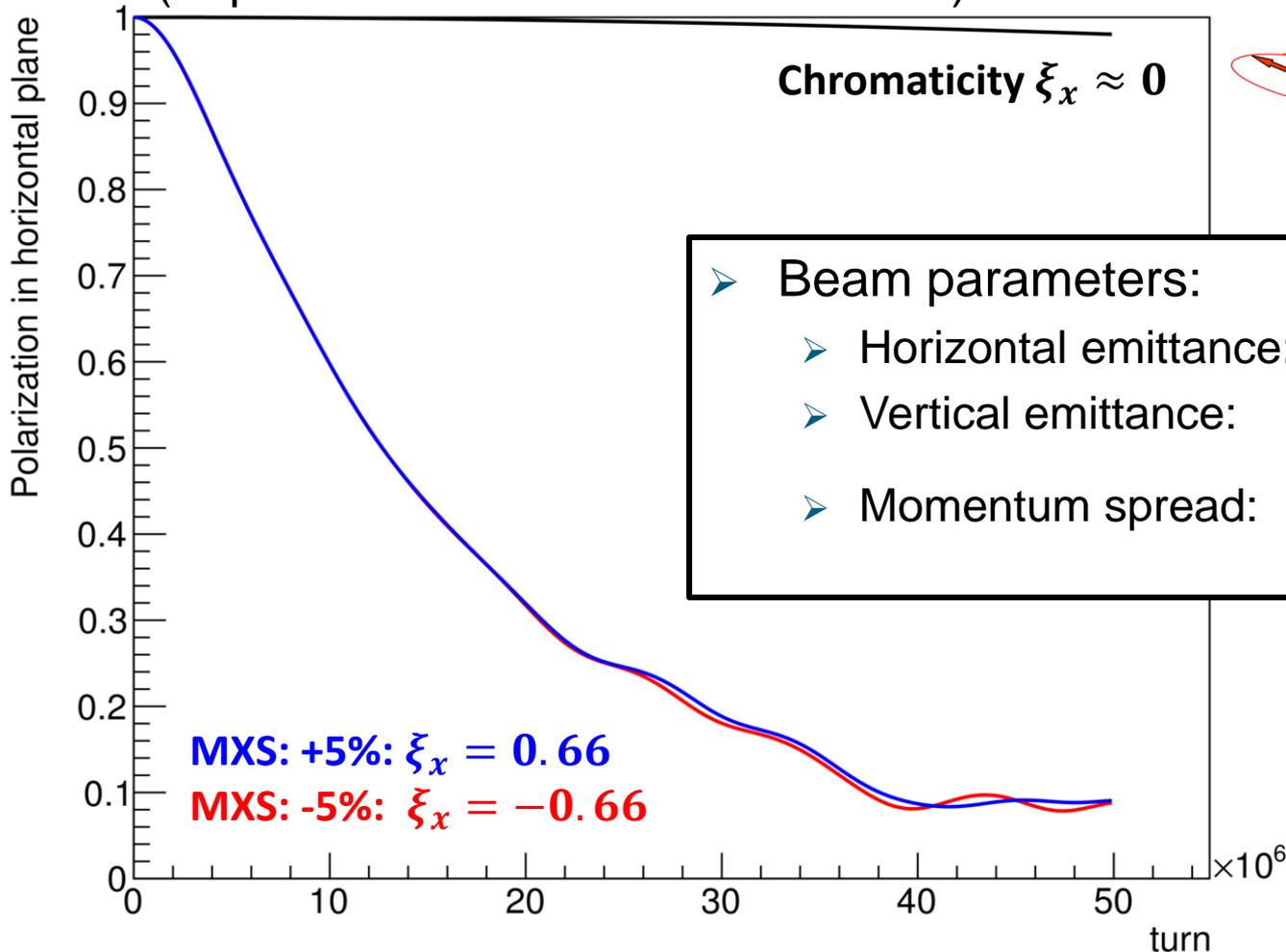
- Straights:
  - Minimized dispersion
- Arcs:
  - Large dispersion
  - 3 sextupole families

# Sextupoles at COSY

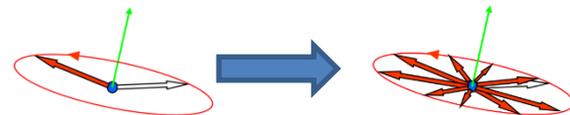


- Linear equation system to minimize  $\xi_x$ ,  $\xi_y$  and  $\Delta$  at the same time
- MXL dominates  $\xi_y$  change
- MXG dominates  $\Delta$  change

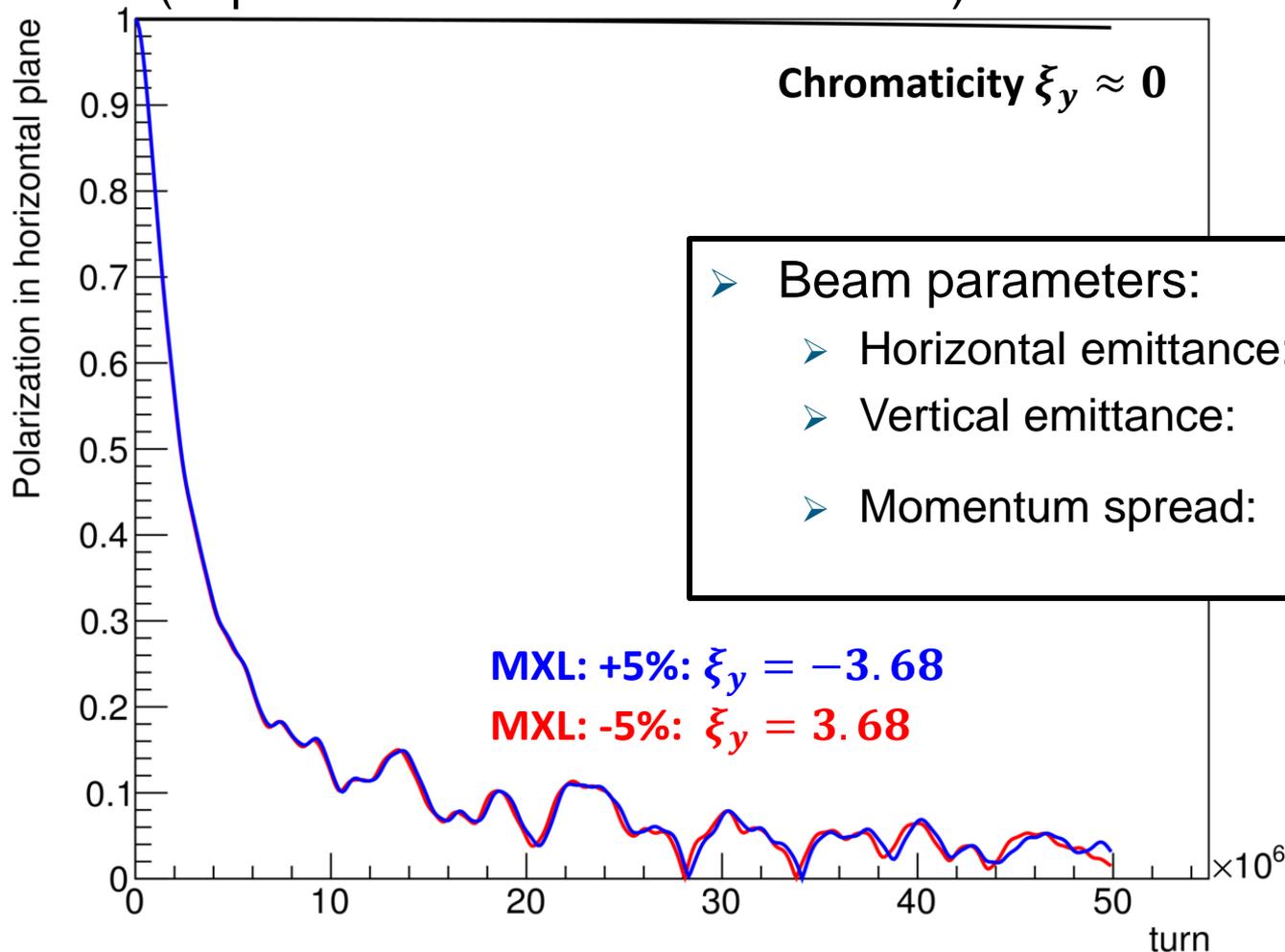
- Deuterons,  $p = 970 \text{ MeV}/c$ , initially radial polarized  
(→ precession around vertical axis)



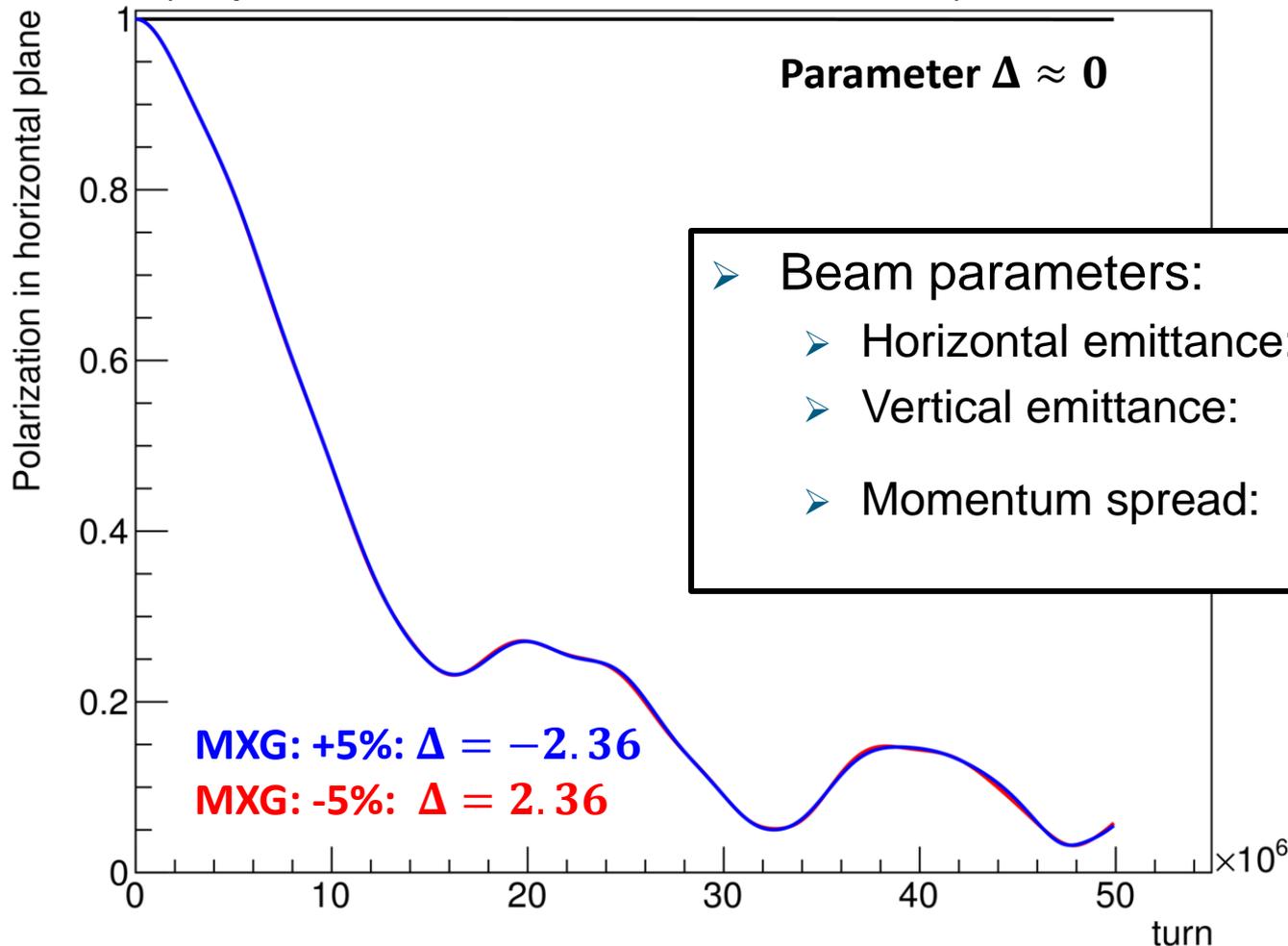
- **Beam parameters:**
  - Horizontal emittance:  $\epsilon_x = 5 \text{ mm mrad}$
  - Vertical emittance:  $\epsilon_y = 0 \text{ mm mrad}$
  - Momentum spread:  $\left(\frac{\Delta p}{p}\right)_{rms} = 0$



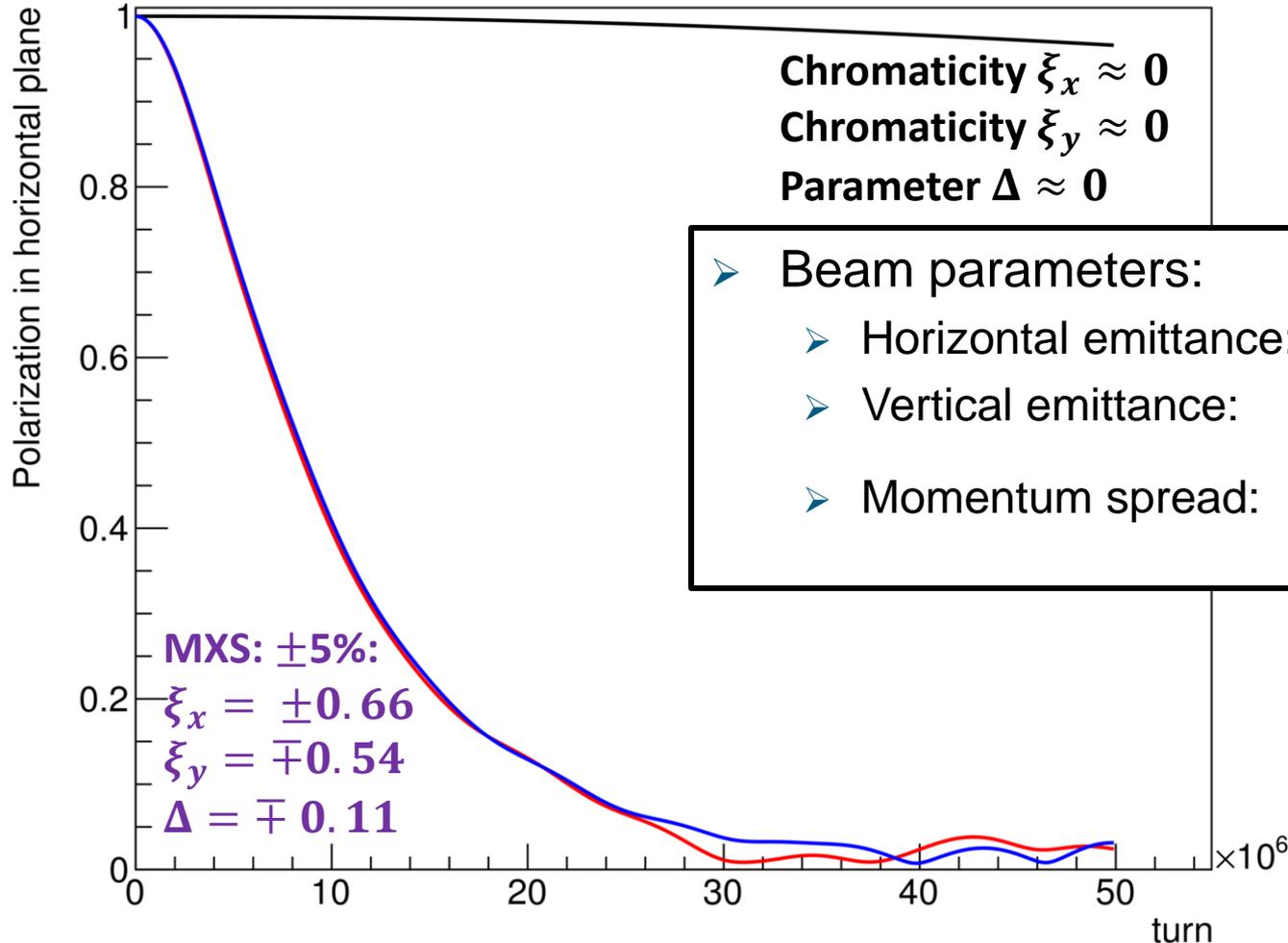
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- Deuterons,  $p = 970 \text{ MeV}/c$ , initially radial polarized  
(→ precession around vertical axis)



- Up to now only the ideal ring was considered.
- Misalignments and field errors of dipoles and quadrupoles introduce additional field components.

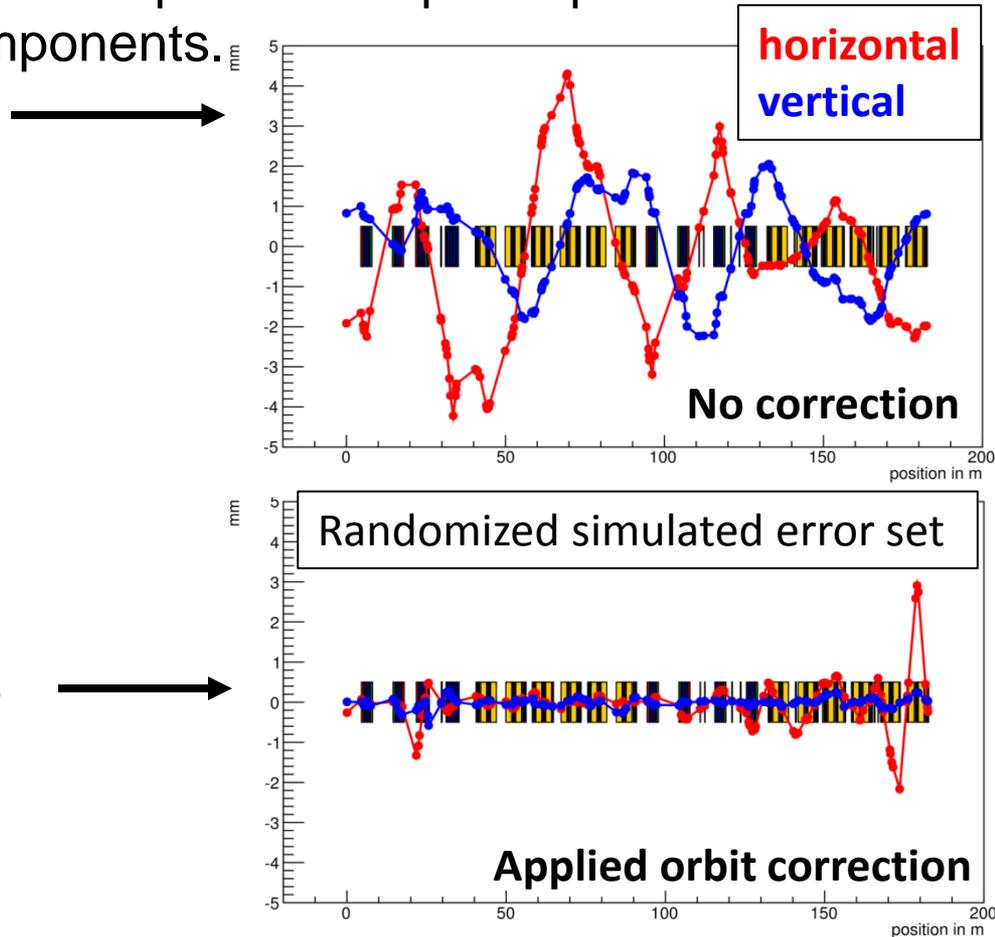
- Beam orbit displacement
- Spin rotations

- COSY main magnets:

- 24 dipoles
- 56 quadrupoles

- Orbit diagnosis & correction:

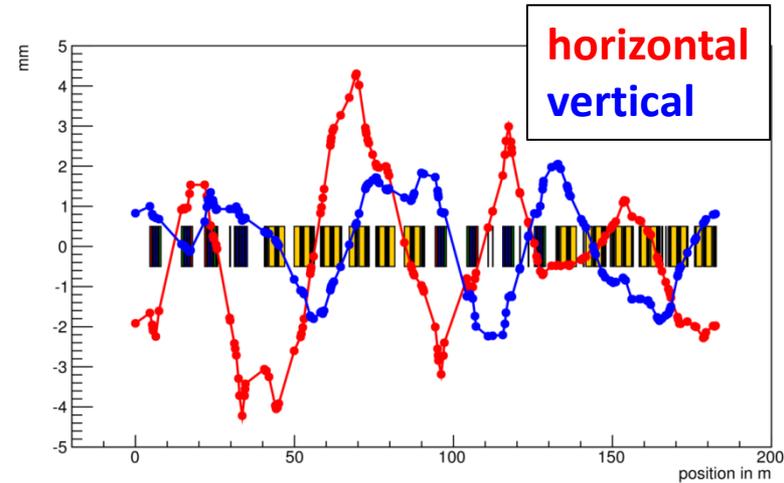
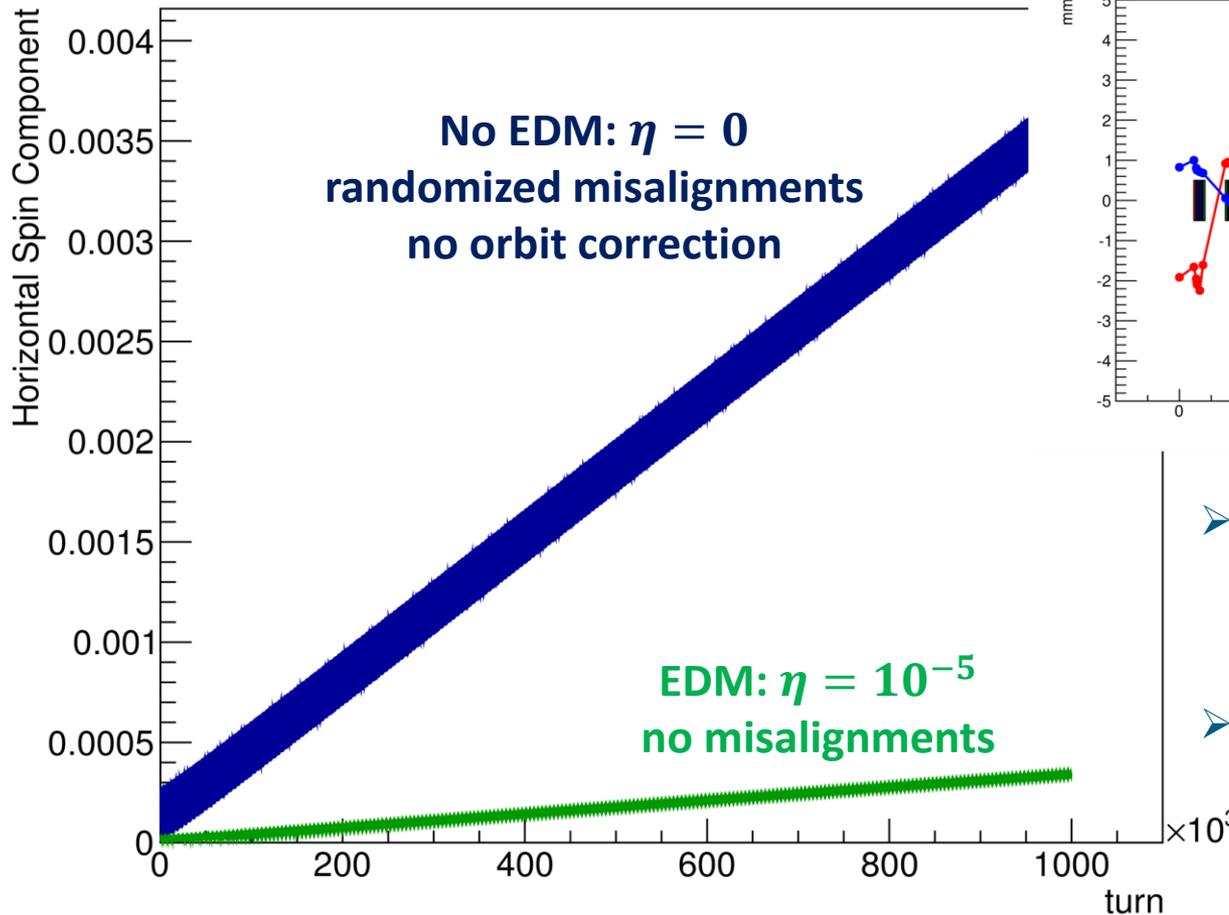
- ~60 beam position monitors
- ~40 correctors



# False “EDM signal“

- No EDM, but misalignments.  
What will happen to the polarization?

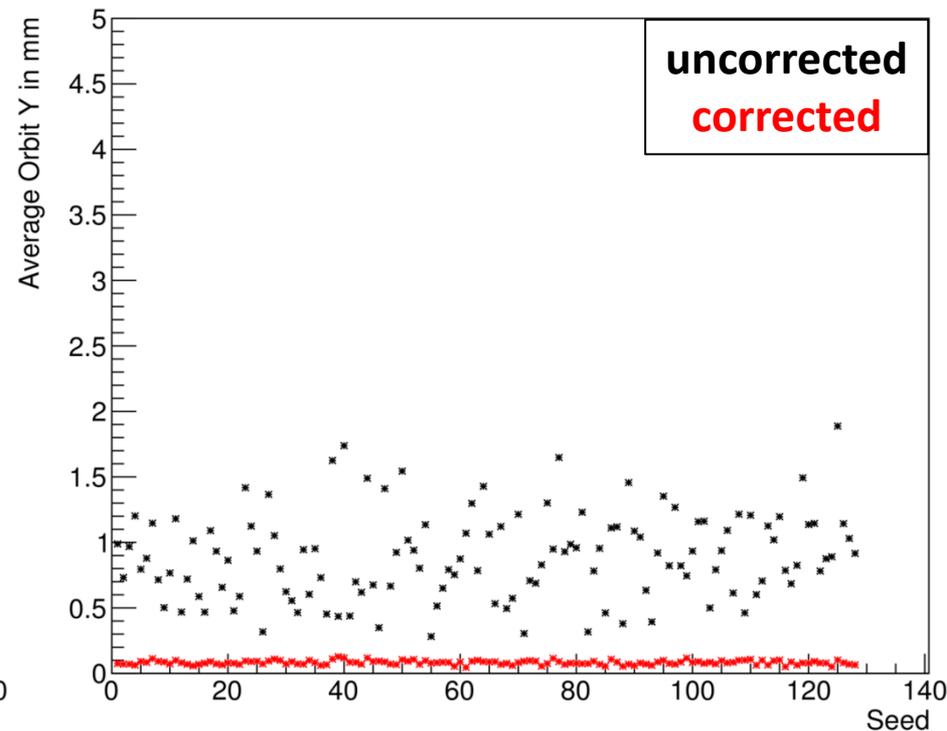
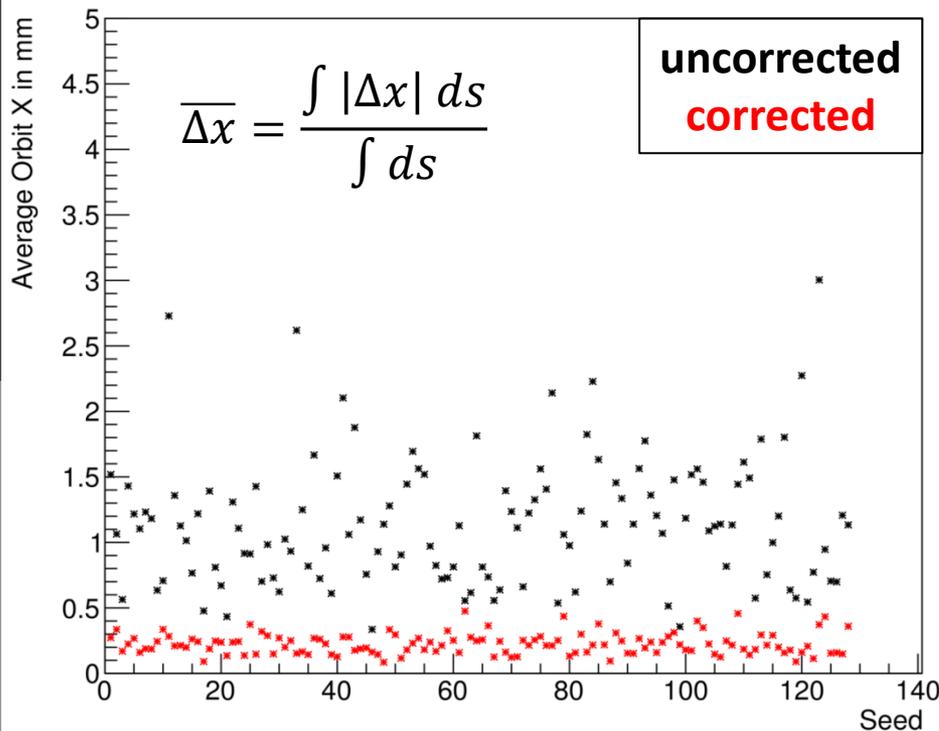
RF-E×B-Dipole strength: 0.3 mT



- Large contribution due to uncorrected misalignments
- Apply corrections to minimize

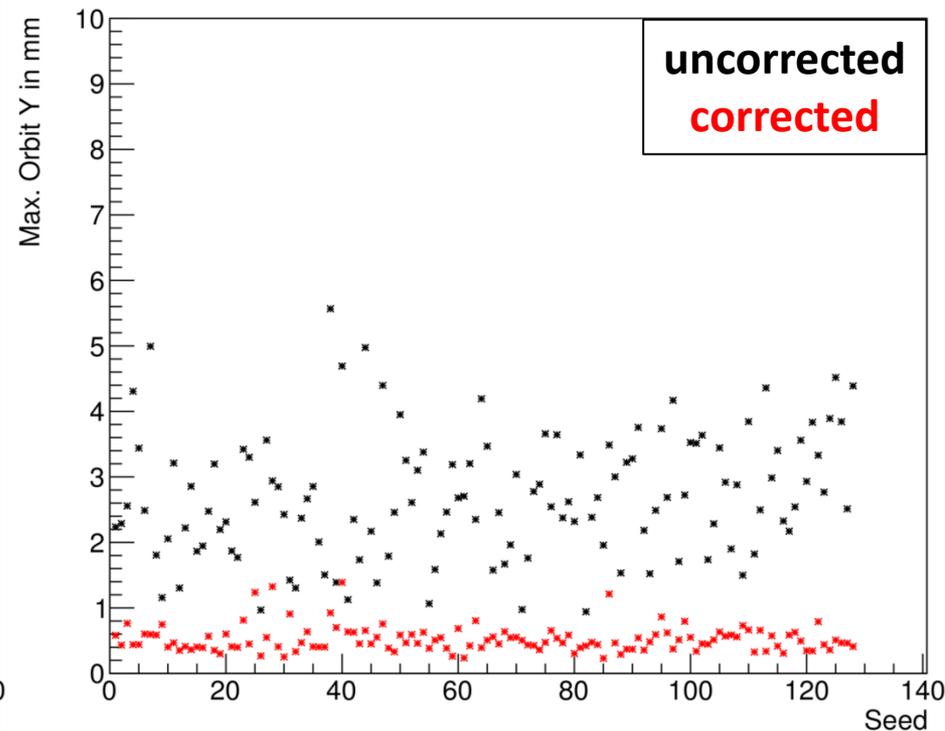
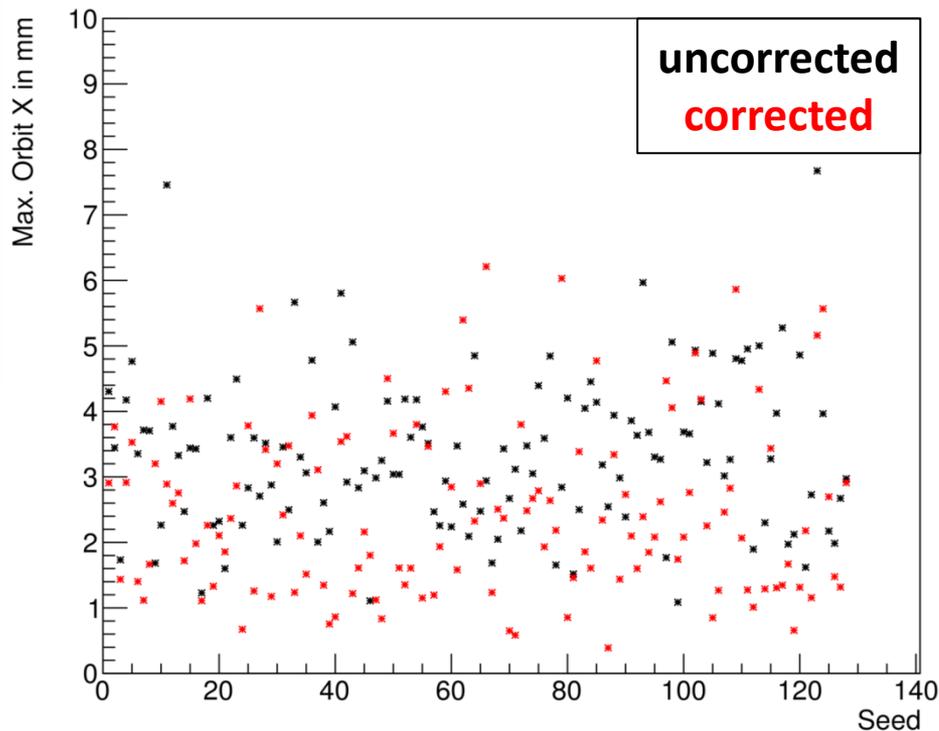
# Simulations of Misalignments

- Randomize different sets of misalignments and field errors.
- Shifts of elements in all three directions with  $\sigma = 10^{-4} m$
- Rotations of elements around all three axes with  $\sigma = 10^{-4} rad$
- Relative magnetic field error of main dipoles with  $\sigma = 10^{-4}$

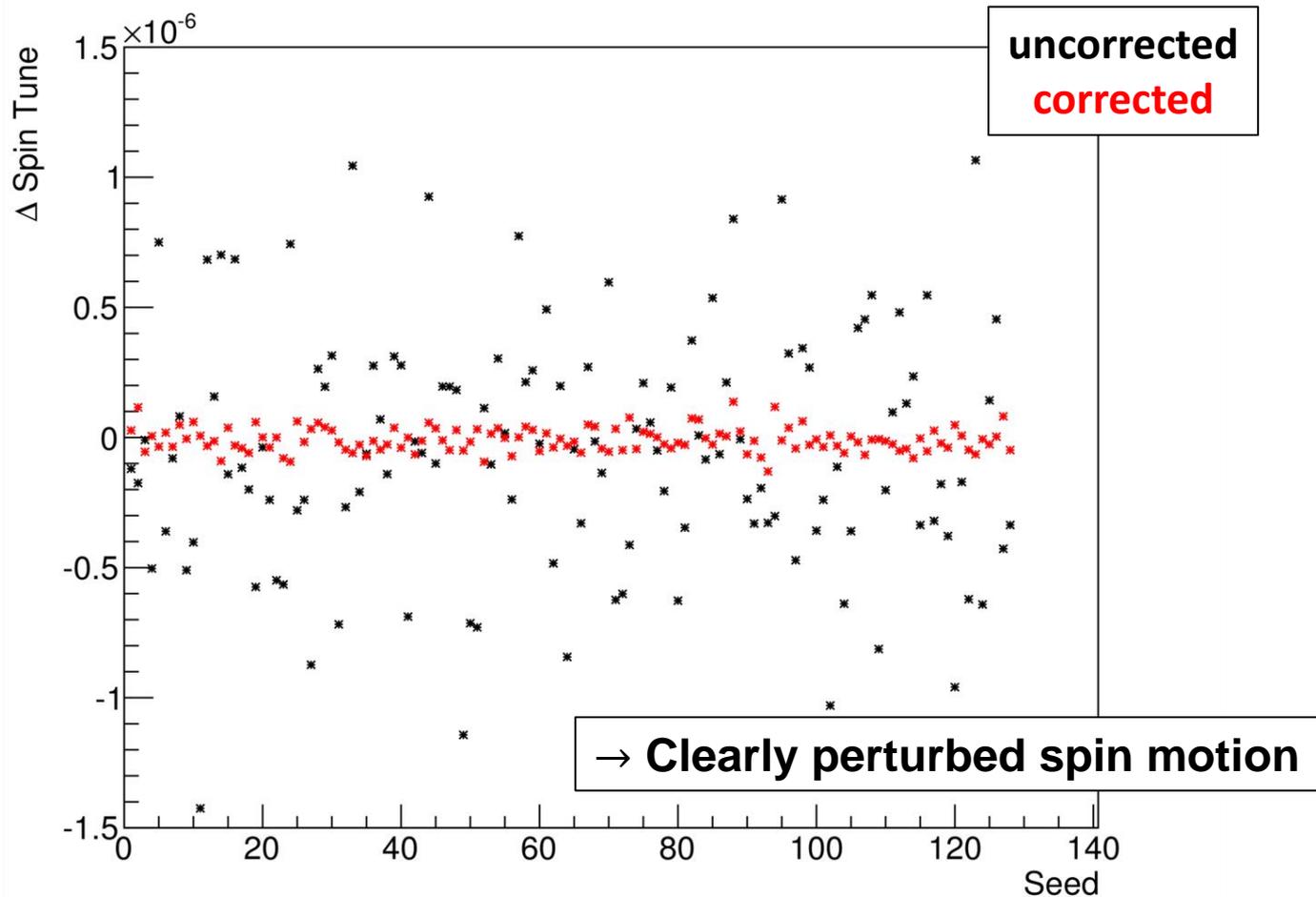


# Simulations of Misalignments

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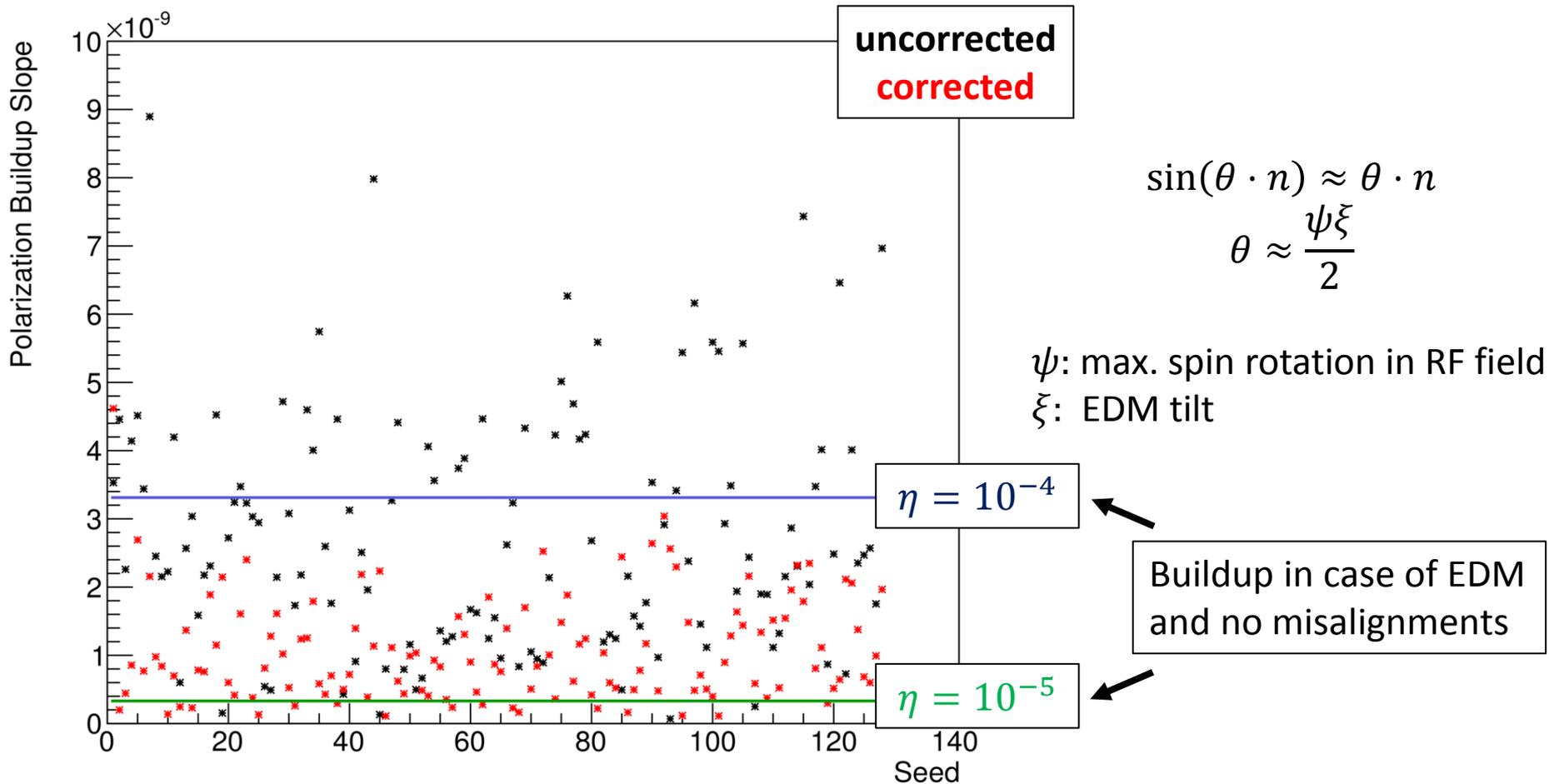


- Non commutative spin rotations change the spin tune  $\nu_s = G\gamma$
- Local orbit corrections partially compensate this effect



# Horizontal Polarization Buildup

- Polarization buildup caused by misalignments is in the order of an EDM:  $\eta \sim 10^{-4} \rightarrow d \sim 5.3 \cdot 10^{-19} \text{ e cm}$

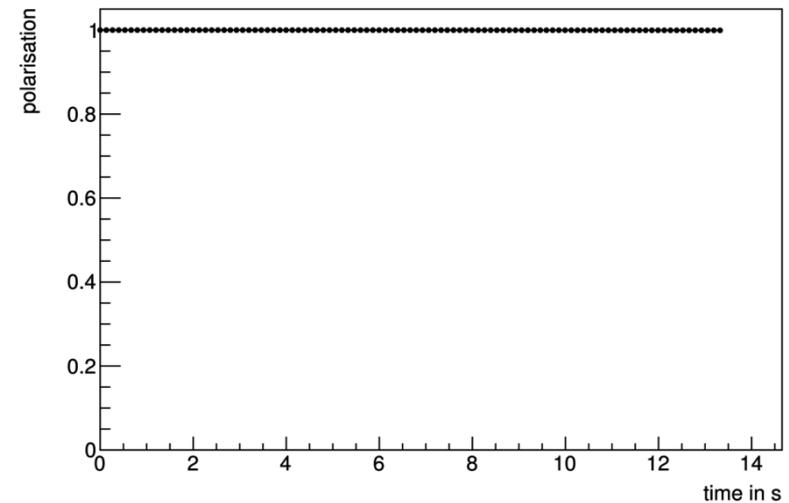
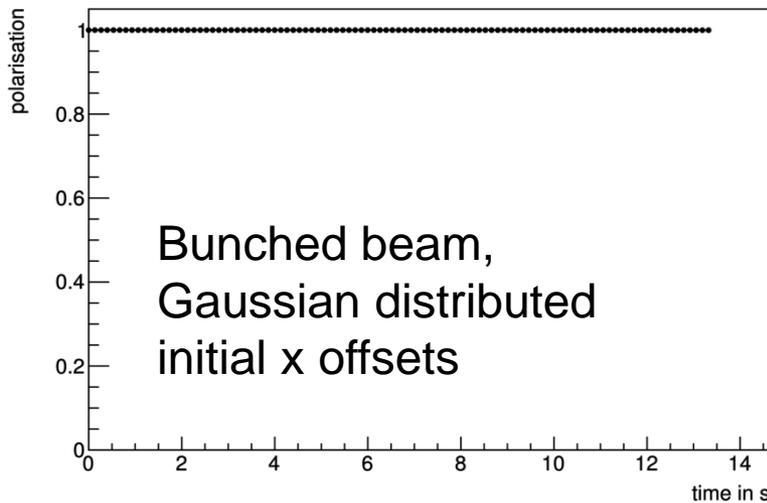


- EDM measurements in storage rings require a long polarization lifetime in vertical direction as well as in the horizontal plane.
- Chromaticities  $\xi_x$ ,  $\xi_y$  and  $\Delta \equiv \left[ \alpha_1 + \frac{3}{2\gamma_0^2} \left( \beta_0^2 - \left( \alpha_0 - \frac{1}{\gamma_0^2} \right) \right) \right]$  are crucial parameters for minimization of energy spread inside a bunched beam.
- Large set of measured horizontal depolarization for various sextupole configurations has to be understood using tracking simulations.
- Misalignments create false EDM signal in resonant buildup method. Method of local orbit corrections needs to be studied to minimize this signal, while not compensating the EDM effect itself.

# SPARES

- Several settings leading to vanishing path-lengthening:

#	MXS / $m^{-3}$	MXL / $m^{-3}$	MXG / $m^{-3}$
1	-0.002 (0)	-0.506	0.775
2	0.999 (1)	-0.522	0.497
3	2.001 (2)	-0.538	0.218
4	3.002 (3)	-0.554	-0.061
5	4.004 (4)	-0.570	-0.340



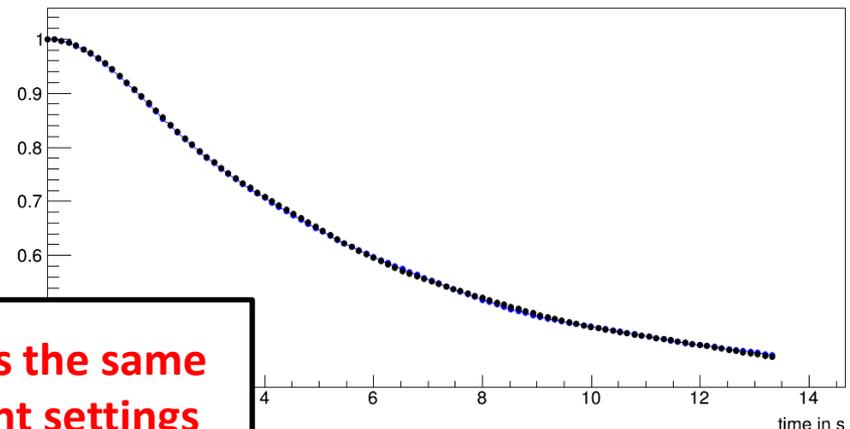
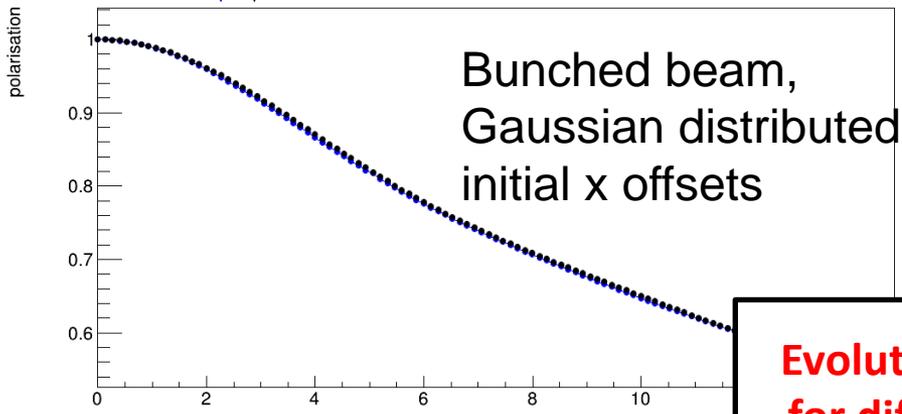
# Path-lengthening (horizontal)

- Deviations from path-lengthening cancellation:

#	MXS / m <sup>-3</sup>	MXL / m <sup>-3</sup>	MXG / m <sup>-3</sup>
1	-0.002 (0)	-0.506	0.775
2	0.999 (1)	-0.522	0.497
3	2.001 (2)	-0.538	0.218
4	3.002 (3)	-0.554	-0.061
5	4.004 (4)	-0.570	-0.340

MXS: +1

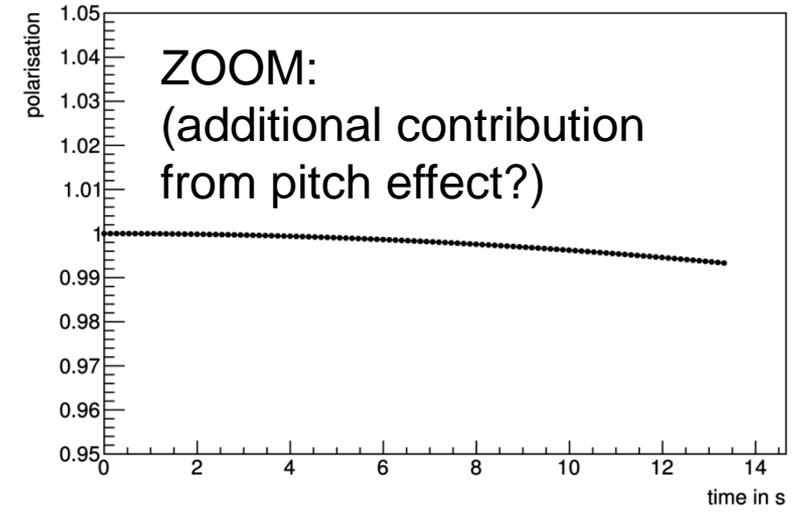
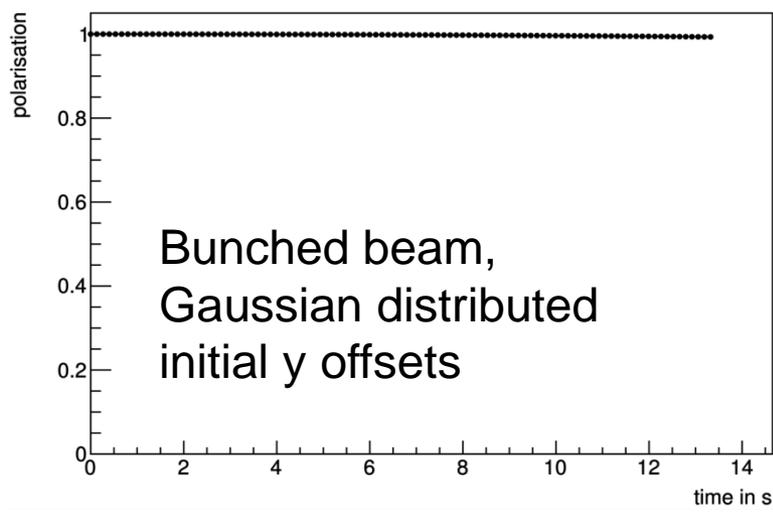
MXS: +2



**Evolution is the same  
for different settings**

- Several settings leading to vanishing path-lengthening:

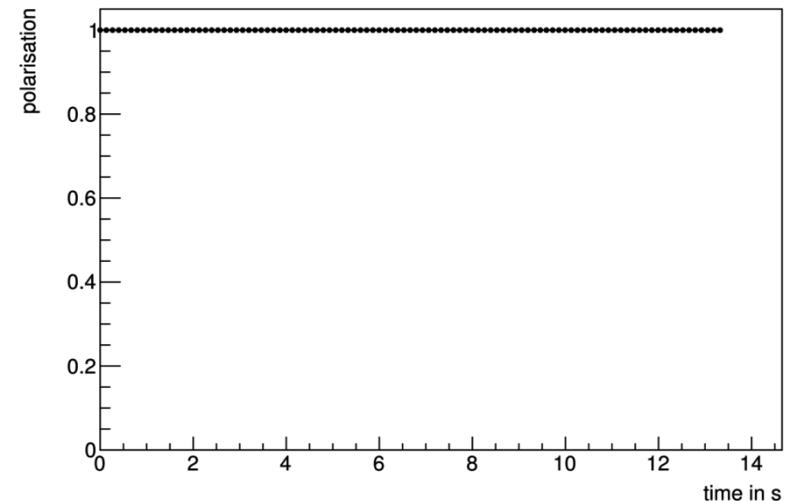
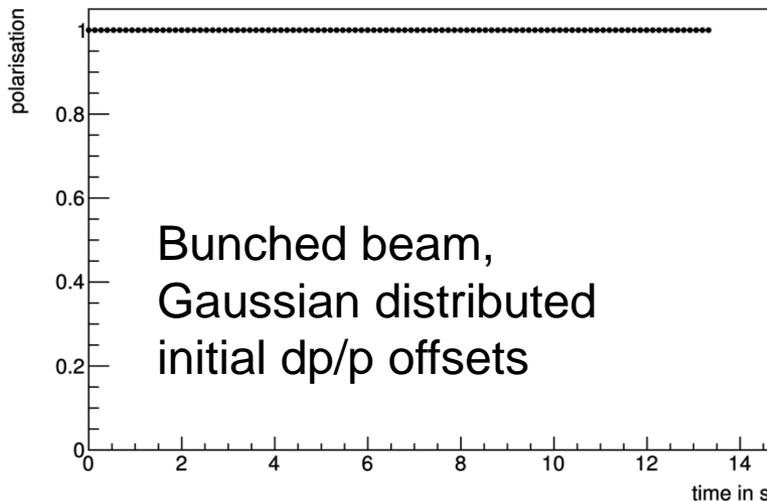
#	MXS / m <sup>-3</sup>	MXL / m <sup>-3</sup>	MXG / m <sup>-3</sup>
1	0	-0.506 (-0.506)	0.775
2	1	-0.522 (-0.522)	0.497
3	2	-0.538 (-0.538)	0.218
4	3	-0.554 (-0.554)	-0.061
5	4	-0.570 (-0.570)	-0.340



- Several settings to fulfill the condition:

$$\left[ \alpha_1 + \frac{3}{2\gamma_0^2} \left( \beta_0^2 - \left( \alpha_0 - \frac{1}{\gamma_0^2} \right) \right) \right] = 0$$

#	MXS / m <sup>-3</sup>	MXL / m <sup>-3</sup>	MXG / m <sup>-3</sup>
1	2	2	-0.284
2	-2	-2	0.825



# COSY Toolbox

COSY INFINITY

COSY Lattice

Armadillo

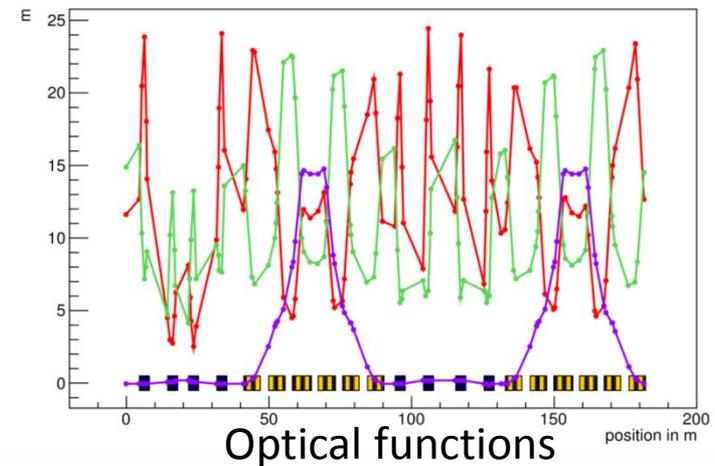
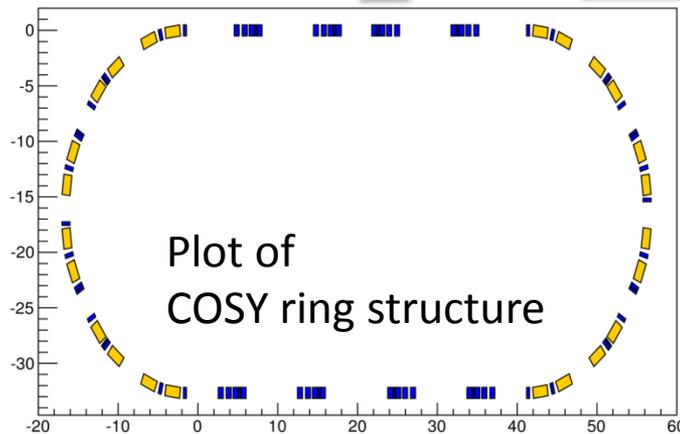
COSY Toolbox

ROOT

- Calculator:
  - Optical functions
  - Orbit
  - Spintune ...
- Tracker:
  - Static maps
  - RF maps (\*)

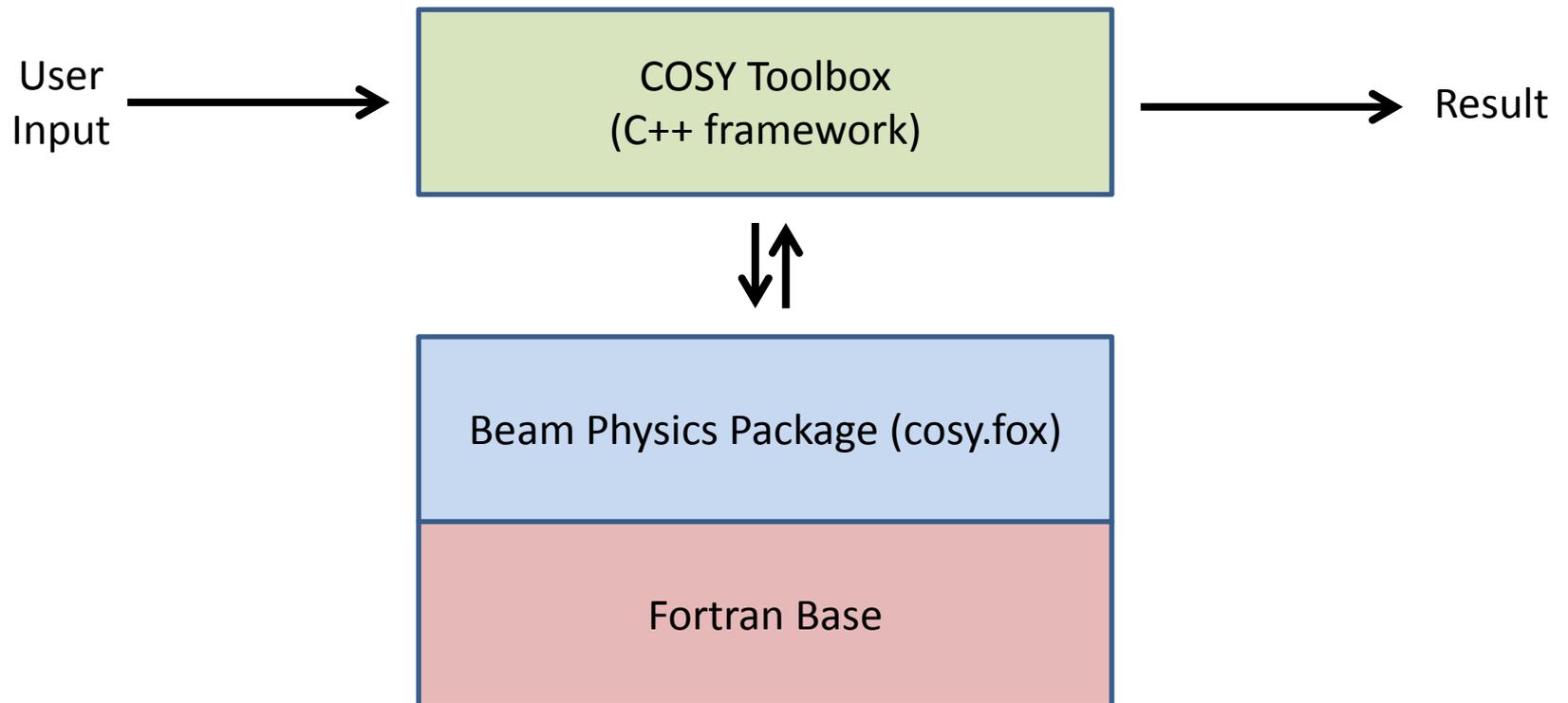
- linear algebra operations
- SVD, pseudo-inverse

- plotting
- storage (ROOT files/Trees)

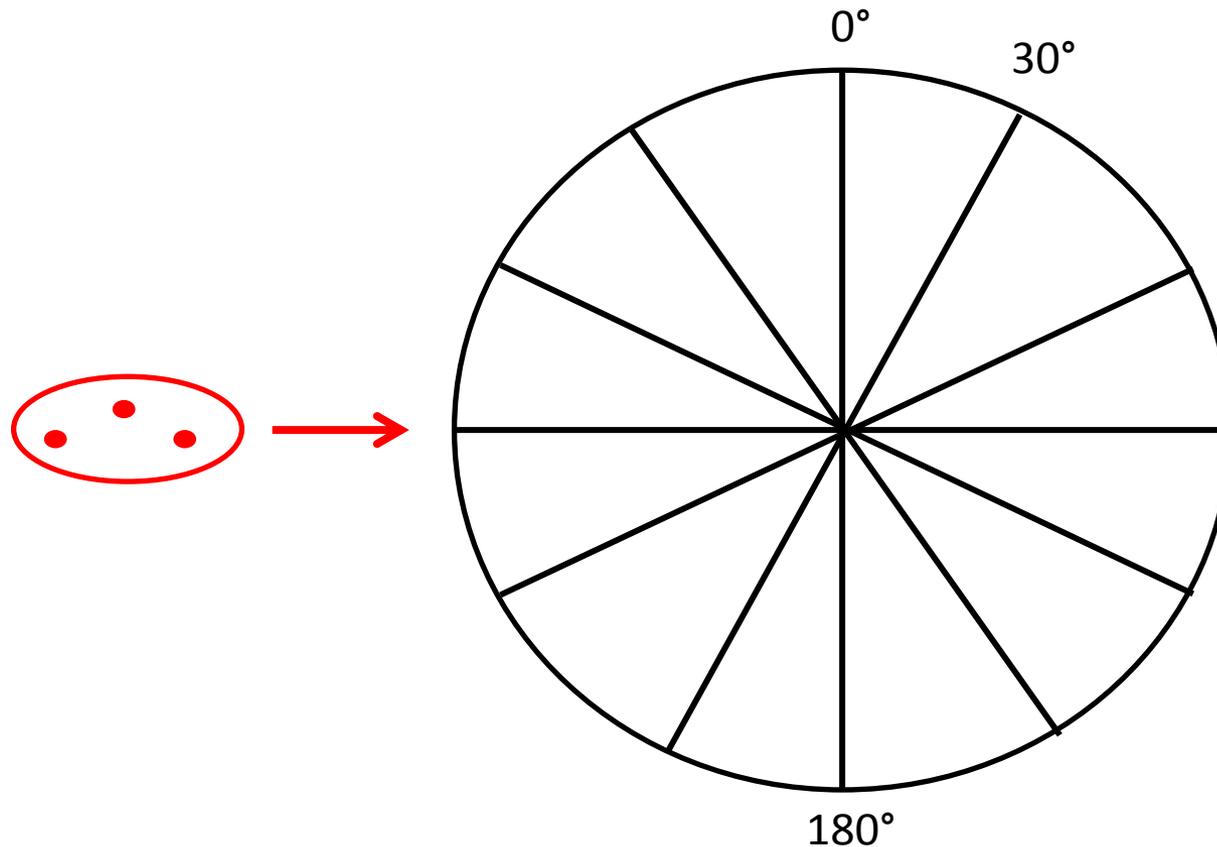


# New method

- New level on top:



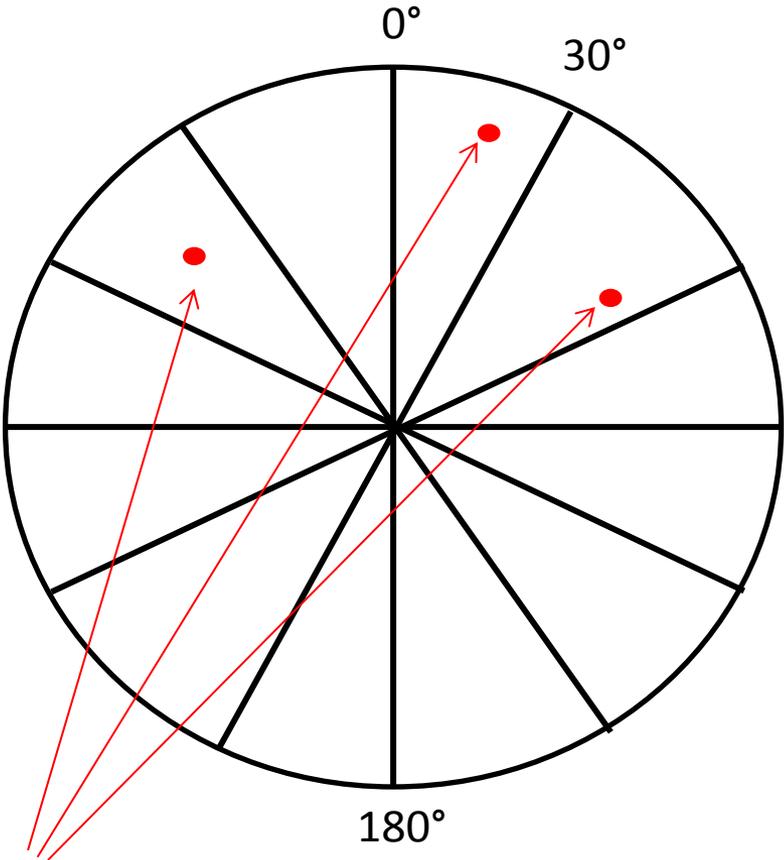
- Example with 12 maps:



**Bunch** approaching RF map

# RF-Tracking illustration

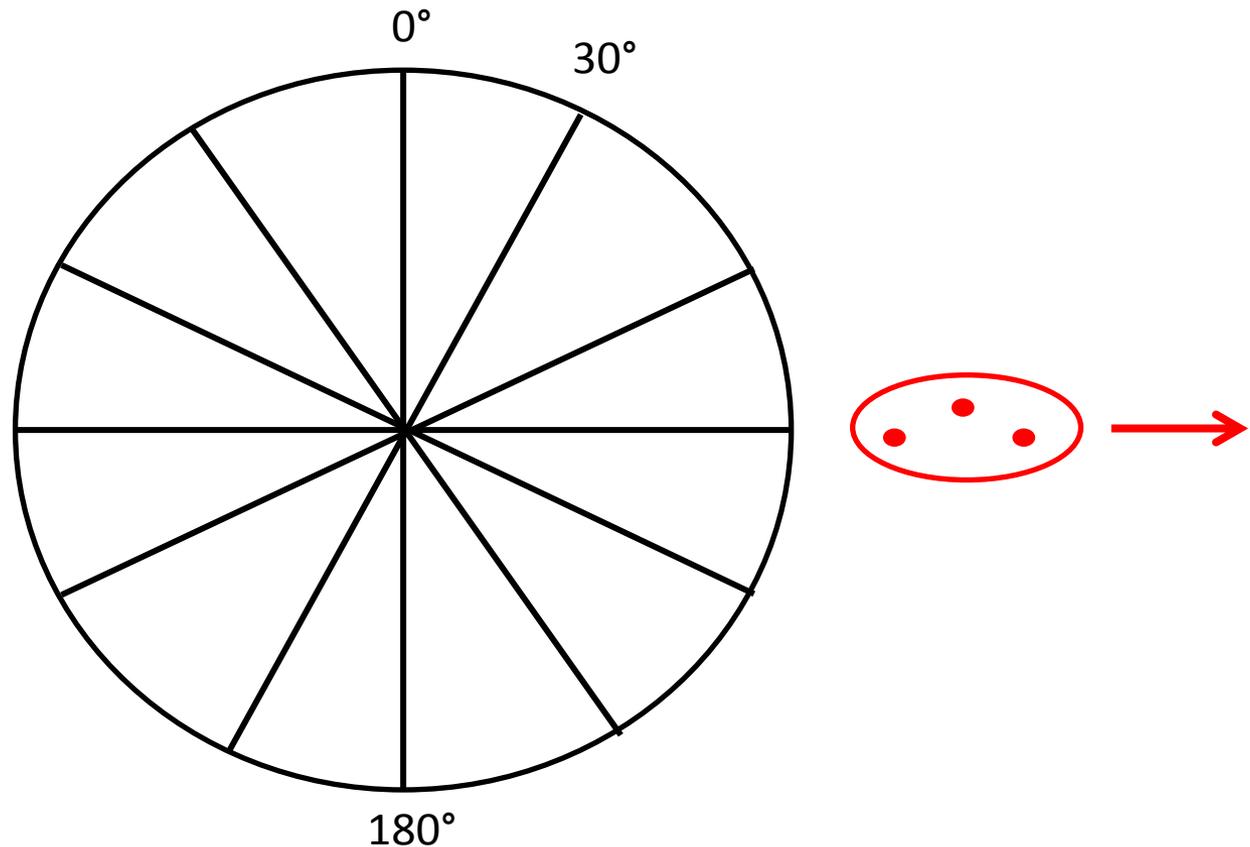
➤ Example with 12 maps:



**Particles** are tracked though map corresponding to phase

# RF-Tracking illustration

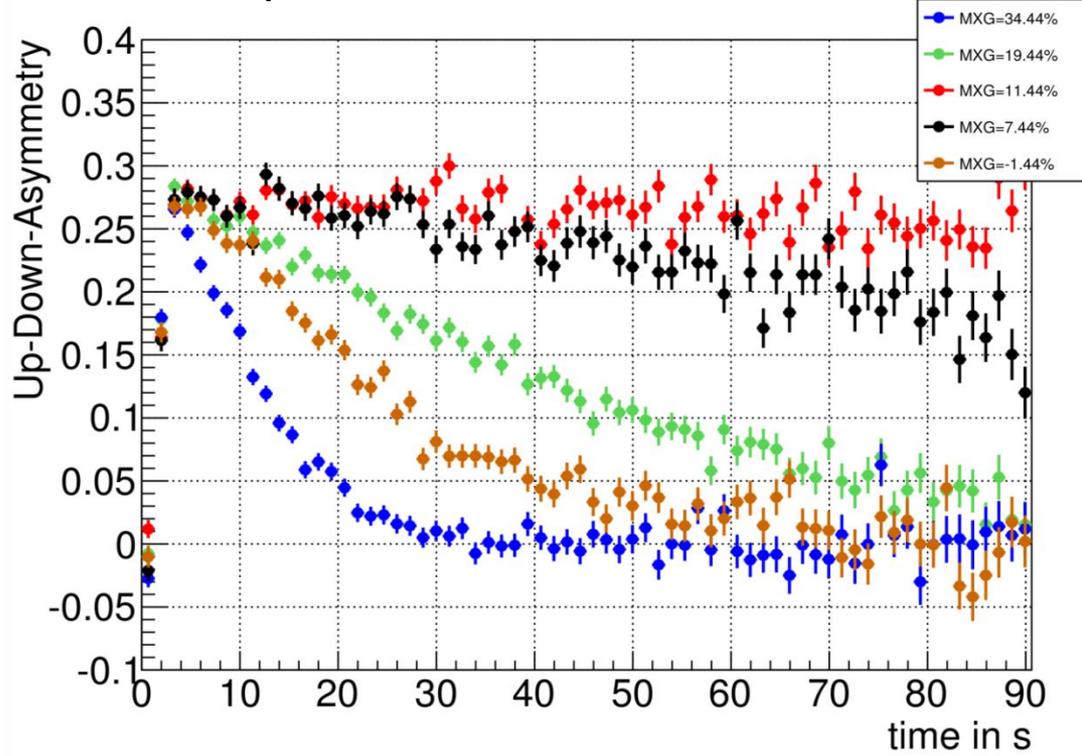
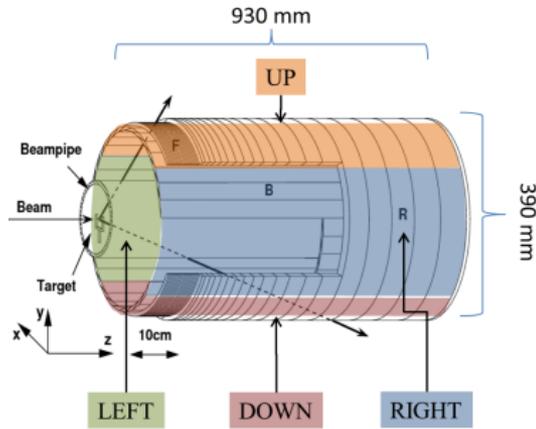
- Example with 12 maps:



**Bunch** passed RF map

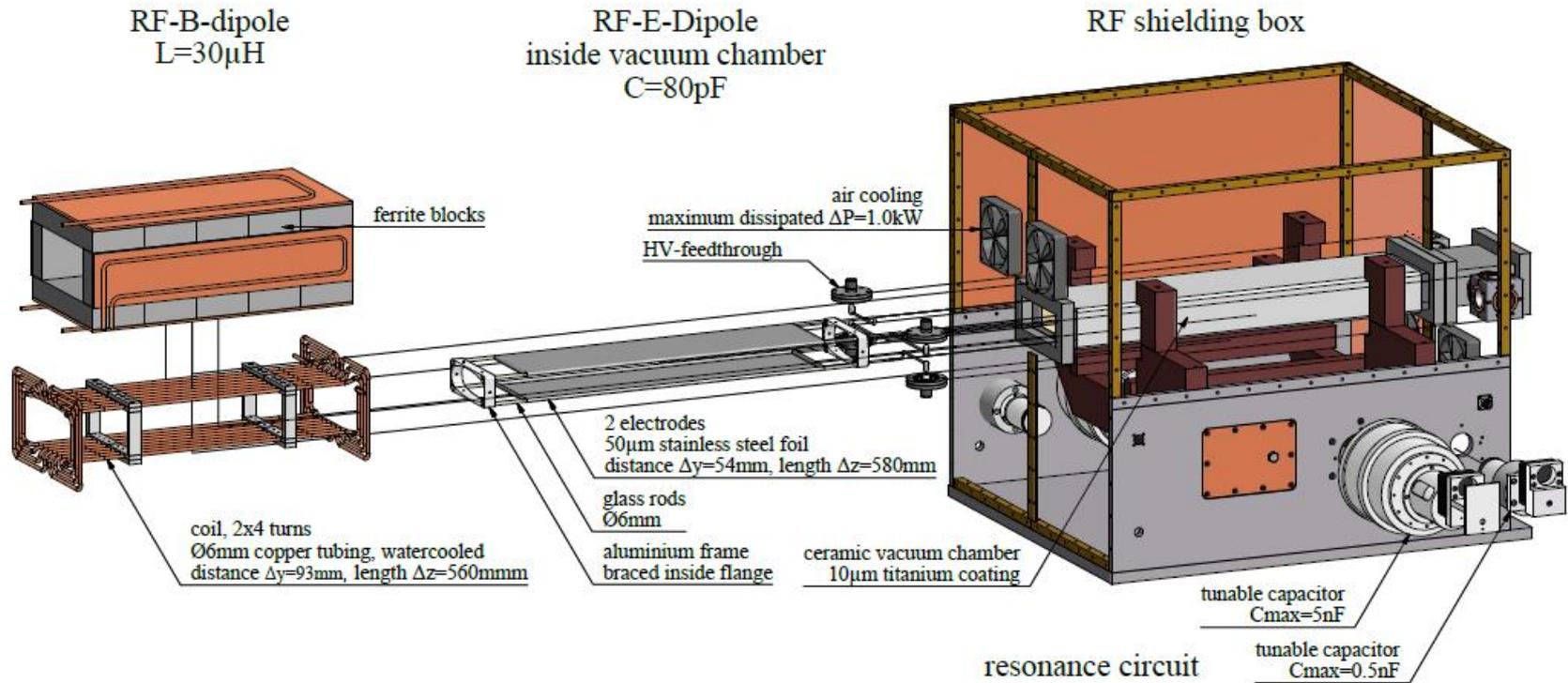
# Spin Coherence Time II

- SCT studies performed during last beam time:
  - Polarised deuterons @ 970 MeV/c
  - Electron-cooled
  - „Heated“ in 1 direction (horizontally or longitudinally)
  - Beam steered on target to measure polarisation over time



- Polarisation evolution for different sextupole configurations

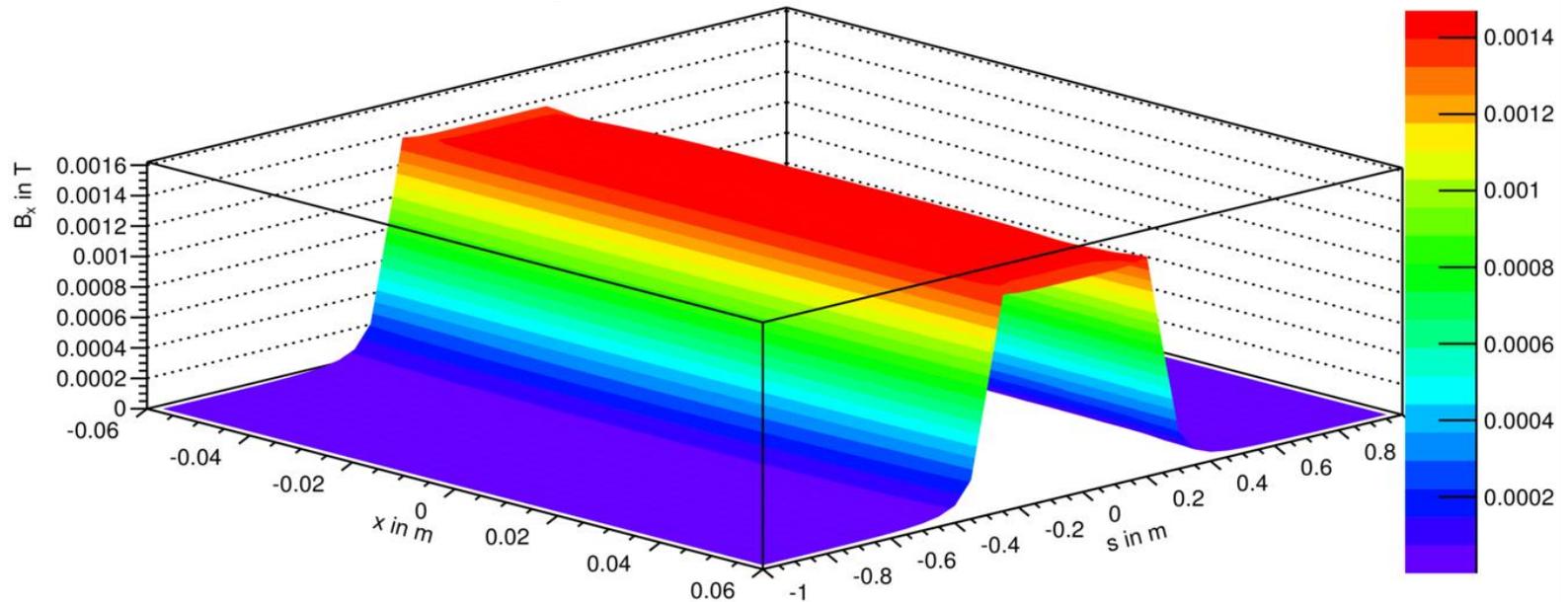
# The RF-E×B-Dipole



Courtesy: Sebastian Mey

# 1st step: Field Maps

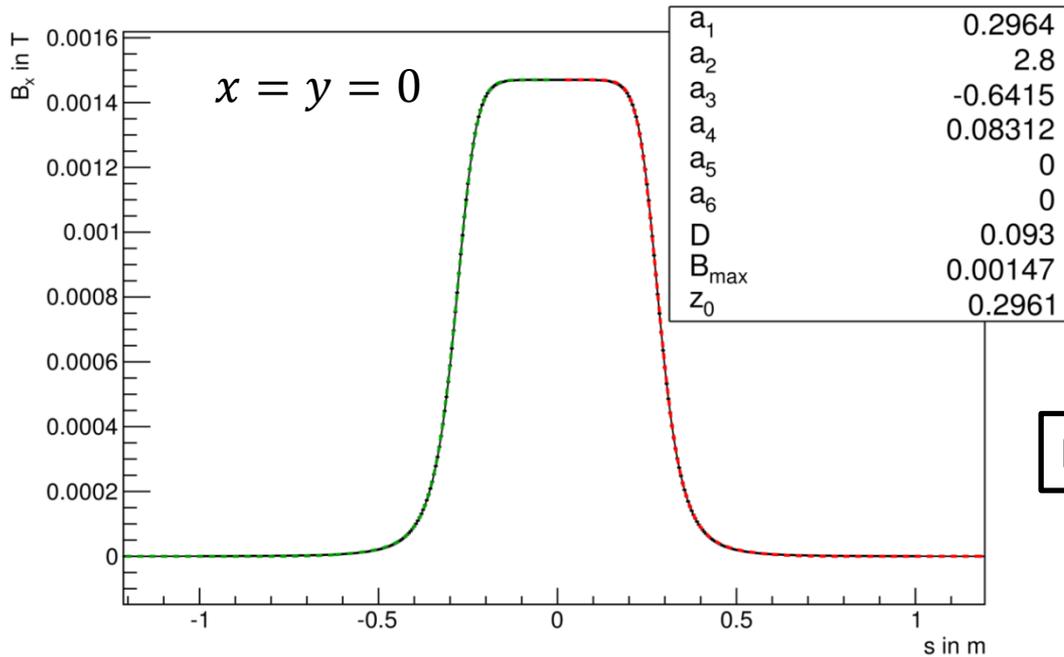
- This Version: Field maps provided by R. Gebel and S. Mey for the existing RF-WienFilter
  - Static field for maximum current / voltage
  
- Example for radial magnetic field



# 1st step: Field Maps

- Simplification (quasi-static approach):
  - $F(x, s, t) = F(x, s) \cdot \cos(\omega t)$
  - $F(x, s) = F_{max} \cdot f(x) \cdot f(s)$
  - $f(x)$ : Polynomial function
  - $f(s)$ : Enge functions for entrance and exit
- Field change during one pass is included.

# 1st step: Field Maps

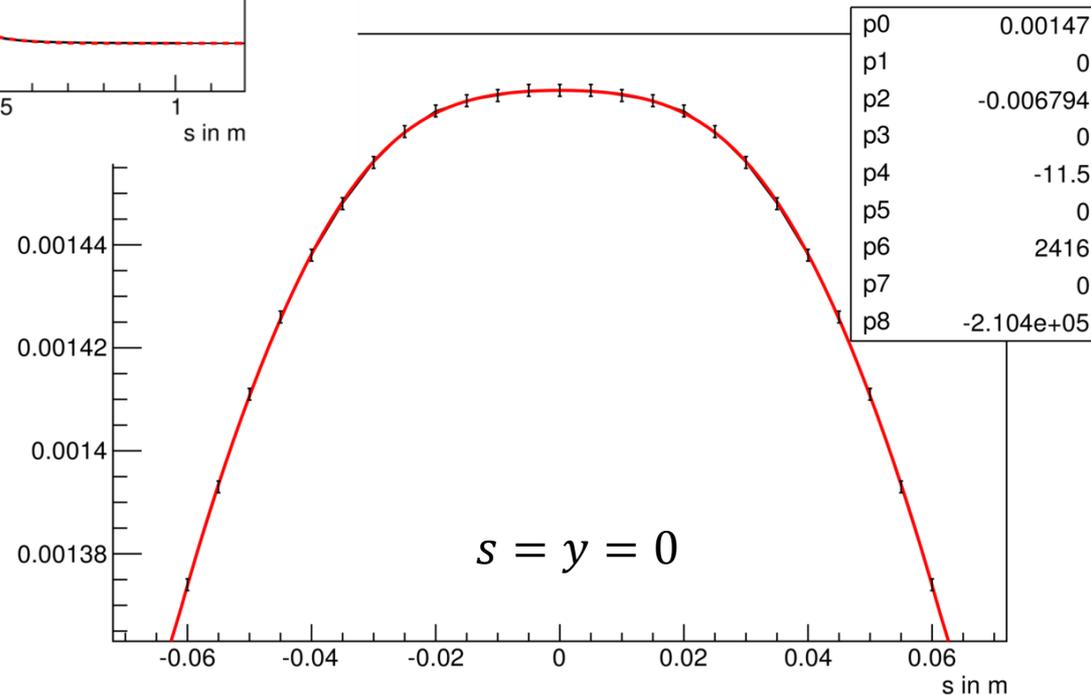


- $F(x, s, t) = F(x, s) \cdot \cos(\omega t)$
- $F(x, s) = F_{max} \cdot f(x) \cdot g(s)$
- $f(x)$ : Polynomial function
- $g(s)$ : Enge functions

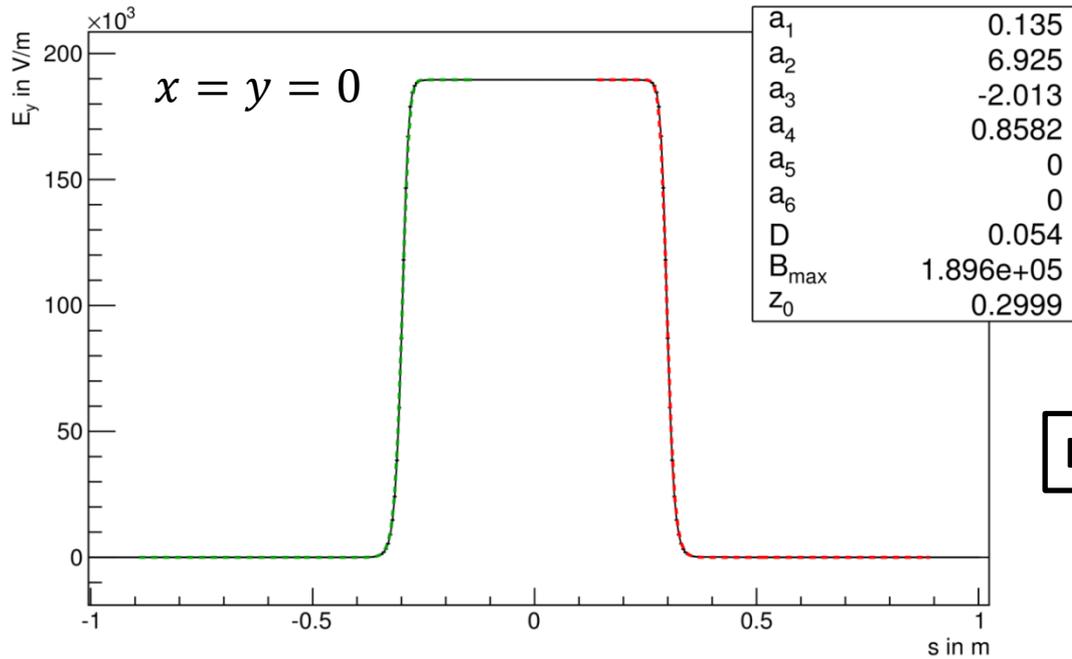
**Magnetic**

Enge-function:

$$h(z) = \frac{1}{1 + \exp(a_1 + a_2 \cdot \left(\frac{z}{D}\right) + \dots + a_6 \left(\frac{z}{D}\right)^5)}$$



# 1st step: Field Maps

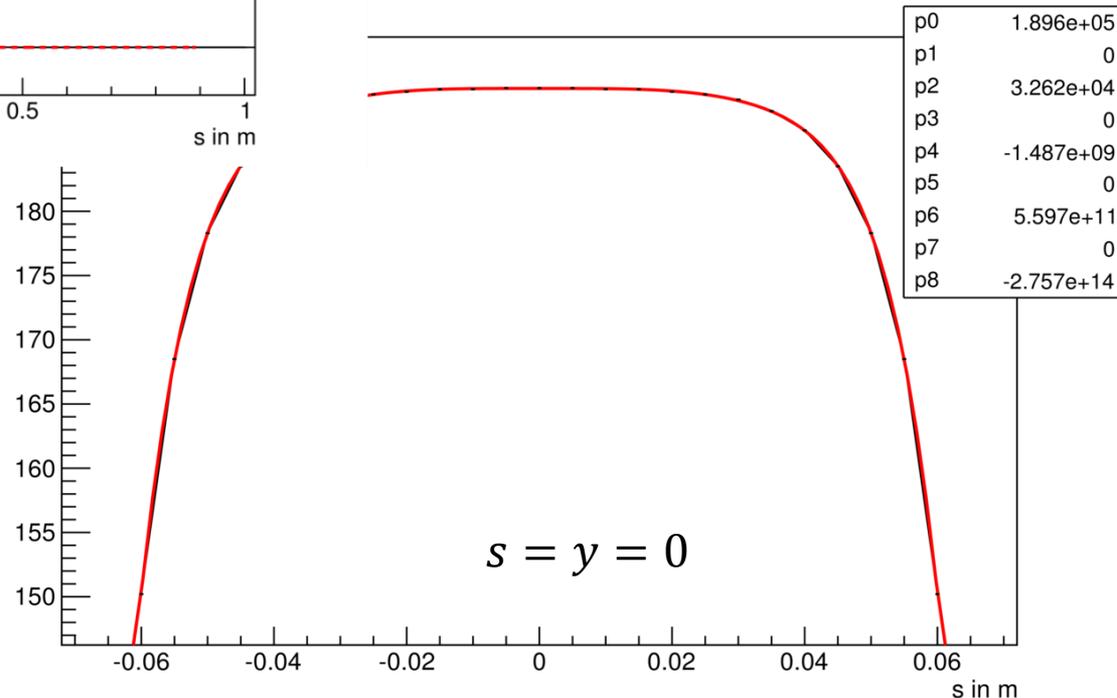


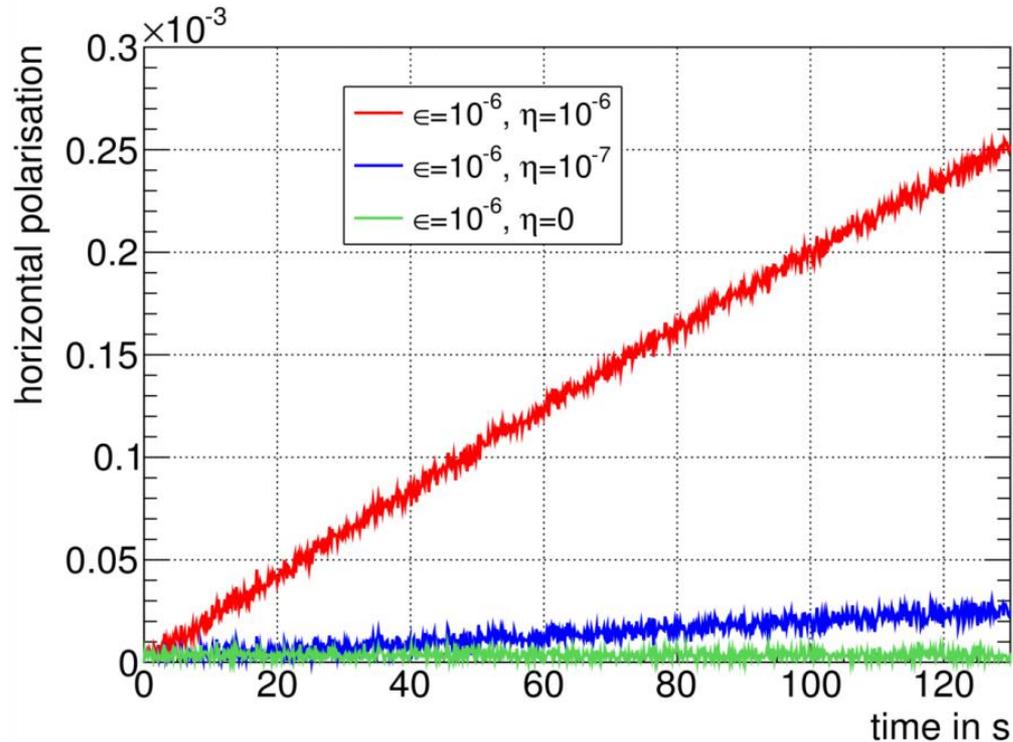
- $F(x, s, t) = F(x, s) \cdot \cos(\omega t)$
- $F(x, s) = F_{max} \cdot f(x) \cdot g(s)$
- $f(x)$ : Polynomial function
- $g(s)$ : Enge functions

**Electric**

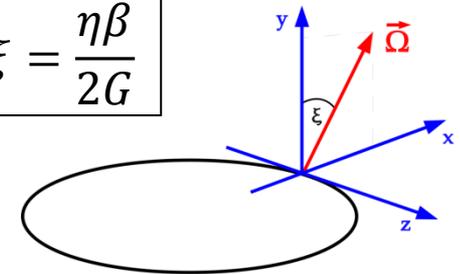
Enge-function:

$$h(z) = \frac{1}{1 + \exp(a_1 + a_2 \cdot \left(\frac{z}{D}\right) + \dots + a_6 \left(\frac{z}{D}\right)^5)}$$





$$\tan \xi = \frac{\eta\beta}{2G}$$



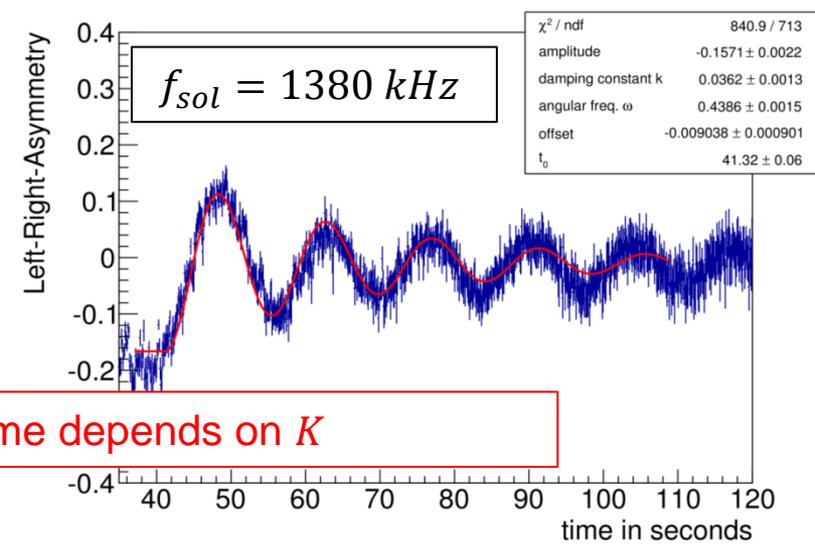
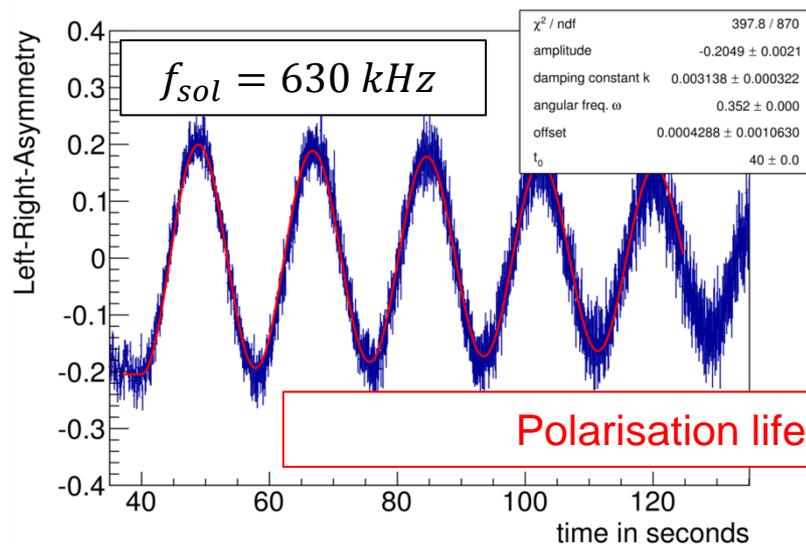
➤ Buildup per revolution:

$$\Theta = \frac{\psi\xi}{2} = 4\pi \cdot \frac{\epsilon\xi}{2}$$

## Deuterons:

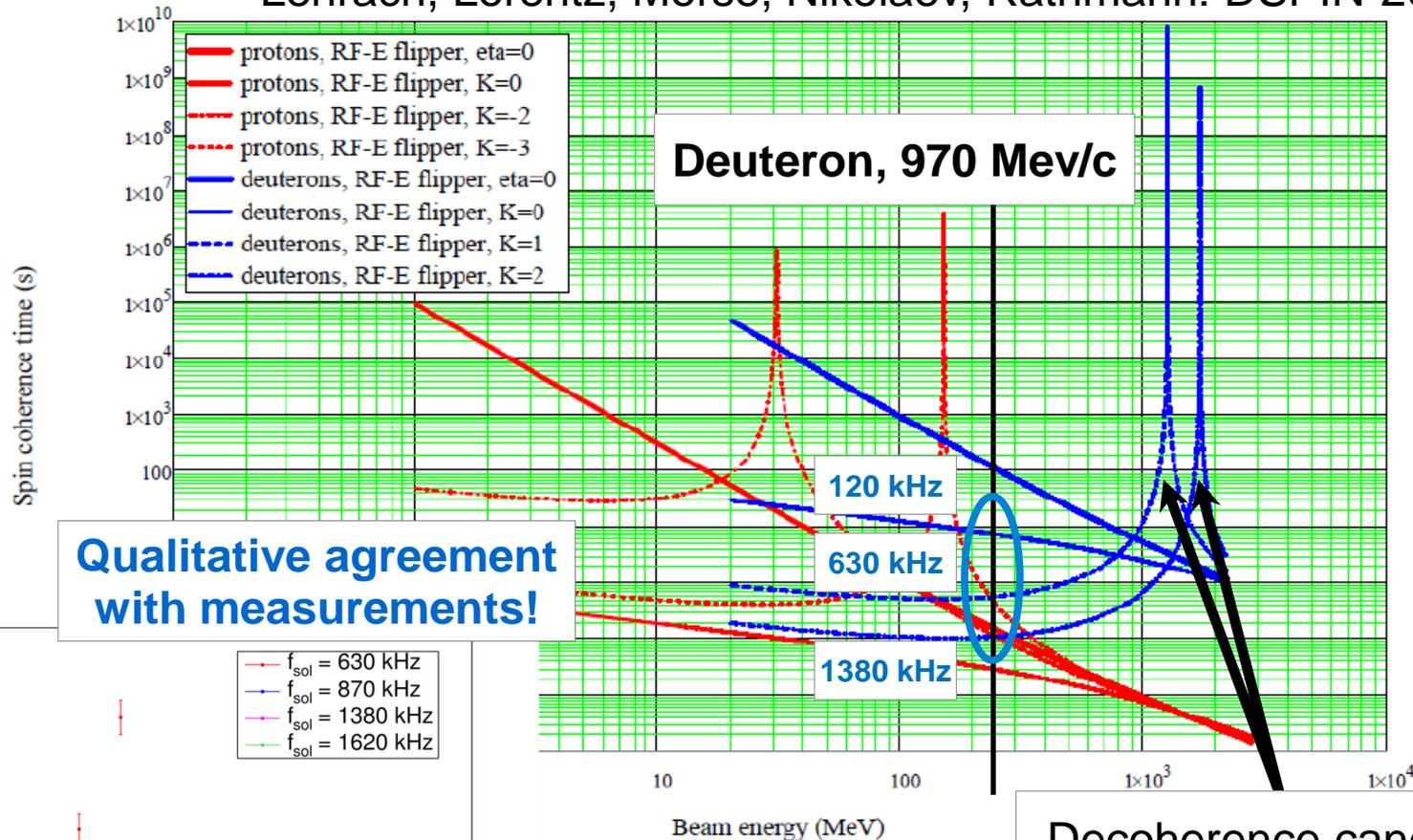
p	$\epsilon$	L	E	B	$\eta$	$\Theta$
970 MeV/c	$10^{-6}$	0.6 m	12.2 kV/m	0.09 mT	$10^{-9}$	$10^{-14}$
970 MeV/c	$10^{-4}$	0.6 m	1.2 MV/m	8.9 mT	$10^{-9}$	$10^{-12}$

- Precursor experiment: RF-ExB-resonance to build-up EDM signal
- First studies using an RF-Solenoid to investigate induced spin resonances
- Resonance condition:  $f_{sol} = |K + G\gamma| \cdot f_{rev}$

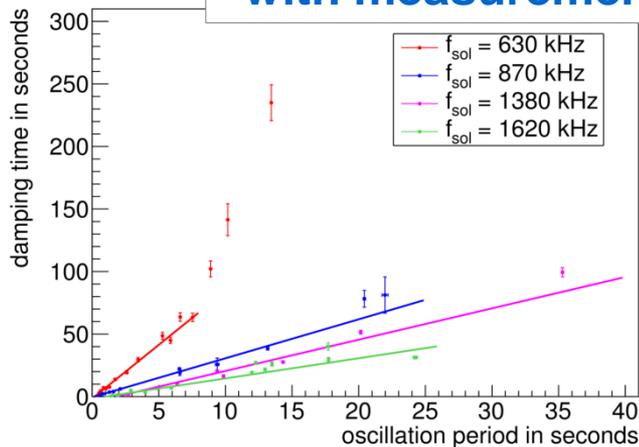


Polarisation lifetime depends on  $K$

Lehrach, Lorentz, Morse, Nikolaev, Rathmann: DSPIN-2011

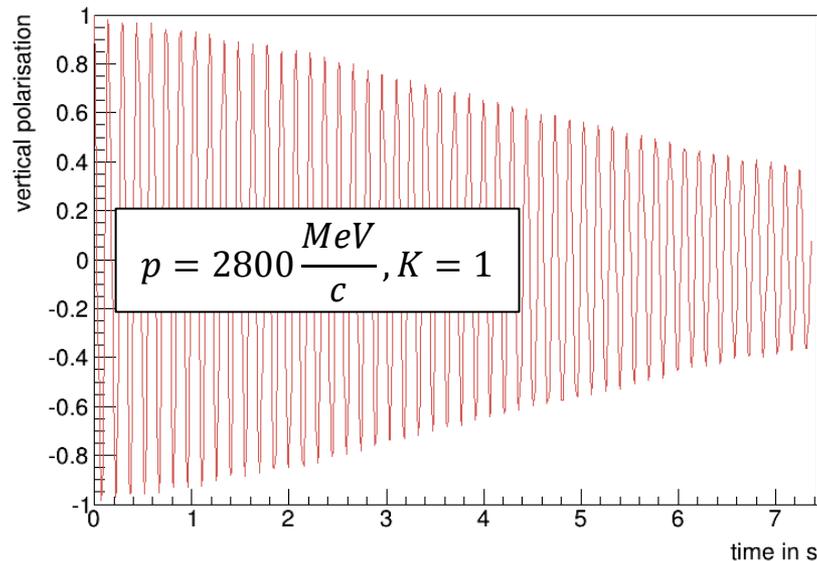
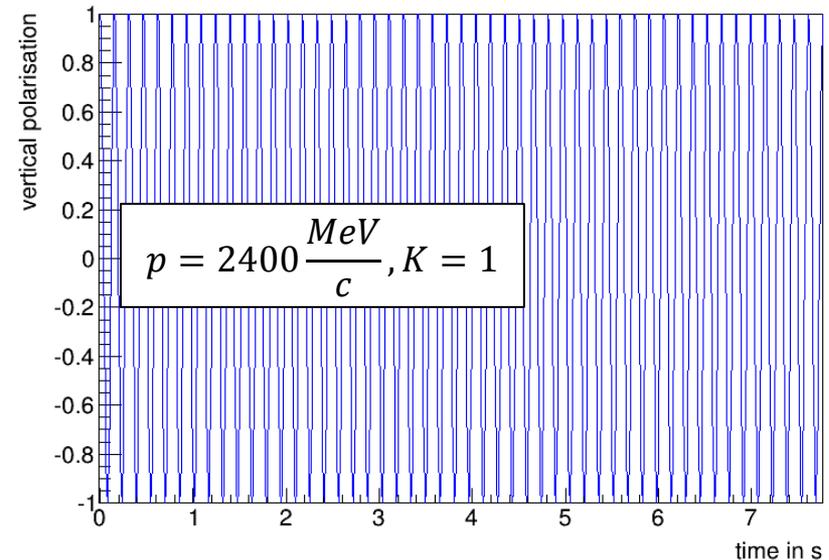
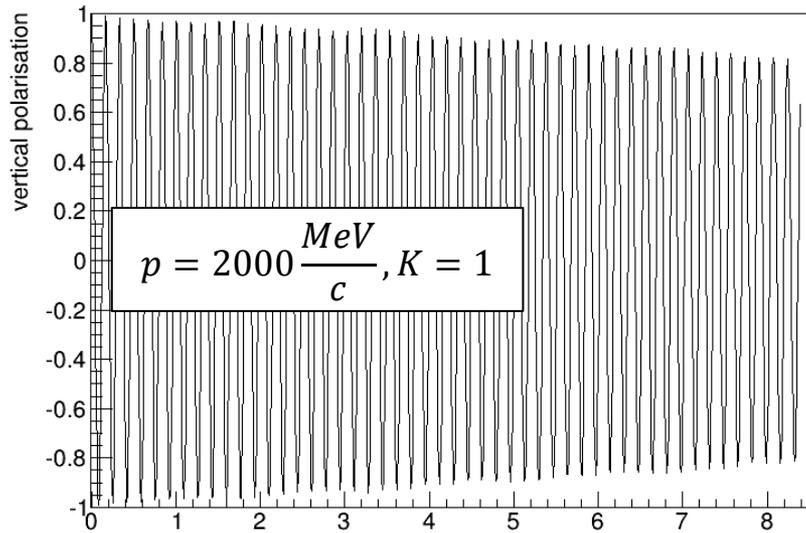


**Qualitative agreement with measurements!**



Decoherence cancelled for

$$C = 1 - \frac{\eta}{\beta^2} \left( 1 + \frac{K}{G\gamma} \right) = 0$$



$$C = 1 - \frac{\eta}{\beta^2} \left( 1 + \frac{K}{G\gamma} \right) = 0$$

- fulfilled for  $p \approx 2402 \text{ MeV}/c$
- Important for different ideas of EDM measurements in storage rings.

# Statistical Sensitivity for electric/combined-ring

$$\sigma \approx \frac{\hbar}{\sqrt{NfT\tau_p PEA}}$$

$P$	beam polarization	0.8
$\tau_p$	Spin coherence time/s	1000
$E$	Electric field/MV/m	10
$A$	Analyzing Power	0.6
$N$	nb. of stored particles/cycle	$4 \times 10^7$
$f$	detection efficiency	0.005
$T$	running time per year/s	$10^7$

$\Rightarrow \sigma \approx 10^{-29} \text{ e}\cdot\text{cm}/\text{year}$  (for magnetic ring  $\approx 10^{-24} \text{ e}\cdot\text{cm}/\text{year}$ )

Expected signal  $\approx 3 \text{ rad/s}$  (for  $d = 10^{-29} \text{ e}\cdot\text{cm}$ )

(BNL proposal)

# Statistical Sensitivity for magnetic ring (COSY)

$$\sigma \approx \frac{\hbar}{2} \frac{G\gamma^2}{G+1} \frac{U}{E \cdot L} \frac{1}{\sqrt{NfT\tau_pPA}}$$

$G$	anomalous magnetic moment	
$\gamma$	relativistic factor	1.13
	$p = 1 \text{ GeV}/c$	
$U$	circumference of COSY	180 m
$E \cdot L$	integrated electric field	$0.1 \cdot 10^6 \text{ V}$
$N$	nb. of stored particles/cycle	$2 \cdot 10^9$

$$\Rightarrow \sigma \approx 10^{-25} \text{ e}\cdot\text{cm}/\text{year}$$

# Systematics

One major source:

Radial  $B$  field mimics an EDM effect:

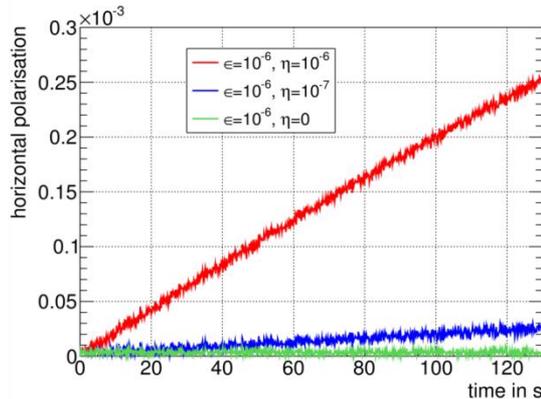
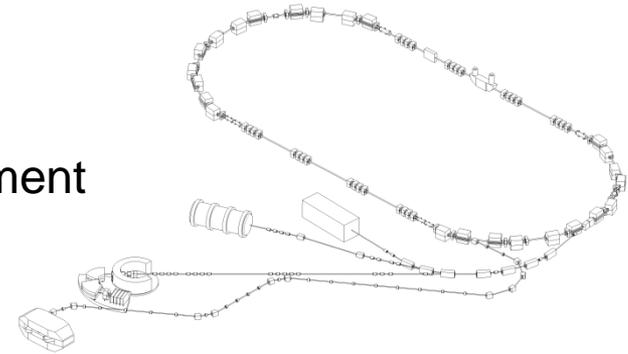
- Difficulty: even small radial magnetic field,  $B_r$  can mimic EDM effect if  $:\mu B_r \approx dE_r$
- Suppose  $d = 10^{-29} \text{ e}\cdot\text{cm}$  in a field of  $E = 10 \text{ MV/m}$
- This corresponds to a magnetic field:

$$B_r = \frac{dE_r}{\mu_N} = \frac{10^{-22} \text{ eV}}{3.1 \cdot 10^{-8} \text{ eV/T}} \approx 3 \cdot 10^{-17} \text{ T}$$

(Earth Magnetic field  $\approx 5 \cdot 10^{-5} \text{ T}$ )

Solution: Use two beams running clockwise and counter clockwise, separation of the two beams is sensitive to  $B_r$

- EDM measurements in storage rings
  - Feasibility studies and precursor experiment at COSY/Jülich



- EDM related polarisation build-up using induced spin resonance
  - RF-ExB-Flipper ( $\rightarrow$  talk: S.Mey)
- *Outlook: systematic studies concerning beam alignment and field quality*

- Preservation of polarisation mandatory
  - Sextupole corrections
- *Outlook: further investigation of SCT*

