

Application of chiral YN interactions to light hypernuclei

Andreas Nogga, Forschungszentrum Jülich

7th LENPIC meeting, March 11-13, 2024, Bonn, Germany



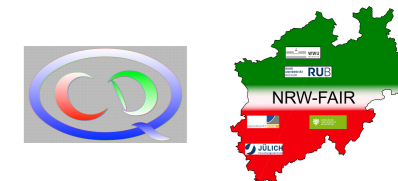
JÜLICH
Forschungszentrum



- Motivation
- YN and YY interactions
- SRG evolution of (hyper-)nuclear interactions
- Determination of CSB contact interactions and Λn scattering length
- Application to $A = 7$ and 8 hypernuclei
- Uncertainty of Λ separation energies and size of chiral 3BF contributions
- Chiral YNN interactions

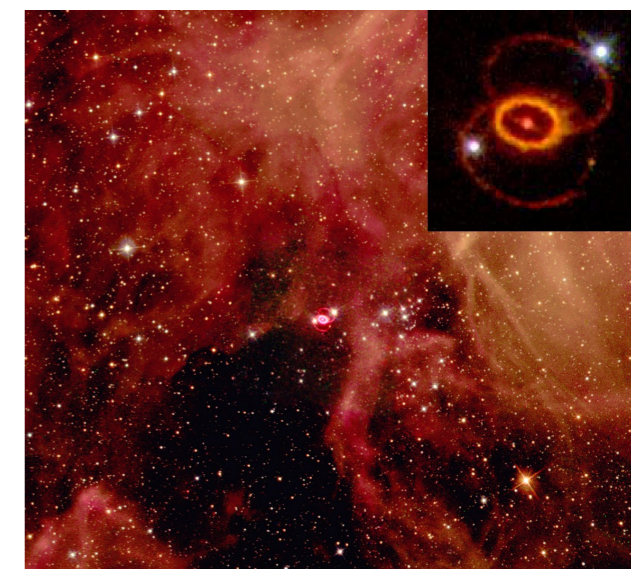
in collaboration with Johann Haidenbauer, **Hoai Le**, Ulf Meißner

Hypernuclear interactions

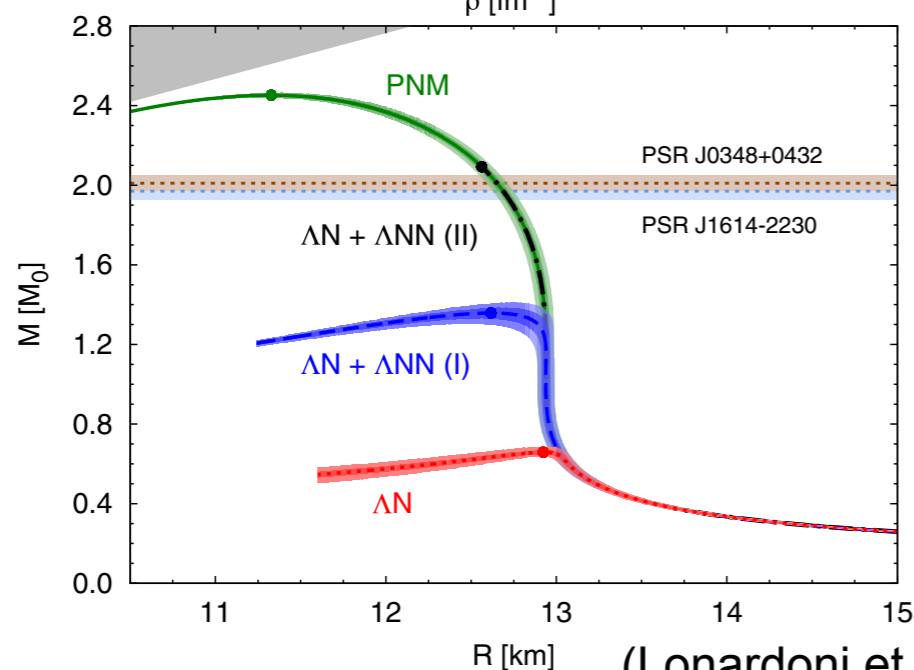
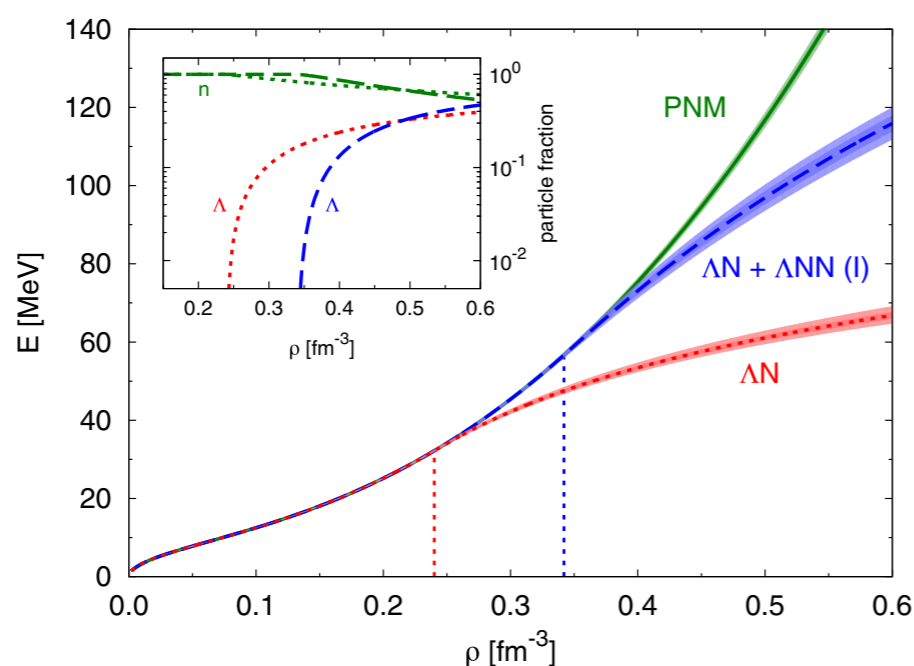


Why is understanding hypernuclear interactions interesting?

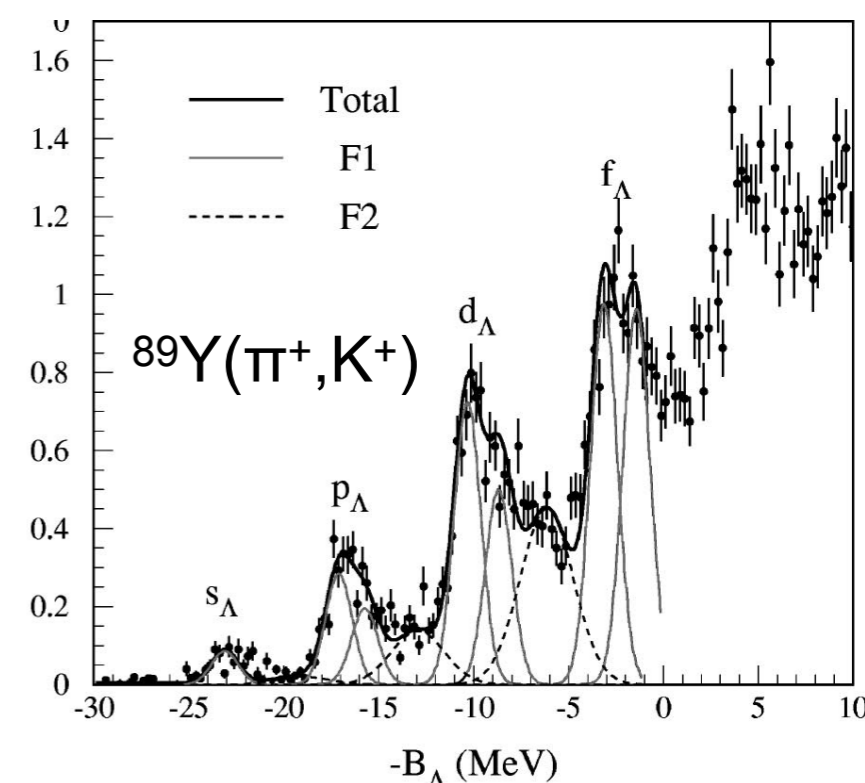
- hyperon contribution to the EOS, neutron stars, supernovae
- "hyperon puzzle"
- Λ as probe to nuclear structure
- flavor dependence of baryon-baryon interactions



(SN1987a, Wikipedia)



(Lonardoni et al. (2015))

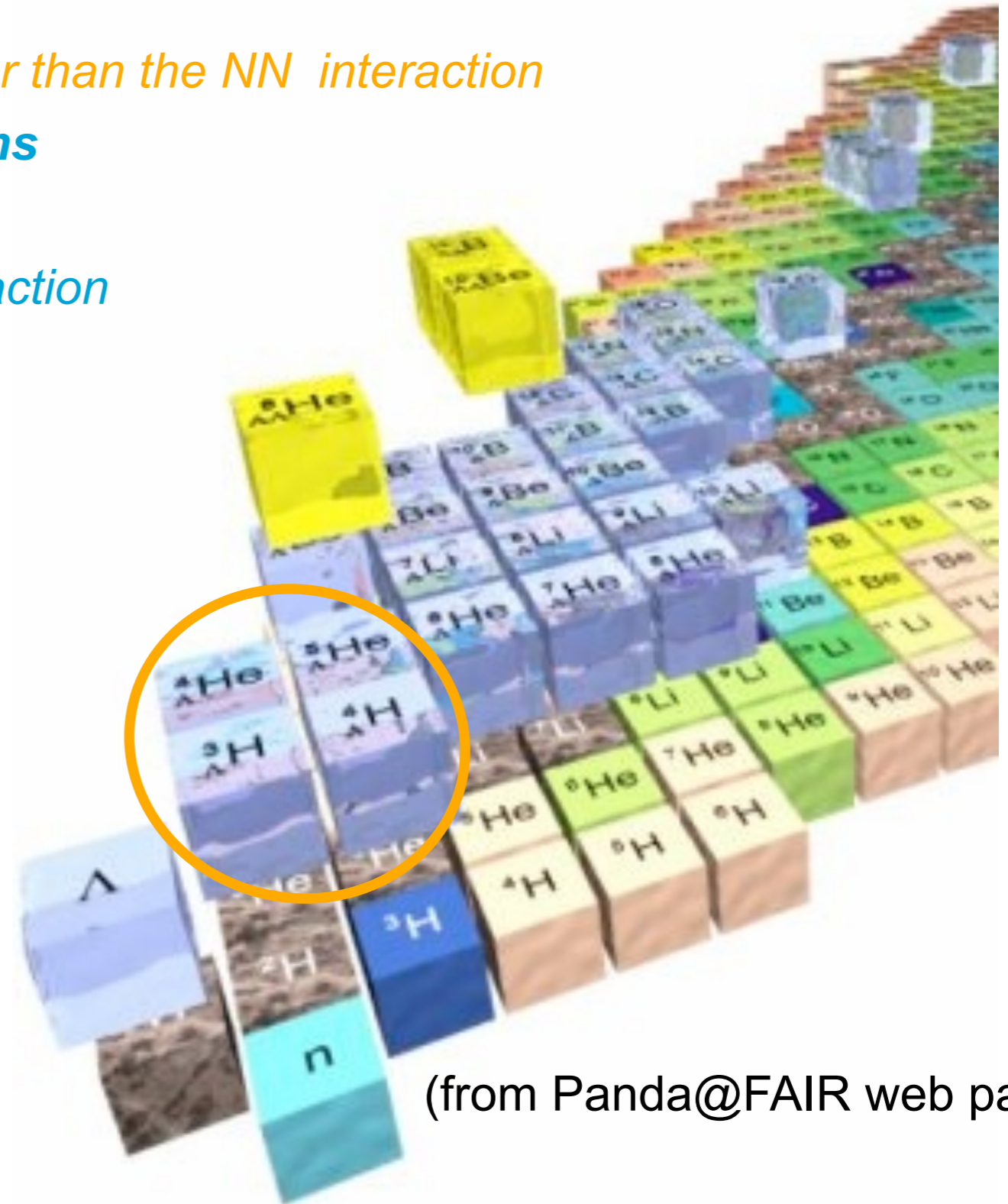


(Hotchi et al. (2001))

Hypernuclei

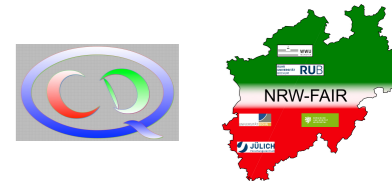
Only few YN data. Hypernuclear data provides additional constraints.

- ΛN interactions are generally weaker than the NN interaction
 - naively: core nucleus + hyperons
 - „separation energies“ are quite independent from NN(+3N) interaction
- no Pauli blocking of Λ in nuclei
 - good to study nuclear structure
 - even light hypernuclei exist in several spin states
- non-trivial constraints on the YN interaction even from lightest ones
- size of YNN interactions?
need to include Λ - Σ conversion!



(from Panda@FAIR web page)

Chiral NN & YN & YY interactions



EFT based approaches

	BB force	3B force	4B force	
LO		—	—	5 (+1) NN/YN (YY) short range parameters
NLO		—	—	23(+5) NN/YN (YY) short range parameters
N ² LO			—	no additional contact terms in NN/YN (YY)

Chiral EFT implements **chiral symmetry of QCD** (adapted from Epelbaum, 2008)

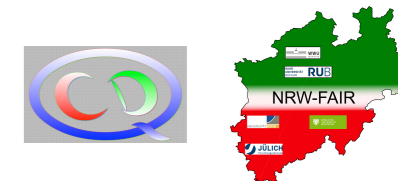
- symmetries constrain exchanges of Goldstone bosons
- relations of two- and three- and more-baryon interactions
- breakdown scale $\approx 600 - 700 \text{ MeV}$
- Semi-local momentum regularization (SMS) up to N²LO (for YN, YY within NRW Fair)

Retain flexibility to adjust to data due to counter terms

Regulator required — cutoff/different orders often used to estimate uncertainty

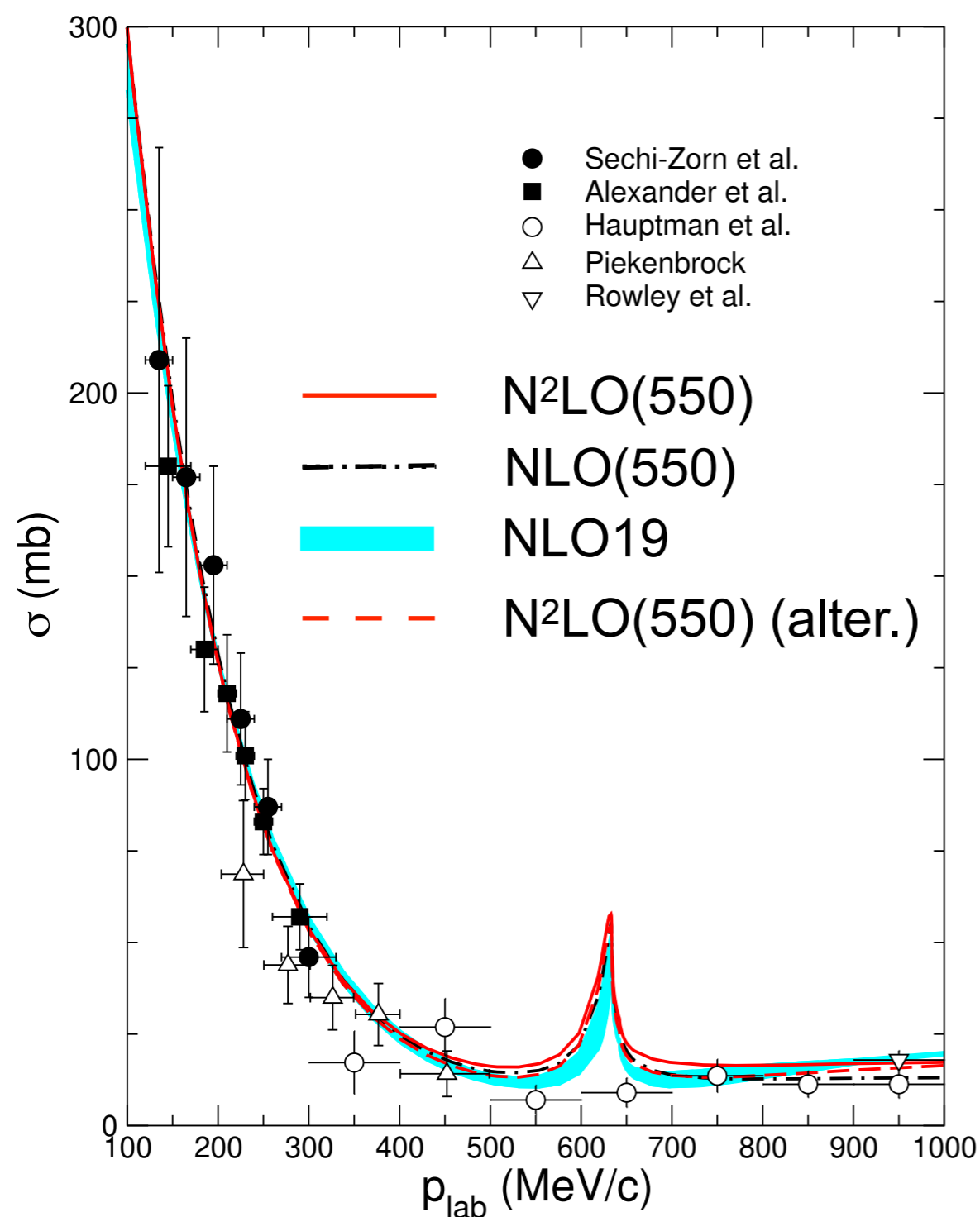
$\Lambda - \Sigma$ and $\Lambda\Lambda - \Sigma\Sigma - \Xi N$ conversion is explicitly included (3BFs only in N²LO)

SMS NLO/N²LO interaction



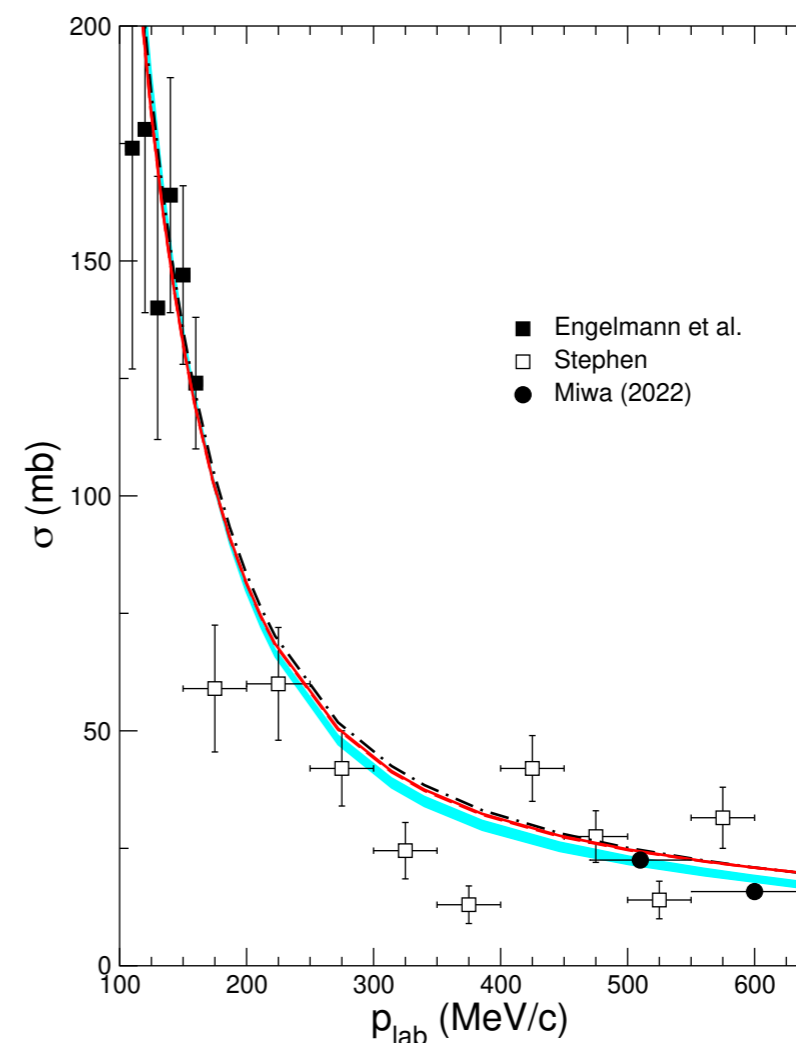
Selected results (show $\Lambda = 550$ MeV, others are very similar in quality)

$\Lambda p \rightarrow \Lambda p$

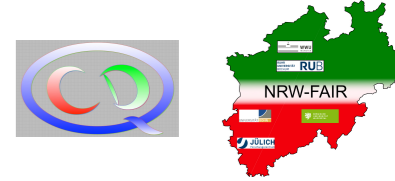


- most relevant cross sections very similar in NLO and N²LO
- similar to NLO19
- alternative fit (see later)

$\Sigma^- p \rightarrow \Lambda n$

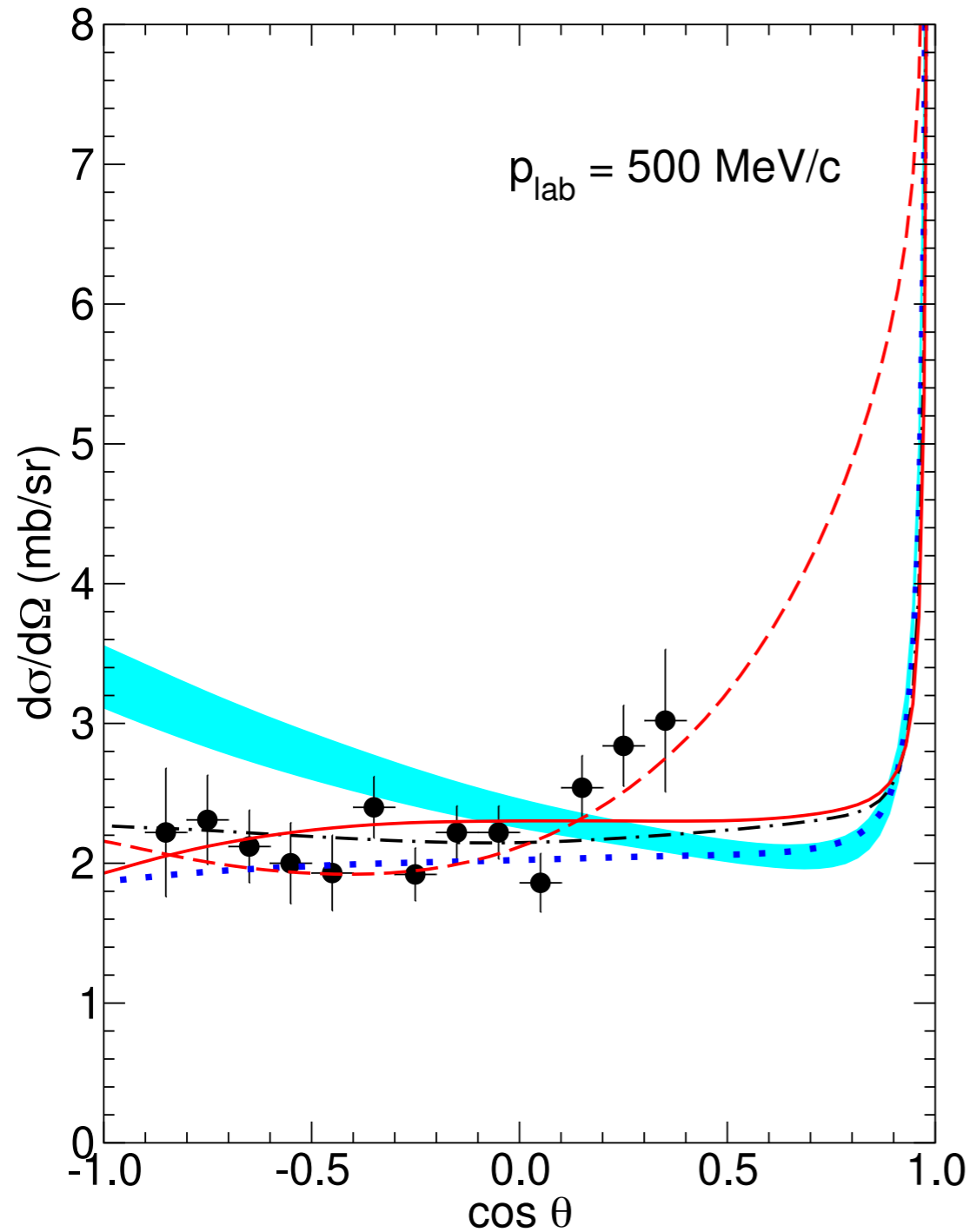


SMS NLO/N²LO interaction

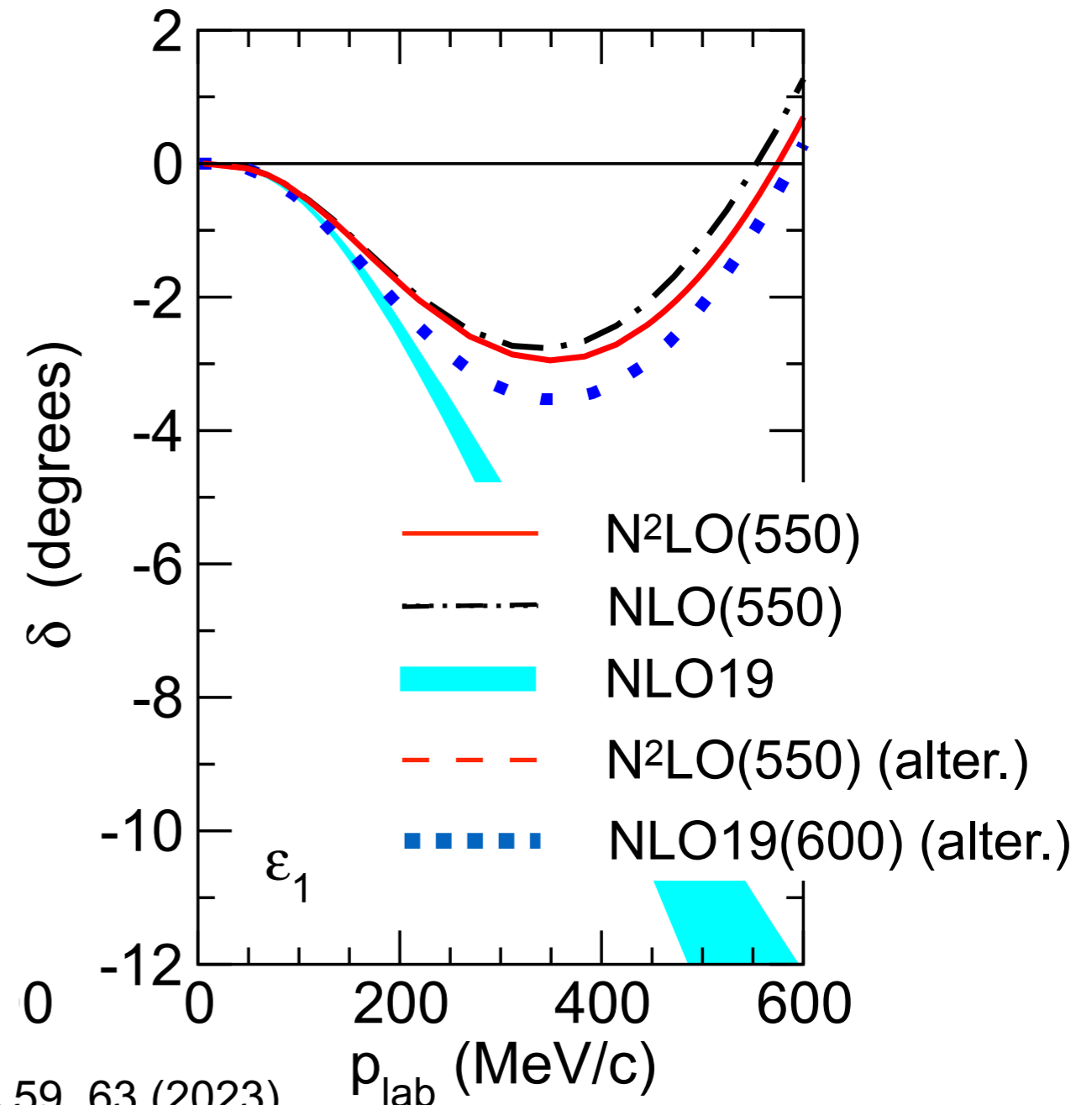


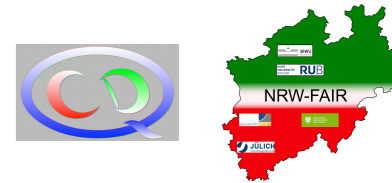
new data (Miwa(2022)) at higher energies provides new constraints!

$$\Sigma^+ p \rightarrow \Sigma^+ p$$



J. Haidenbauer et al. EPJ A 59, 63 (2023).





Similarity renormalization group is by now a **standard tool** to obtain soft effective interactions for various many-body approaches (NCSM, coupled-cluster, MBPT, ...)

Idea: perform a unitary transformation of the NN (and YN interaction) using a cleverly defined "generator"

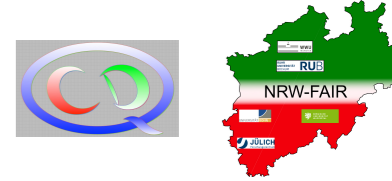
$$\frac{dH_s}{ds} = \left[\underbrace{[T, H(s)]}_{\equiv \eta(s)}, H(s) \right] \quad H(s) = T + V(s)$$

this choice of generator drives $V(s)$ into a diagonal form in momentum space

- $V(s)$ will be **phase equivalent** to original interaction
- short range $V(s)$ will change towards **softer interactions**
- Evolution can be restricted to **2-,3-, ... body level** (approximation)
- $\lambda = \left(\frac{4\mu_{BN}^2}{s} \right)^{1/4}$ is a measure of the width of the interaction in momentum space
- **dependence** of results on λ or s is a measure for **missing terms**

(Bogner et al., 2007)

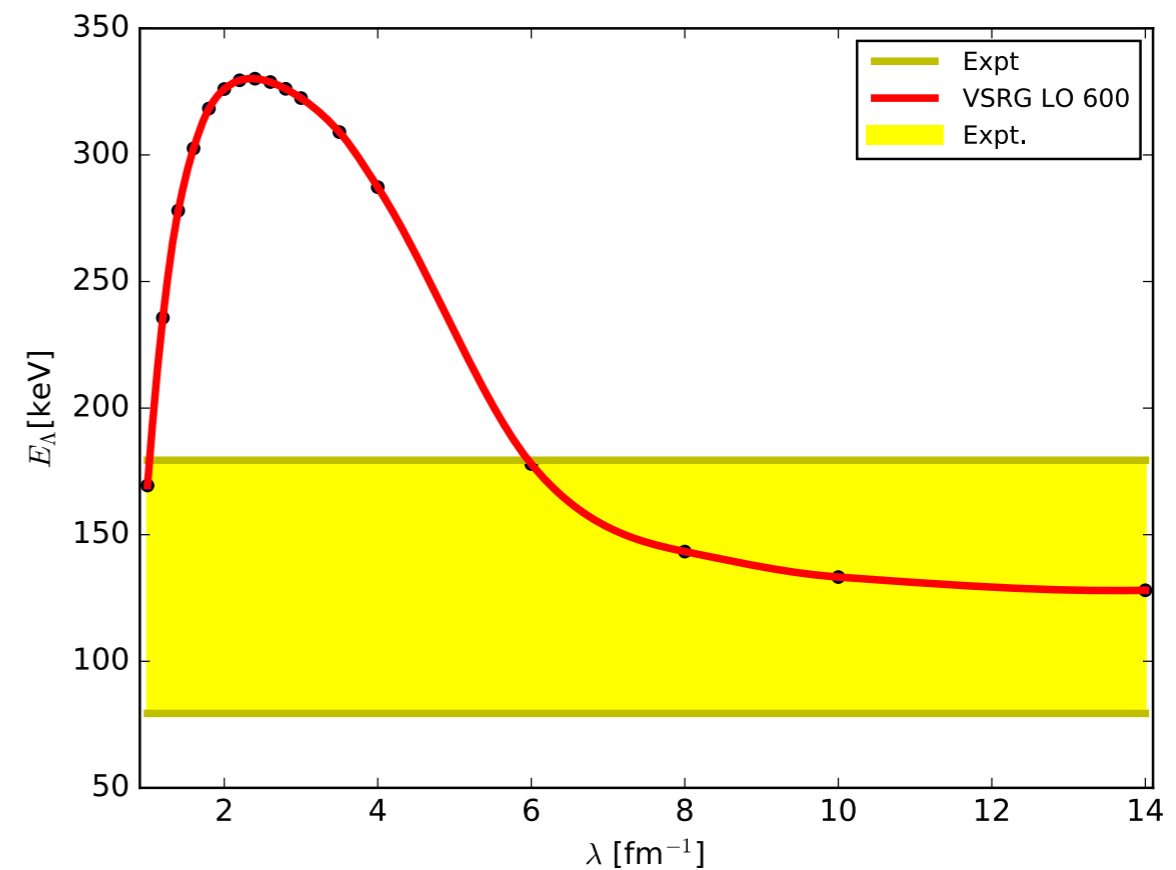
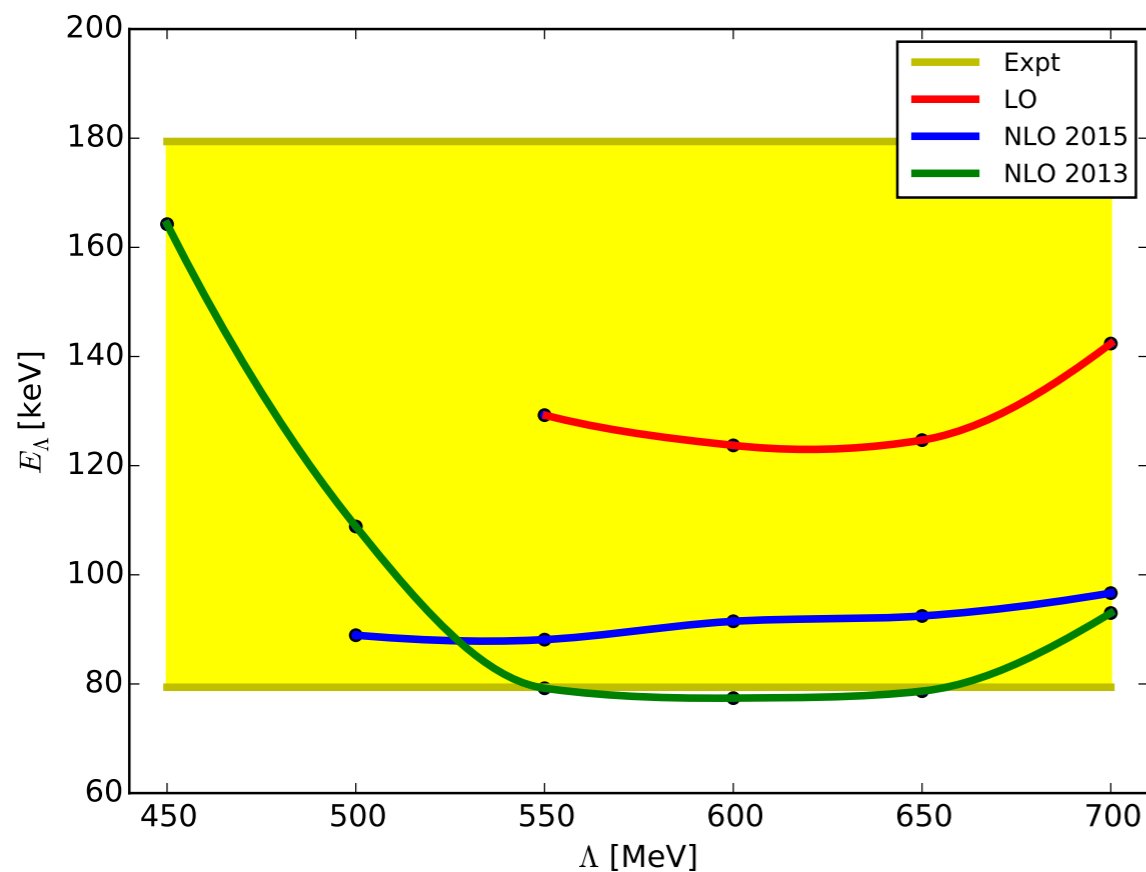
Induced 3BF ...



SRG parameter dependence is significant when NN and YN interactions are evolved

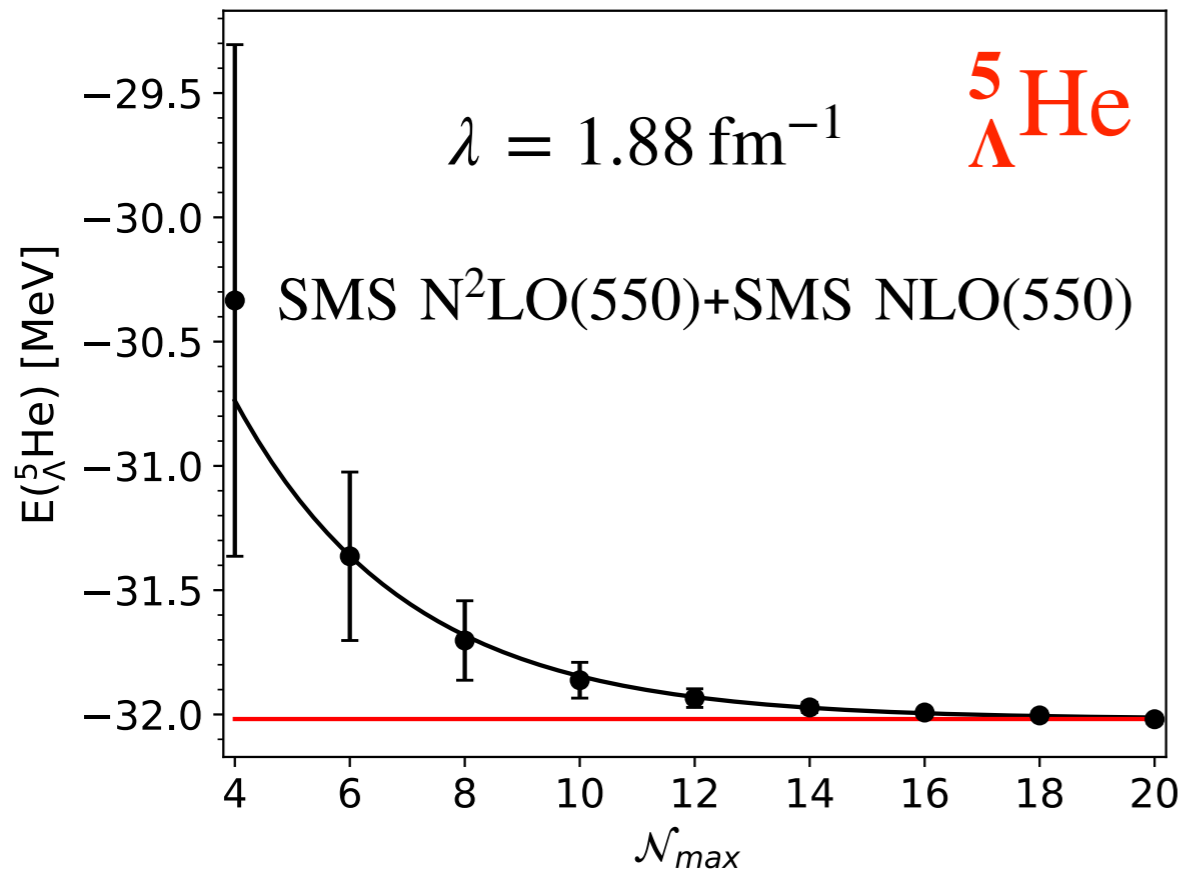
➔ missing 3N and YNN interactions

- 3NF is comparable to chiral 3NF
- YNN is larger than chiral YNN

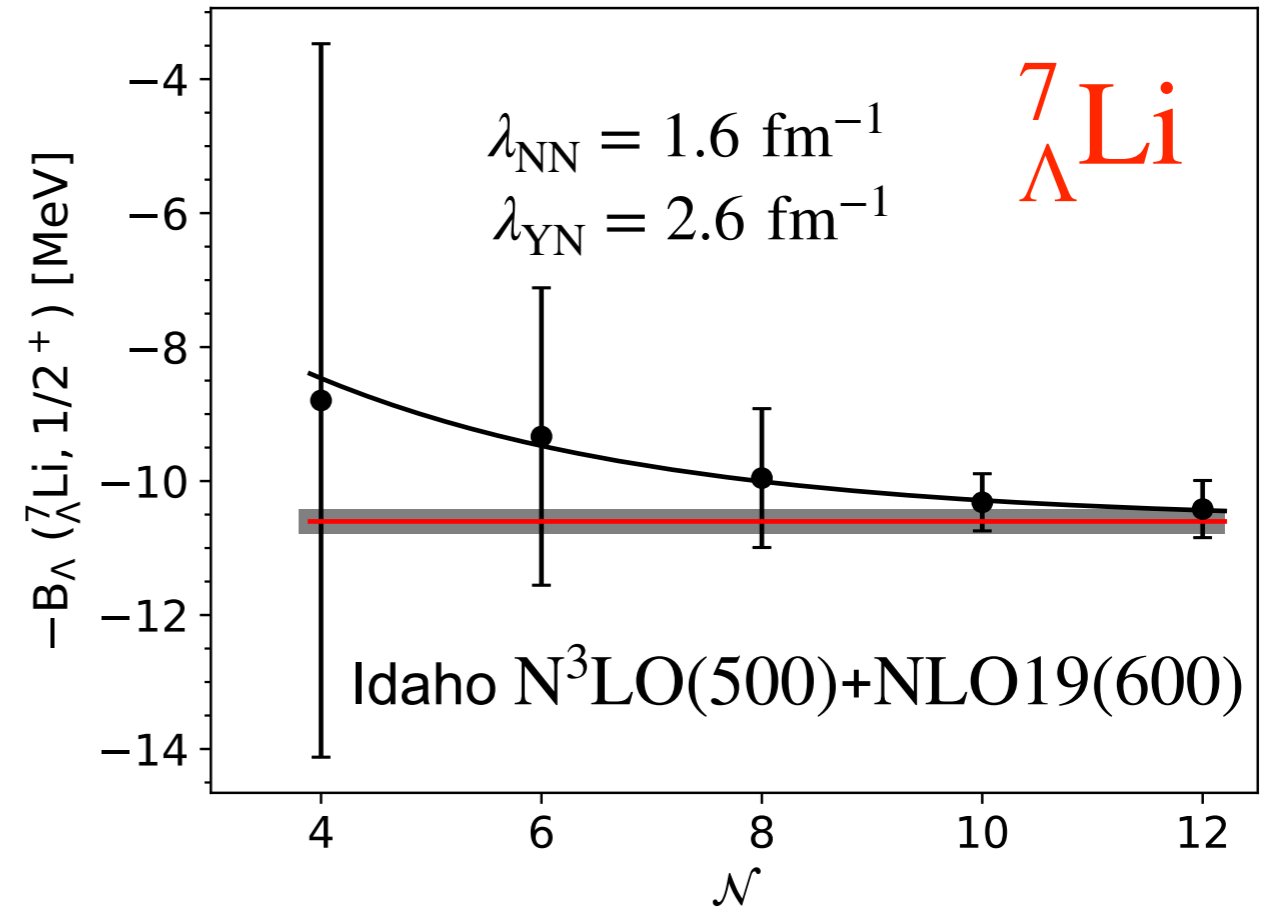


J-NCSM convergence

SRG evolution improves convergence



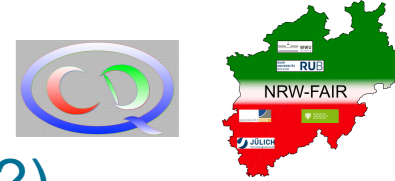
$$E({}^5_{\Lambda}\text{He}) = -32.018 \pm 0.001 \text{ MeV}$$



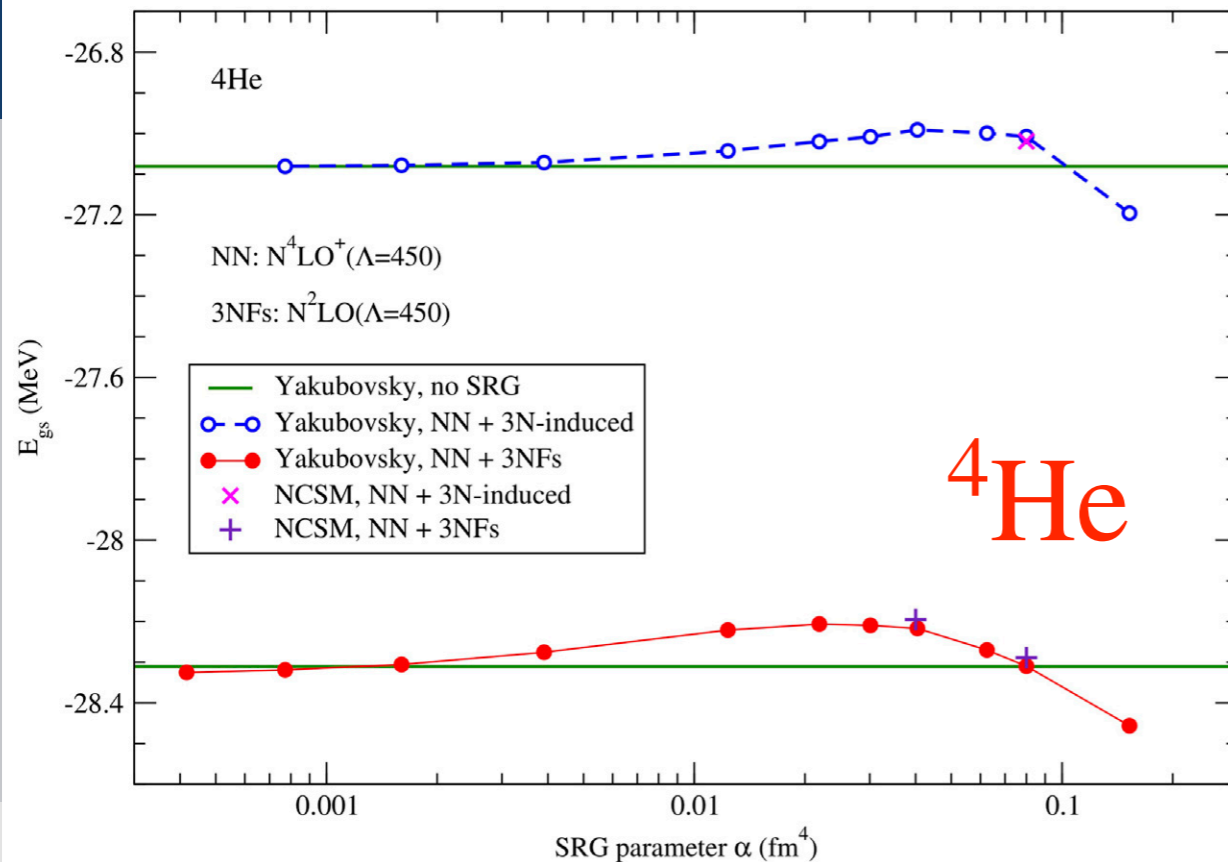
$$E_{\Lambda}({}^7_{\Lambda}\text{Li}) = 10.6 \pm 0.2 \text{ MeV}$$

- for light nuclei and hypernuclei, the numerical uncertainty is negligible.
- for p-shell nuclei/hypernuclei, the uncertainty is visible
- extrapolation of separation energy can reduce uncertainty of this quantity

SRG dependence of results

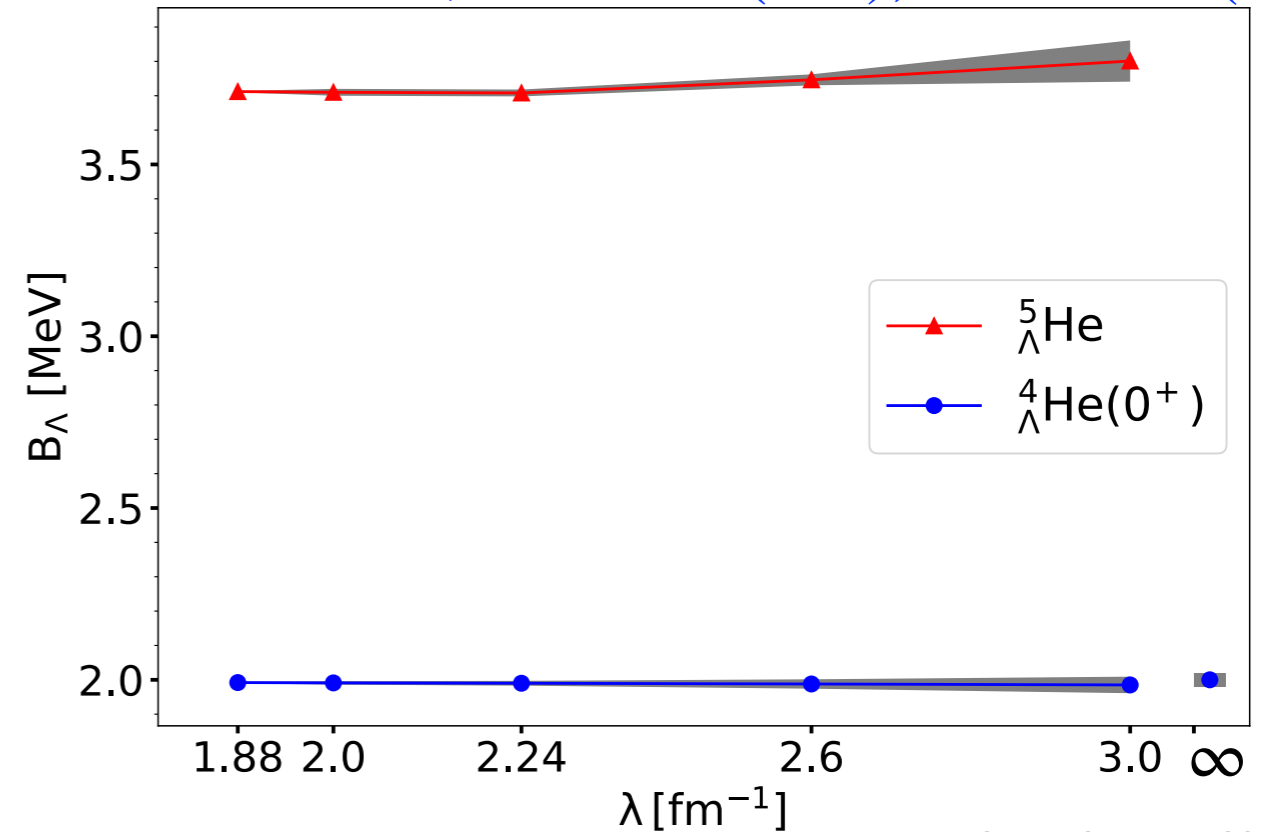


- SRG-induced 3N and YNN interactions
- ^4He binding energies varies by $\approx 100 - 200$ keV (relevant in the future?)
- separation energies are even less dependent (YNNN forces small)



(Maris, Le, Nogga, Roth, Vary (2023))

NN: $N^4\text{LO}^+$, 3N: $N^2\text{LO}(450)$; YN: $N^2\text{LO}(550)$



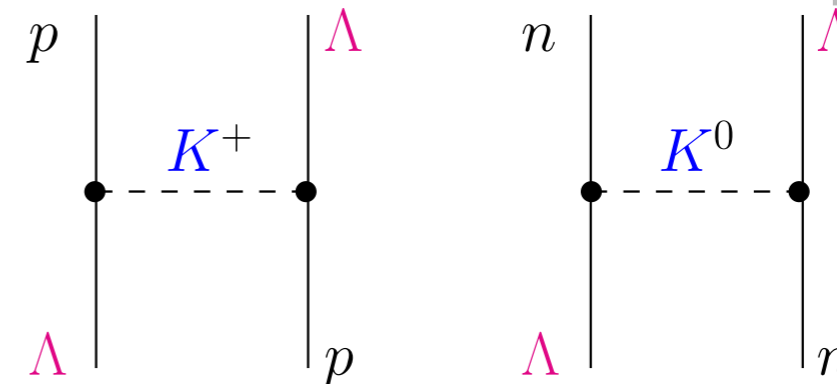
(Le (2023))

For **hypernuclei**, calculations based on SRG induced BB and 3B interactions are sufficiently accurate!

CSB contributions to ΛN interactions

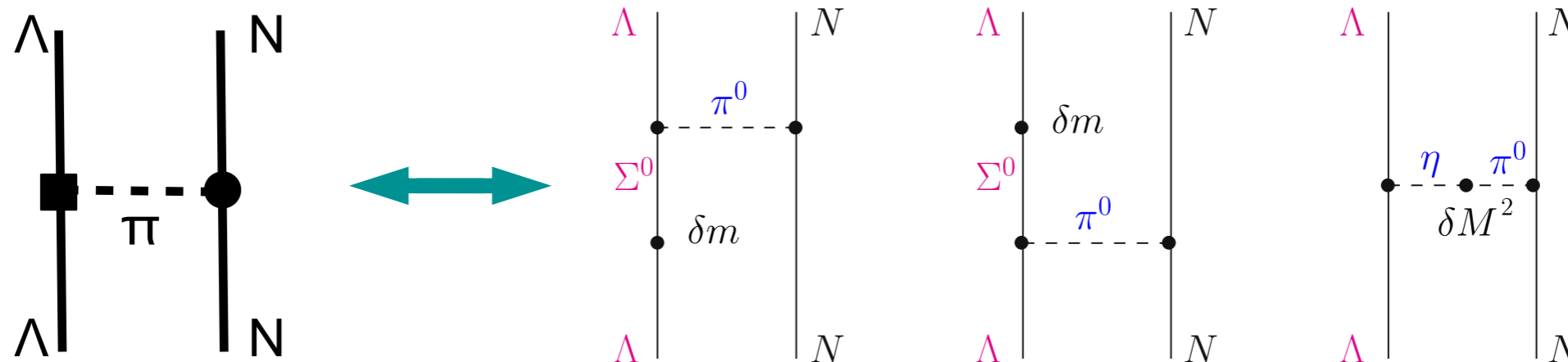


- formally leading contributions: Goldstone boson mass difference
 - very small due to the small relative difference of kaon masses



- subleading but most important
 - effective CSB $\Lambda\Lambda\pi$ coupling constant (Dalitz, van Hippel, 1964)

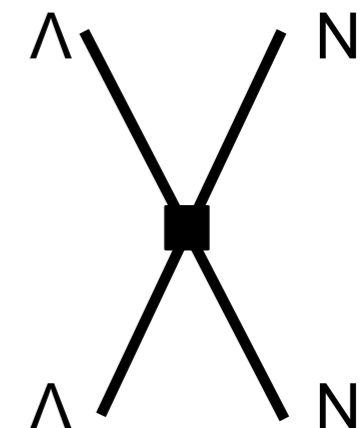
$$f_{\Lambda\Lambda\pi} = \left[-2 \frac{\langle \Sigma^0 | \delta m | \Lambda \rangle}{m_{\Sigma^0} - m_{\Lambda}} + \frac{\langle \pi^0 | \delta M^2 | \eta \rangle}{M_{\eta}^2 - M_{\pi^0}^2} \right] f_{\Lambda\Sigma\pi} \approx (-0.0297 - 0.0106) f_{\Lambda\Sigma\pi}$$



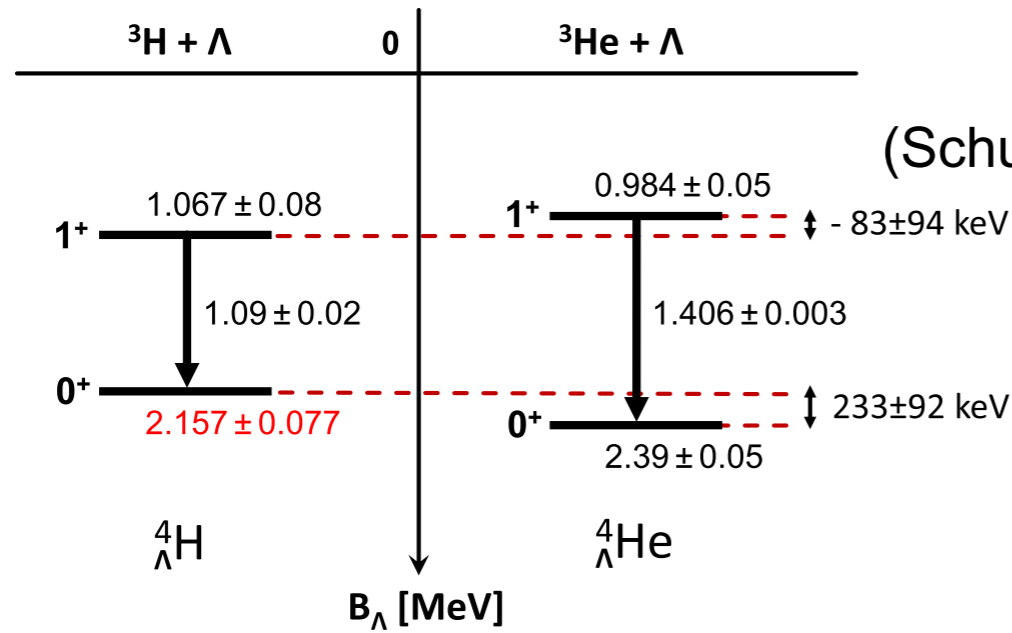
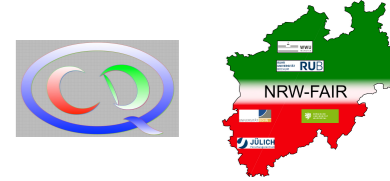
- so far less considered, but equally important
 - CSB contact interactions (for singlet and triplet)

Aim: use $A=4$ hypernuclei to determine the two unknown CSB LECs and predict Λn scattering

(so far: NLO13 and NLO19)



Fit of contact interactions



(Schulz et al., 2016; Yamamoto, 2015)

- Adjust the two CSB contact interactions to one main scenario (**CSB1**)

Λ	NLO13		NLO19	
	C_s^{CSB}	C_t^{CSB}	C_s^{CSB}	C_t^{CSB}
500	4.691×10^{-3}	-9.294×10^{-4}	5.590×10^{-3}	-9.505×10^{-4}
550	6.724×10^{-3}	-8.625×10^{-4}	6.863×10^{-3}	-1.260×10^{-3}
600	9.960×10^{-3}	-9.870×10^{-4}	9.217×10^{-3}	-1.305×10^{-3}
650	1.500×10^{-2}	-1.142×10^{-3}	1.240×10^{-2}	-1.395×10^{-3}

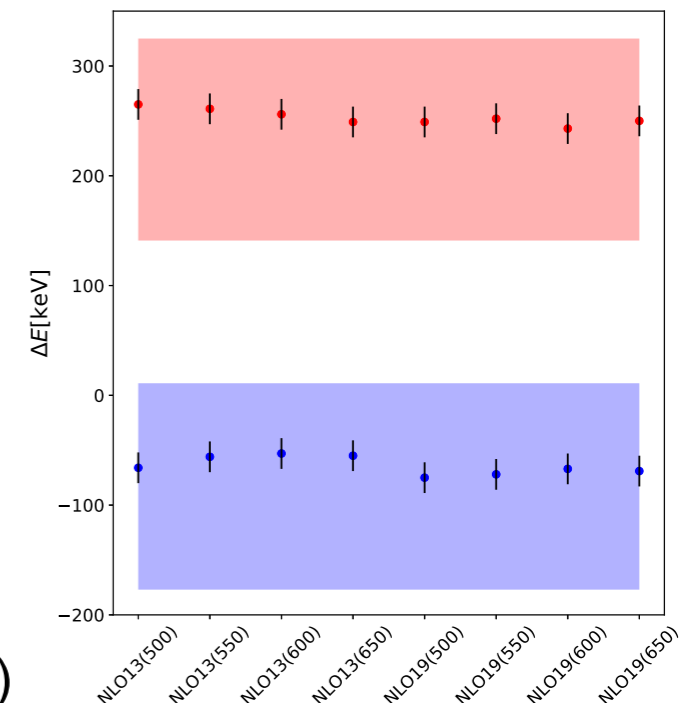
The values of the LECs are in 10^4 GeV^{-2}

- Size of LECs as expected by power counting

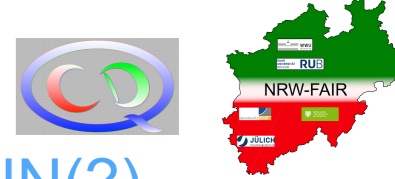
$$\frac{m_d - m_u}{m_u + m_d} \left(\frac{M_\pi}{\Lambda} \right)^2 C_{S,T} \approx 0.3 \cdot 0.04 \cdot 0.5 \cdot 10^4 \text{ GeV}^{-2} \propto 6 \cdot 10^{-3} \cdot 10^4 \text{ GeV}^{-2}$$

- Problem: large experimental uncertainty of experiment
- here only fit to central values to test theoretical uncertainties

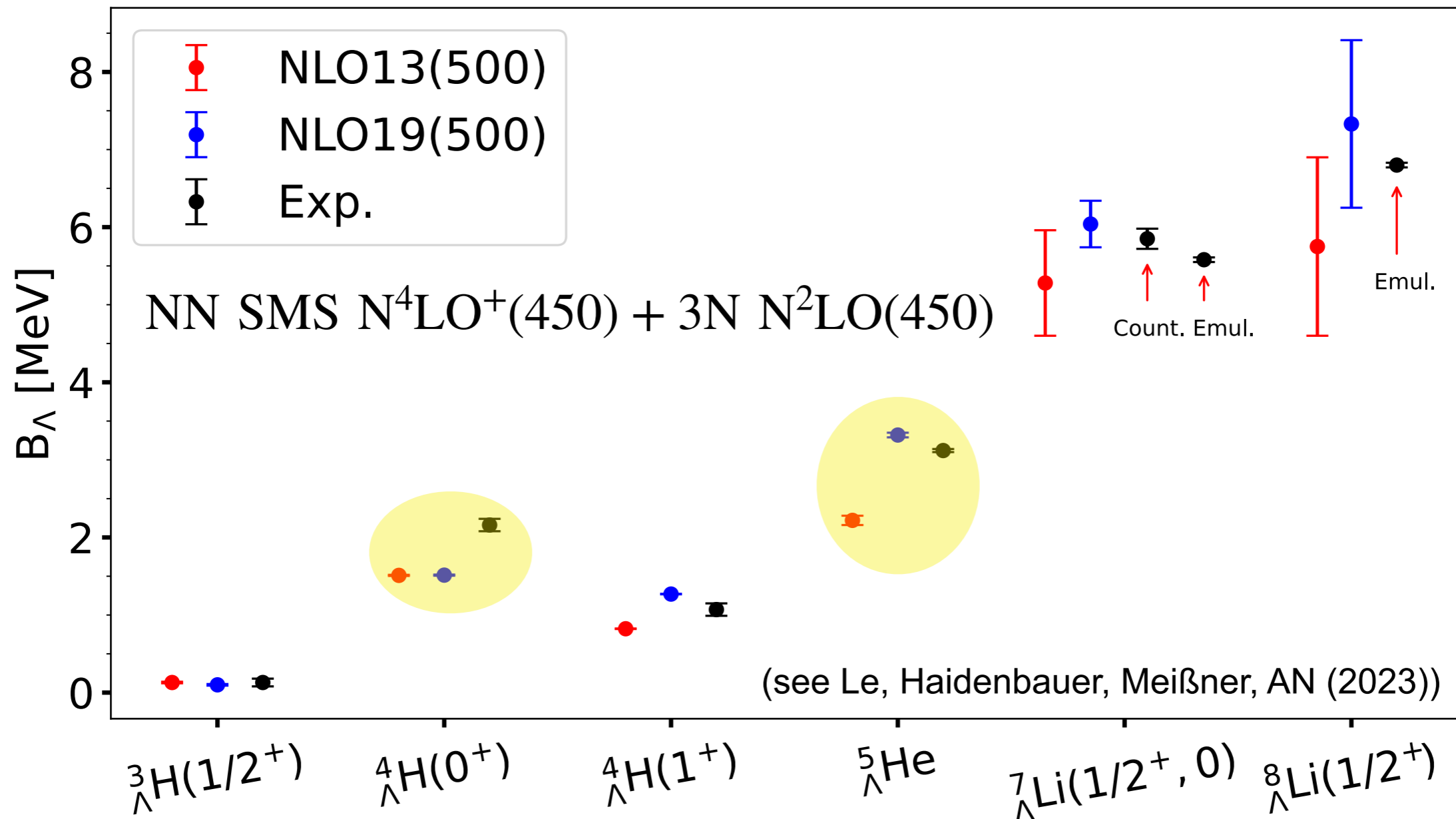
(see Haidenbauer, Meißner, AN (2021))



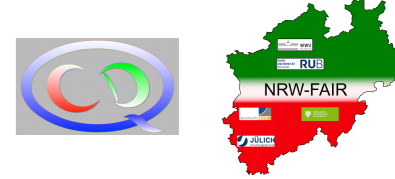
Application to $A = 7$ and 8



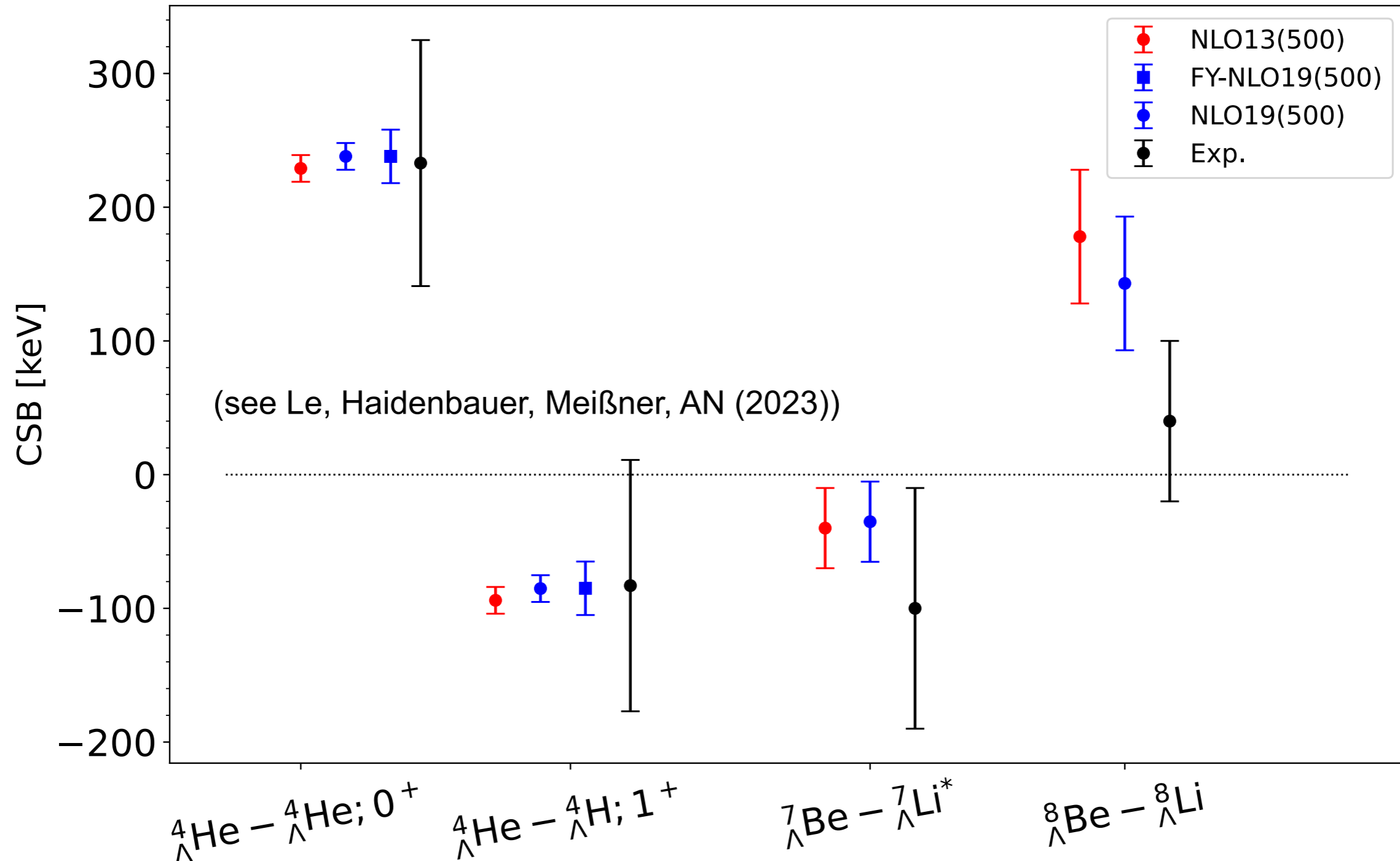
- YN interaction adjusted to the hypertriton — YNN is small
- based only on YN interactions: splitting for ${}^4_{\Lambda}\text{H}$ is not well reproduced — YNN(?)
- NLO19 gives better results for ${}^5_{\Lambda}\text{He}$ and heavier hypernuclei
— accidentally small YNN interaction?
- uncertainties are numerical — no estimate of chiral uncertainties yet



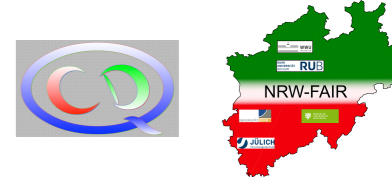
Application to $A = 7$ and 8



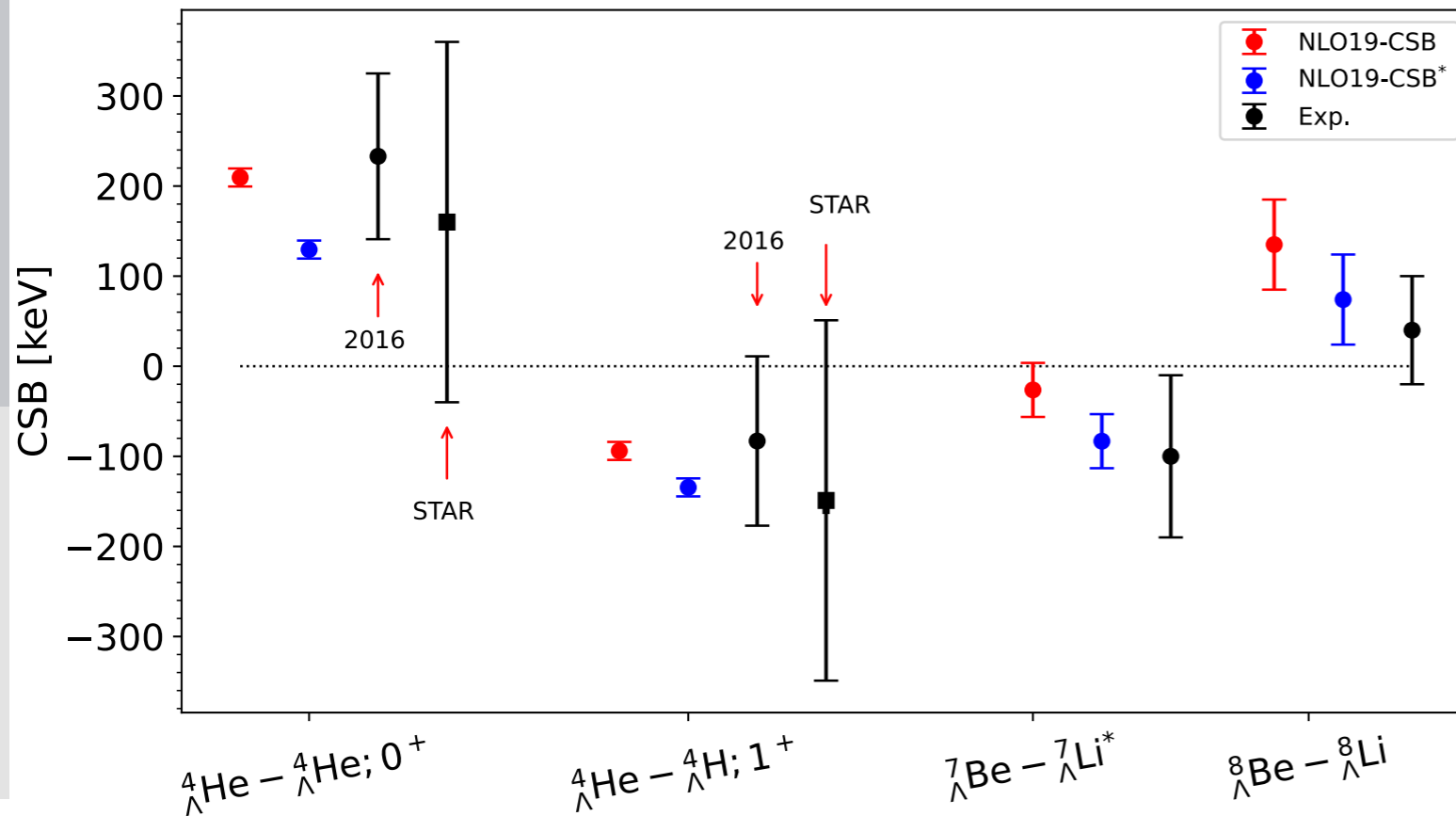
- CSB of singlet and triplet states interferes differently
- CSB still not fixed — experimental uncertainty is large
- scenario studied here is only **marginally consistent** with CSB in $A = 8$



New STAR data for $A = 4$ CSB



- fit to STAR data only
- only slight adjustment required
- improves description to p-shell CSB
- higher experimental accuracy is desirable
- good example of using hypernuclei to determine YN interactions



	NLO19(500)	CSB	CSB*
$a_s^{\Lambda p}$	-2.91	-2.65	-2.58
$a_s^{\Lambda n}$	-2.91	-3.20	-3.29
δa_s	0	0.55	0.71
$a_t^{\Lambda p}$	-1.42	-1.57	-1.52
$a_t^{\Lambda n}$	-1.41	-1.45	-1.49
δa_t	-0.01	-0.12	-0.03

(see Le, Haidenbauer, Meißner, AN (2023))

Uncertainty analysis to $A = 3$ to 5



Order N²LO requires combination of chiral NN, YN, 3N and **YNN** interaction

Need calculation of separation energies (use Faddeev, Yakubovsky eq. or J-NCSM)
and use **different orders** for uncertainty estimate.

Assuming a negligible numerical uncertainty and the following ansatz for the order by order convergence

$$X_K = X_{ref} \sum_{k=0}^K c_k Q^k \quad \text{where} \quad Q = M_{\pi}^{eff} / \Lambda_b \quad (X_{ref} \text{ LO, exp., max, ...})$$

a **Bayesian analysis** of the uncertainty is possible (see Melendez et al. 2017,2019)

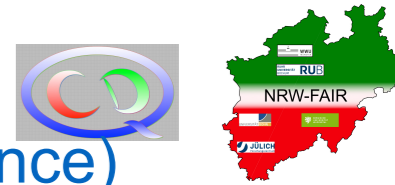
Extracting c_k for $k \leq K$ from calculations and assuming identical probability distributions for c_k for $k > K$ the uncertainty is given by the distribution of

$$\delta X_K = X_{ref} \sum_{k=K+1}^{\infty} c_k Q^k$$

Numerical uncertainties negligible (carefully checked!).

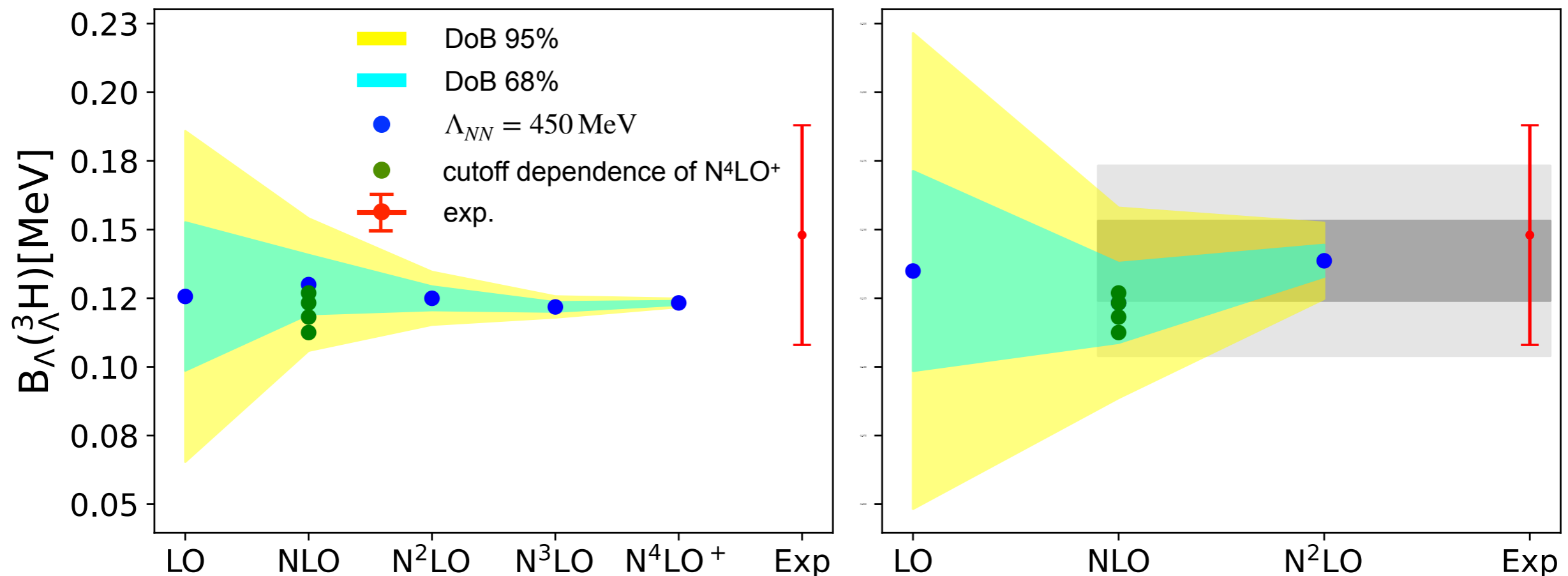
Uncertainty due to missing higher orders is most relevant!

Application to ${}^3_{\Lambda}\text{H}$

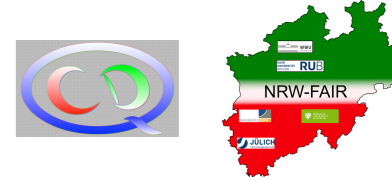


- Q , ν_0 and τ_0 are chosen using all available data (NN and YN convergence)
- uncertainties are extracted using c_k for NN or YN convergence
- use c_k of individual hypernuclei

➔ individual uncertainties for NN and YN convergence for each separation energy
consistent with experimental data
cutoff dependence always at least NLO (YNN missing!)

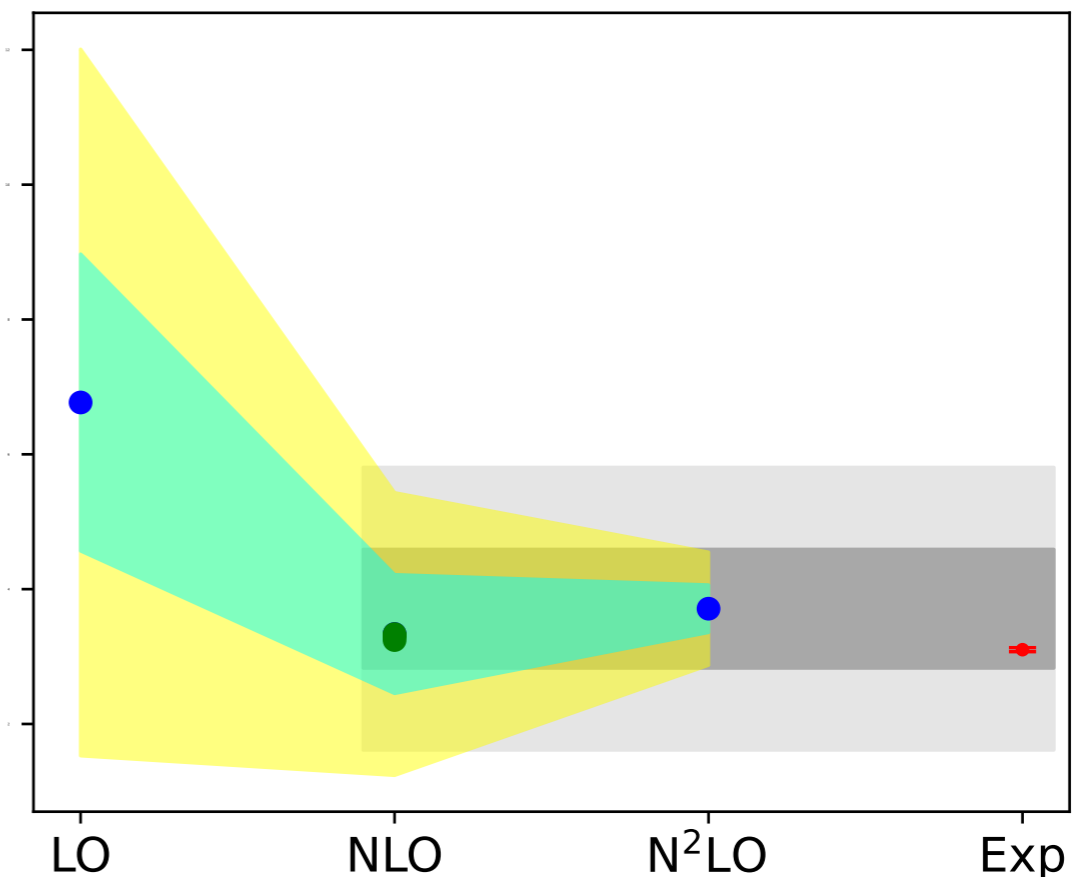
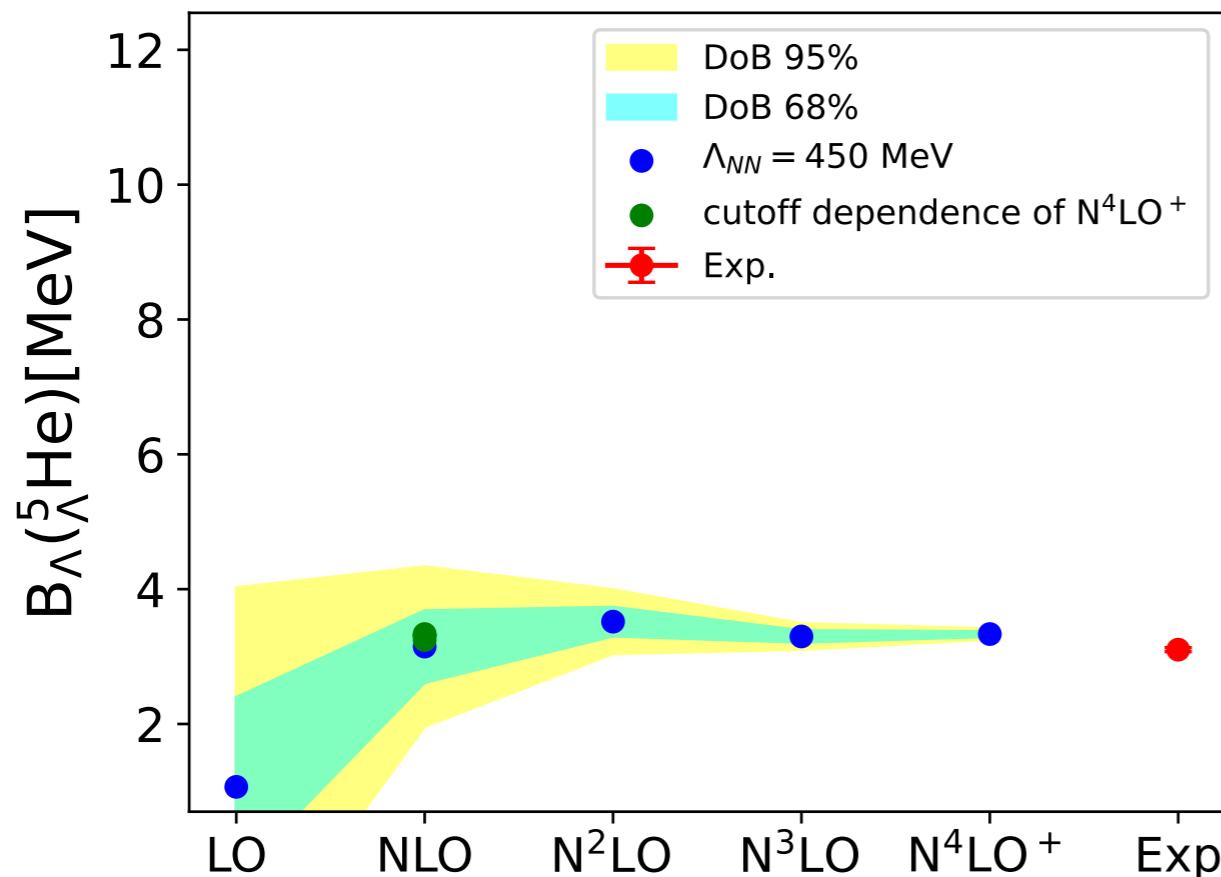


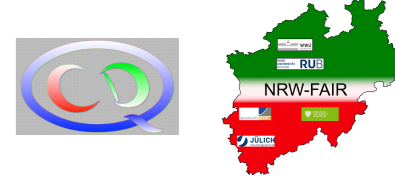
Application to ${}^5_{\Lambda}\text{He}$ and summary



- without YNN: sizable uncertainties at $A = 4$ and 5
- $A = 3$ sufficiently accurate
- NN/YN dependence small at least for $A = 3$

nucleus	$\Delta_{68}(NN)$	$\Delta_{68}(YN)$
${}^3_{\Lambda}\text{H}$	0.011	0.015
${}^4_{\Lambda}\text{He} (0^+)$	0.157	0.239
${}^4_{\Lambda}\text{He} (1^+)$	0.114	0.214
${}^5_{\Lambda}\text{He}$	0.529	0.881

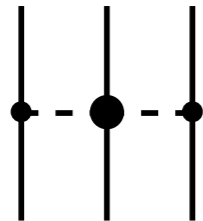




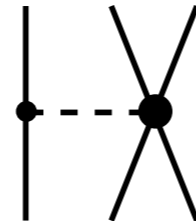
YNN (Λ NN) interactions

Leading 3BF with the usual topologies (see Petschauer et al., 2016 & 2017)

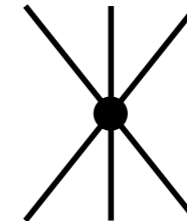
ChPT \longrightarrow all octet mesons contribute \longrightarrow only take π explicitly into account



2 LECs in Λ NN
(up to 10)



2 LECs in Λ NN
(up to 14)



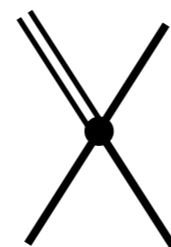
3 LECs in Λ NN
5 LECs in Σ NN + 1 Λ - Σ transition

only few data \longrightarrow need to keep the # of LECs small

Decuplet baryons (Σ^* ...) enhances YNN partly to NLO (see Petschauer et al., 2017)

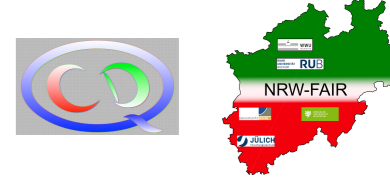
By decuplet saturation all LECs can be related to the following leading octet-decuplet transitions (Petschauer et al., 2020)

$$\propto C = \frac{3}{4} g_A$$

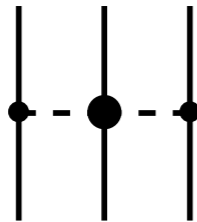


$$\propto G_1, G_2 \longrightarrow \text{reduction to 2 LECs}$$

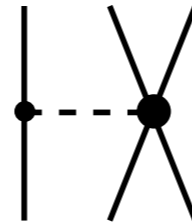
YNN (Λ NN) interactions



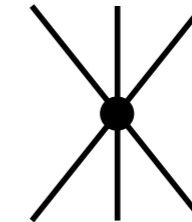
Decuplet saturation relates all LECs to G_1 and G_2



$$\propto C^2$$



$$\propto CG_1, CG_2$$



$$\propto (G_1)^2, (G_2)^2, G_1G_2$$

For Λ NN: $\propto C^2$

$$\propto C(G_1 + 3G_2)$$

$$\propto (G_1 + 3G_2)^2 \quad \mathbf{1 \text{ LEC}}$$

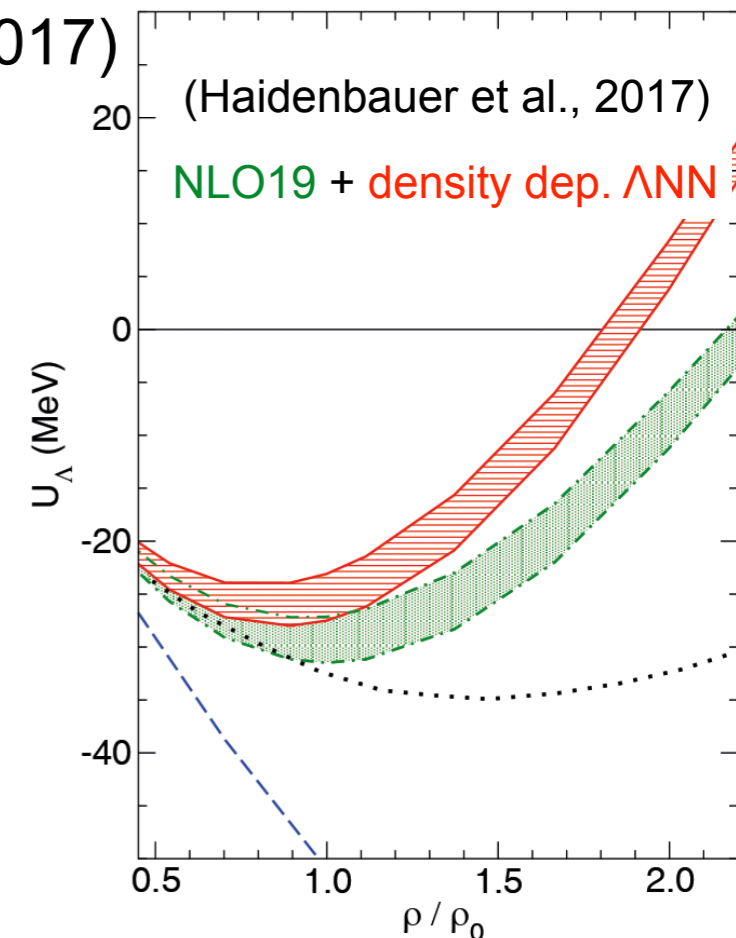
➔ density dependent BB interactions (Petschauer et al., 2017)

➔ application to nuclear matter (Haidenbauer et al., 2017)

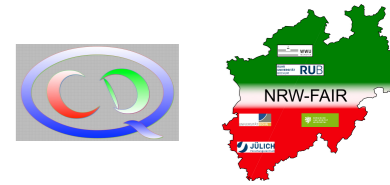
neutron stars (Logoteta et al., 2019)

- contribution on the single particle potentials can be large
- realistic results seem to require partly cancelations of 2π and 1π exchange (sign of $G_1 + 3G_2$!)

Results for hypertriton of Kohno et al., 2023/2024 small
but need correction ... benchmark ongoing



YNN (Λ NN) interactions



On the way to fit G_1/G_2 to ${}^4_\Lambda\text{He}/\text{H} \dots$

Recalculate 2π , 1π and contact terms of Λ NN using old **non-local** regularization

to compare to Kohno et al. (use fixed constant $G_1 = G_2 = \frac{1}{4f_\pi^2}$, $G_1 + 3G_2 = +\frac{1}{f_\pi^2}$)

→ Λ NN matrix elements ✓

Comparison of separation energies is ongoing (SMS $\text{N}^4\text{LO}^+(550)/\text{N}^2\text{LO} + \text{NLO19}$):

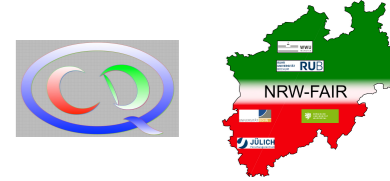
	<i>w/o</i> YNN	<i>w/</i> 2π	<i>w/</i> $2\pi/1\pi$	<i>w/</i> $2\pi/1\pi/\text{ct}$
${}^3_\Lambda\text{H}$ <i>w/o</i> Σ NN	0.080	0.151	0.215	0.208
${}^3_\Lambda\text{H}$		0.241	0.564	0.549
${}^4_\Lambda\text{He}(0^+)$	1.432	2.412		
${}^4_\Lambda\text{He}(1^+)$	1.164	2.623		
${}^5_\Lambda\text{He}$	3.174	7.139		

Large contribution to all light hypernuclei (larger than estimate!)

— consistent description requires larger cancelation of 2π and 1π part

— contact terms negligible for ${}^3_\Lambda\text{H}$

YNN (Λ NN) interactions



On the way to fit G_1/G_2 to ${}^4_{\Lambda}\text{He}/\text{H} \dots$

apply locally regularized YNN including subtractions

inclusion of ΣNN gives dependence on G_1/G_2 independently

Test results (SMS $N^4\text{LO}^+(550)/N^2\text{LO} + \text{SMS NLO}(550)$):

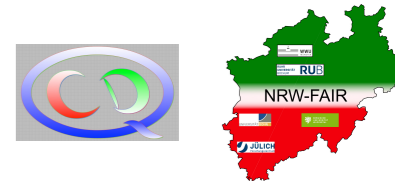
	<i>w/o</i> YNN	<i>w/</i> 2π	<i>w/</i> $2\pi/1\pi$	<i>w/</i> $2\pi/1\pi/\text{ct}$
${}^3_{\Lambda}\text{H}$ <i>w/o</i> subtr	0.107	0.149		
${}^3_{\Lambda}\text{H}$ only subtr		0.086		
${}^3_{\Lambda}\text{H}$ Λ NN compl		0.124		
${}^3_{\Lambda}\text{H}$		0.159	0.238	
${}^4_{\Lambda}\text{He}(0^+)$	1.969	2.333		
${}^4_{\Lambda}\text{He}(1^+)$	1.063	1.367		
${}^5_{\Lambda}\text{He}$	3.247	4.294		

SMS regularization leads to much more **natural results**.

Consistent regularization of NN/3N and YN/YNN forces?

Sensitivity of hypernuclear binding to G_1/G_2 ?

Conclusions & Outlook



- **YN (& YY) interactions not well understood**
 - *scarce YN data (almost no YY data)*
 - *more information necessary to solve "hyperon puzzle"*
- **Hypernuclei provide important constraints**
 - *CSB of ΛN scattering & ${}^4_{\Lambda}\text{He}$ / ${}^4_{\Lambda}\text{H}$*
 - *${}^3_{\Lambda}\text{H}$ is used to constrain the spin dependence*
 - *new experiments & analyses planned at J-PARC, MAMI, J-Lab, FAIR,...*
- **New SMS YN interactions**
 - *give an accurate description low energy YN data*
 - *order LO, NLO and N²LO allow uncertainty quantification*
 - *have a non-unique determination of contact interactions (data necessary)*
- **Chiral 3BF need to be included**
 - *NLO uncertainty is sizable in $A = 4$ and beyond*
 - *chiral 3BFs are now available*
 - *fitting to ${}^4_{\Lambda}\text{He}$ in progress*
 - *regularization affects size of individual contributions*