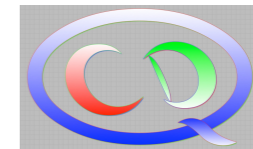


Application of SMS chiral interactions to light hypernuclei



Andreas Nogga, Forschungszentrum Jülich

ROCKSTAR Workshop, ECT*, Trento, Italy

- Motivation
- New chiral YN interactions
- Theoretical uncertainties of Λ -separation energies
- Estimates of 3BF contributions for ${}^3_{\Lambda}\text{H}$, ${}^4_{\Lambda}\text{H}$ / ${}^4_{\Lambda}\text{He}$ and ${}^5_{\Lambda}\text{He}$
- CSB of the YN interaction
- Conclusions & Outlook

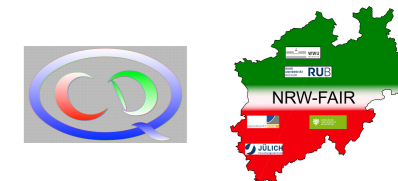
in collaboration with Hoai Le, Johann Haidenbauer and Ulf Meißner

Hoai Le et al. arXiv:2308.01756.

J. Haidenbauer et al. EPJ A 59, 63 (2023).

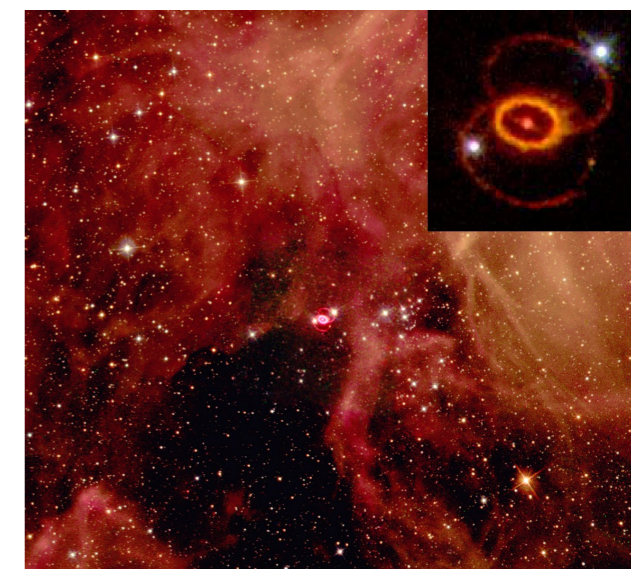
J. Haidenbauer et al. FBS 62, 105 (2021).

Hypernuclear interactions

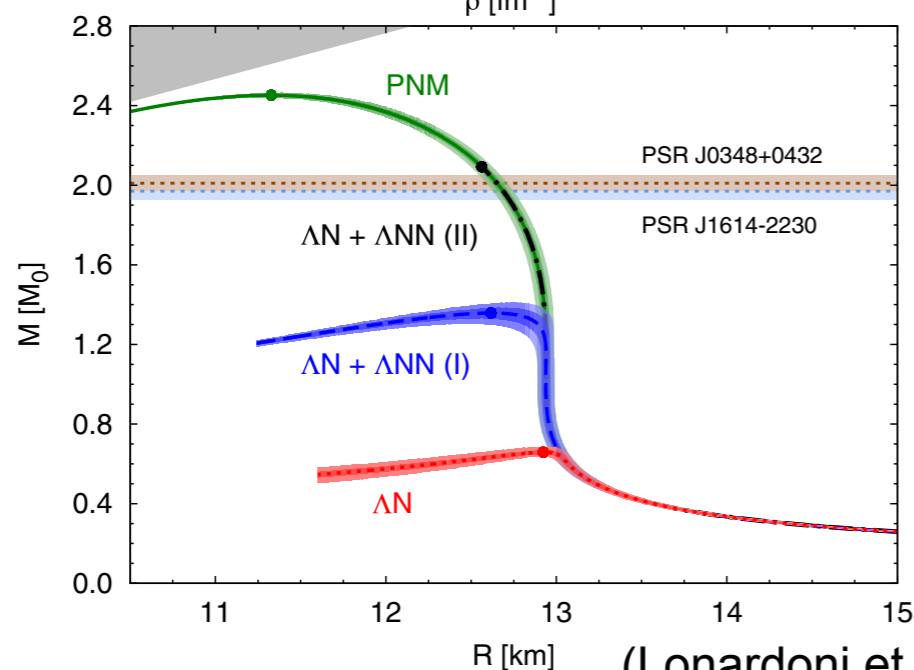
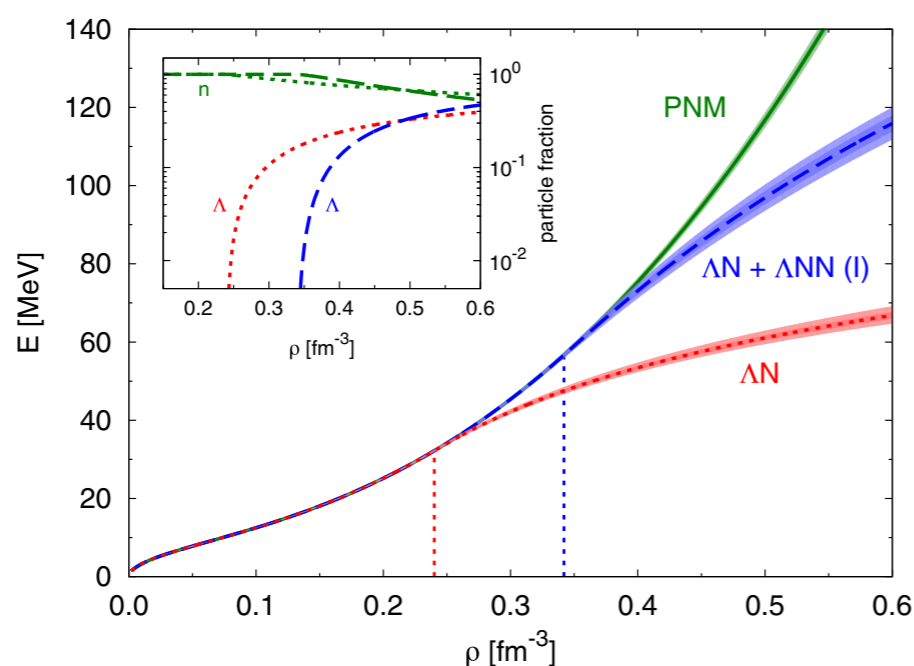


Why is understanding hypernuclear interactions interesting?

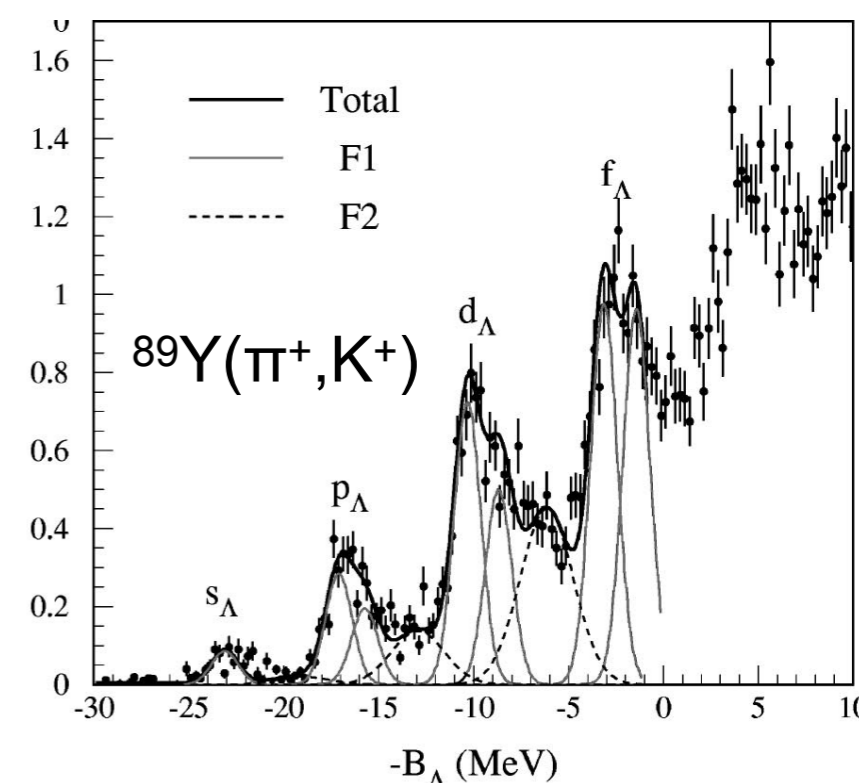
- *hyperon contribution to the EOS, neutron stars, supernovae*
- *"hyperon puzzle"*
- *Λ as probe to nuclear structure*
- *flavor dependence of baryon-baryon interactions*



(SN1987a, Wikipedia)



(Lonardoni et al. (2015))

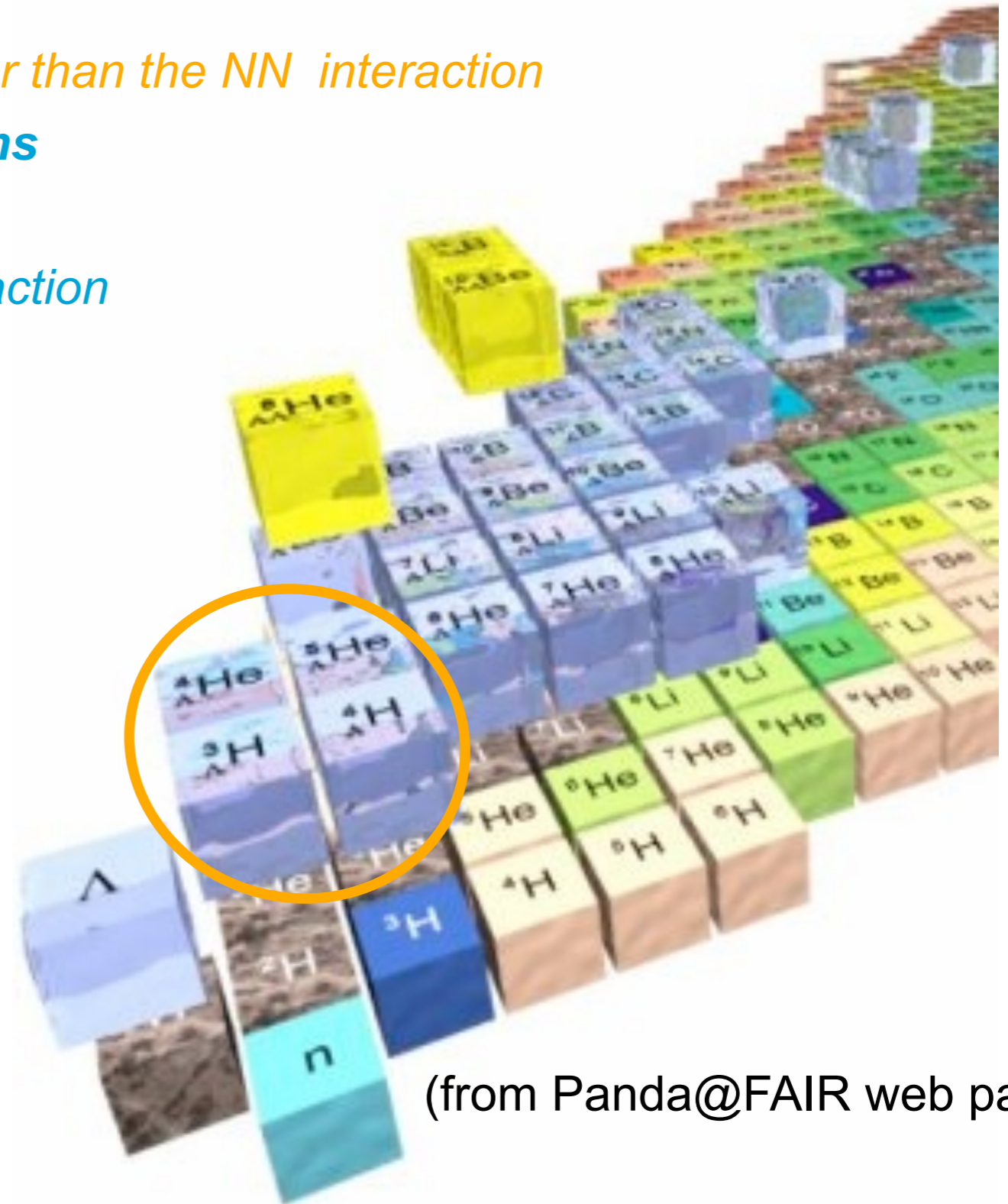


(Hotchi et al. (2001))

Hypernuclei

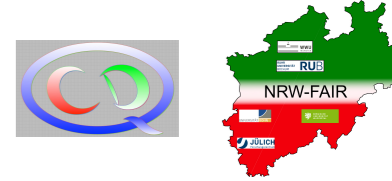
Only few YN data. Hypernuclear data provides additional constraints.

- ΛN interactions are generally weaker than the NN interaction
 - naively: core nucleus + hyperons
 - „separation energies“ are quite independent from NN(+3N) interaction
- no Pauli blocking of Λ in nuclei
 - good to study nuclear structure
 - even light hypernuclei exist in several spin states
- non-trivial constraints on the YN interaction even from lightest ones
- size of YNN interactions?
need to include Λ - Σ conversion!



(from Panda@FAIR web page)

Hypernuclei



- new data (J-PARC, Star, Alice, ...)
- world averages compiled by Mainz group (Eckert et al.)

hypernuclei.kph.uni-mainz.de

CHART OF HYPERNUCLIDES – Hypernuclear Structure and Decay Data

${}^3_{\Lambda}\text{H}$ Hydrogen

- Non-strange core: ${}^2\text{H}$
- mass: $m_{\text{GS}} = 1875.613 \text{ MeV}/c^2$
- mean life time: stable
- ground state spin/parity: 1^+

- Hyperon Content: Λ
- mass: $m_{\text{GS}} = 1115.683 \text{ MeV}/c^2$
- mean life time: $\tau = 263.1 \text{ ps}$
- spin/parity: $\frac{1}{2}^+$

Chart Legend - available data

- - less than 6 values
- - less than 20 values
- - at least 20 values

${}^3_{\Lambda}\text{H}$: Λ binding energy

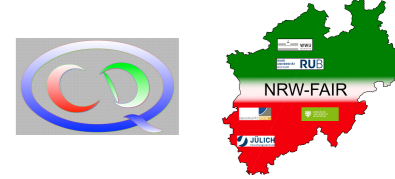
our value: $0.164 \pm 0.043 \text{ MeV}$

${}^3_{\Lambda}\text{H}$

Ground State: Λ Binding Energy our value: $0.164 \pm 0.043 \text{ MeV}$

B_{Λ} [MeV]	Weight	$\chi^2, \Sigma = 4.69$	Author	Year	Method	Comment	More...
<input checked="" type="checkbox"/> $0.102 \pm 0.063 \text{ (stat.)} \pm 0.067 \text{ (syst.)}$	0.21	0.46	ALICE	2023	Heavy Ion Coll.	-	Info Ref.
<input checked="" type="checkbox"/> $0.406 \pm 0.120 \text{ (stat.)} \pm 0.110 \text{ (syst.)}$	0.07	2.21	J. Adam	2020	Heavy Ion Coll.	combined	Info Ref.
<input type="checkbox"/> $0.346 \pm 0.130 \text{ (stat.)} \pm 0.121 \text{ (syst.)}$	0	0	J. Adam	2020	Heavy Ion Coll.	matter	Info Ref.
<input type="checkbox"/> $0.696 \pm 0.280 \text{ (stat.)} \pm 0.261 \text{ (syst.)}$	0	0	J. Adam	2020	Heavy Ion Coll.	antimatter	Info Ref.
<input type="checkbox"/> $-0.004 \pm 0.400 \text{ (stat.)}$	0	0	S. Acharya	2019	Heavy Ion Coll.	from lifetime measurem...	Info Ref.
<input type="checkbox"/> $0.296 \pm 1.000 \text{ (stat.)} \pm 3.000 \text{ (syst.)}$	0	0	J. Adam	2016	Heavy Ion Coll.	from lifetime measurem...	Info Ref.
<input type="checkbox"/> $-6.104 \pm 1.200 \text{ (stat.)}$	0	0	C. Rappold	2013	Heavy Ion Coll.	from lifetime measurem...	Info Ref.
<input type="checkbox"/> $2.296 \pm 1.000 \text{ (stat.)} \pm 2.000 \text{ (syst.)}$	0	0	B. I. Abelev	2010	Heavy Ion Coll.	from lifetime measurem...	Info Ref.
<input type="checkbox"/> $0.296 \pm 1.000 \text{ (stat.)} \pm 2.000 \text{ (syst.)}$	0	0	B. I. Abelev	2010	Heavy Ion Coll.	from lifetime measurem...	Info Ref.
<input checked="" type="checkbox"/> $0.130 \pm 0.050 \text{ (stat.)} \pm 0.040 \text{ (syst.)}$	0.44	0.29	M. Juric	1973	Emulsion	combined	Info Ref.
<input type="checkbox"/> $0.150 \pm 0.080 \text{ (stat.)} \pm 0.040 \text{ (syst.)}$	0	0	M. Juric	1973	Emulsion	combined	Info Ref.

Hypernuclear interactions



We will use here chiral interactions, start with a brief summary of other approaches to the ΛN (YN) interaction

Long history of ΛN interaction models

- **early models** (Downs, Iddings, Brown, Dalitz, before 1970)
- **Nijmegen group** (Nijm D, Nijm F, SC89, SC97 and ESC(ESC16), 1973-2016)
- **Jülich model** (Jülich 1994, Jülich 2004)
- **RGM model of Fujiwara, fss2** (1995, 2002)

models have successfully used to understand binding mechanism

important role of $\Lambda - \Sigma$ conversion

EFT based approaches

pionless (=Goldstone boson less) EFT

- **application to ${}^3_{\Lambda}\text{H}$ and Λnn** (Hammer 2002, Hildenbrand et al. 2019,2020)
- **application to $A = 3 - 5$ hypernuclei** (Contessi et al. PRL 2018)

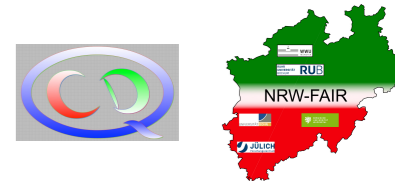
interactions given by contact interactions, usually in leading order

only Λ explicitly considered

EFT requires **3BFs** in leading order (three additional parameters)

expansion parameter

Chiral NN & YN interactions



EFT based approaches (cont')

	BB force	3B force	4B force	
LO		—	—	5 (+1) NN/YN (YY) short range parameters
NLO		—	—	23(+5) NN/YN (YY) short range parameters
N ² LO			—	no additional contact terms in NN/YN (YY)

Chiral EFT implements **chiral symmetry of QCD** (adapted from Epelbaum, 2008)

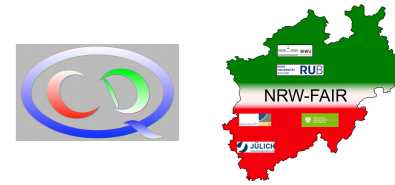
- perturbative expansion for the interaction
- non-perturbative solution of Schrödinger eq.
- symmetries constrain exchanges of Goldstone bosons
- relations of two- and three- and more-baryon interactions
- breakdown scale $\approx 600 - 700 \text{ MeV}$

Retain flexibility to adjust to data due to counter terms

Regulator required — cutoff/different orders often used to estimate uncertainty

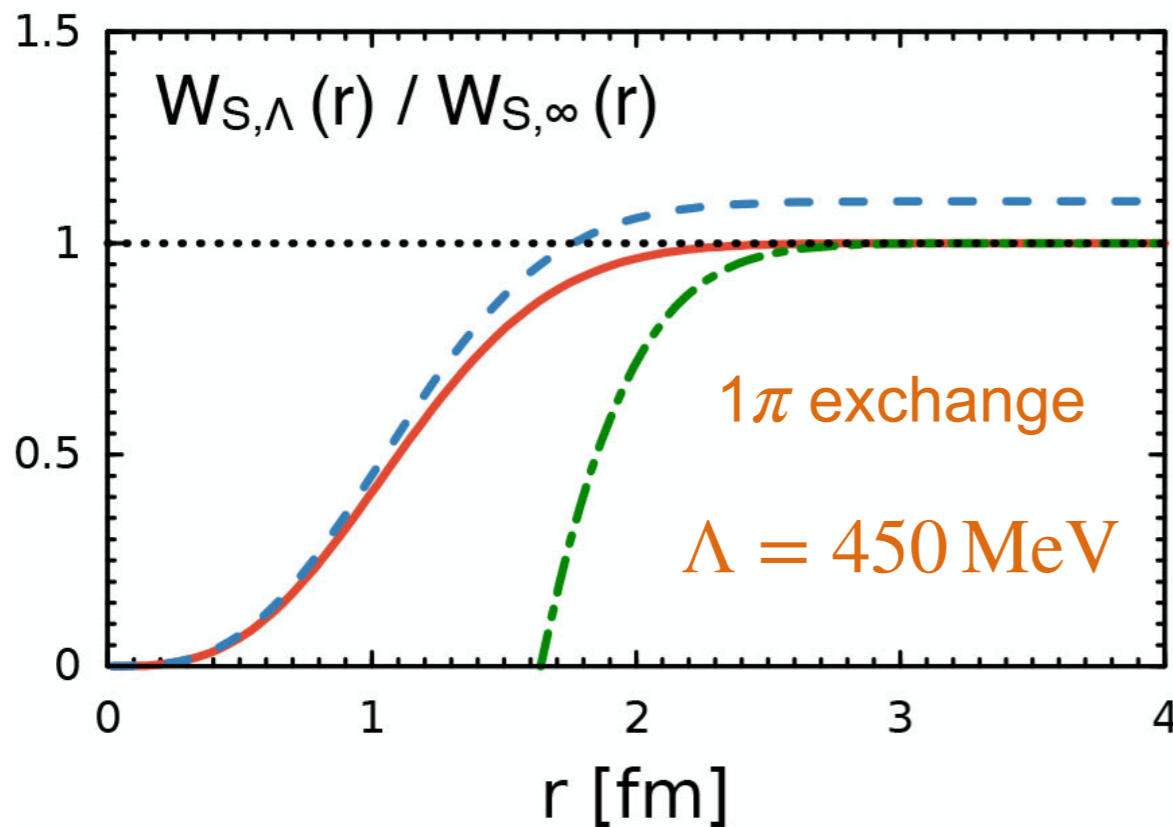
Λ - Σ conversion is explicitly included (size of 3BFs expected to be N²LO)

Choice of regulator

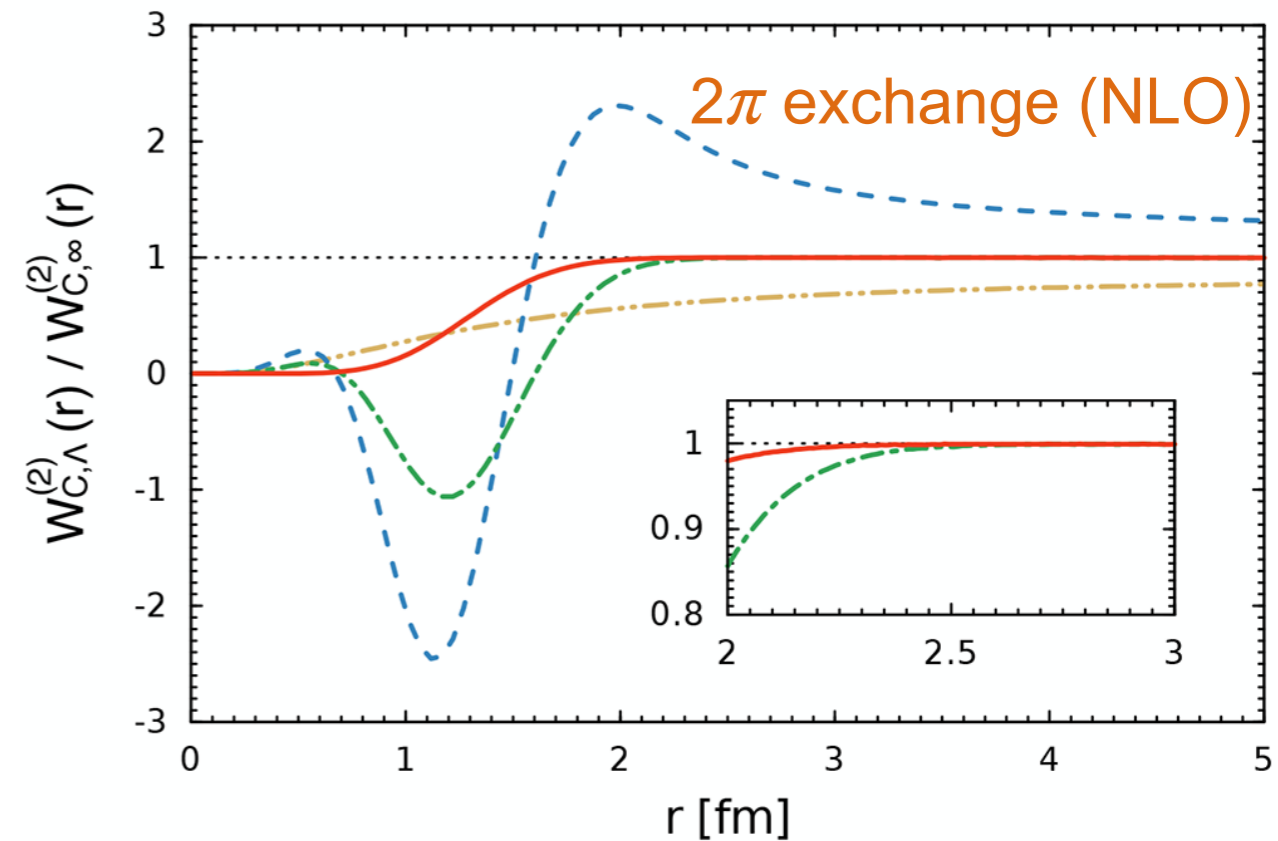


- trad. regularized (Entem et al. 2005, Epelbaum 2001, Ordonez et al. 1994)
- spectral function (SFR) regularization (Epelbaum 2005)
- semilocal coordinate-space (SCS) regularization (Epelbaum et al. 2015)
- semilocal momentum-space (SMS) regularization (Reinert et al. 2018)

$$W_{C,\Lambda}^{(2)}(r) = \frac{1}{2\pi^2 r} \int_{2M_\pi}^{\infty} d\mu \mu e^{-\mu r} \eta_C^{(2)}(\mu) e^{-\frac{q^2 + \mu^2}{2\Lambda^2}} + \text{subtraction}$$



(Reinert et al. 2018)



(Reinert et al. 2018)

NLO13/NLO19 YN based on trad. reg.

N²LO based on SMS (incl. subtractions)

SMS N²LO interaction

In order to check uncertainties realization in

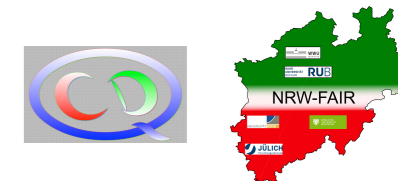
LO(700) NLO(500), **NLO(550)**, NLO(600) and N²LO(500), **N²LO(550)**, N²LO(600)



Details on fitting procedure using partly flavor-SU(3):

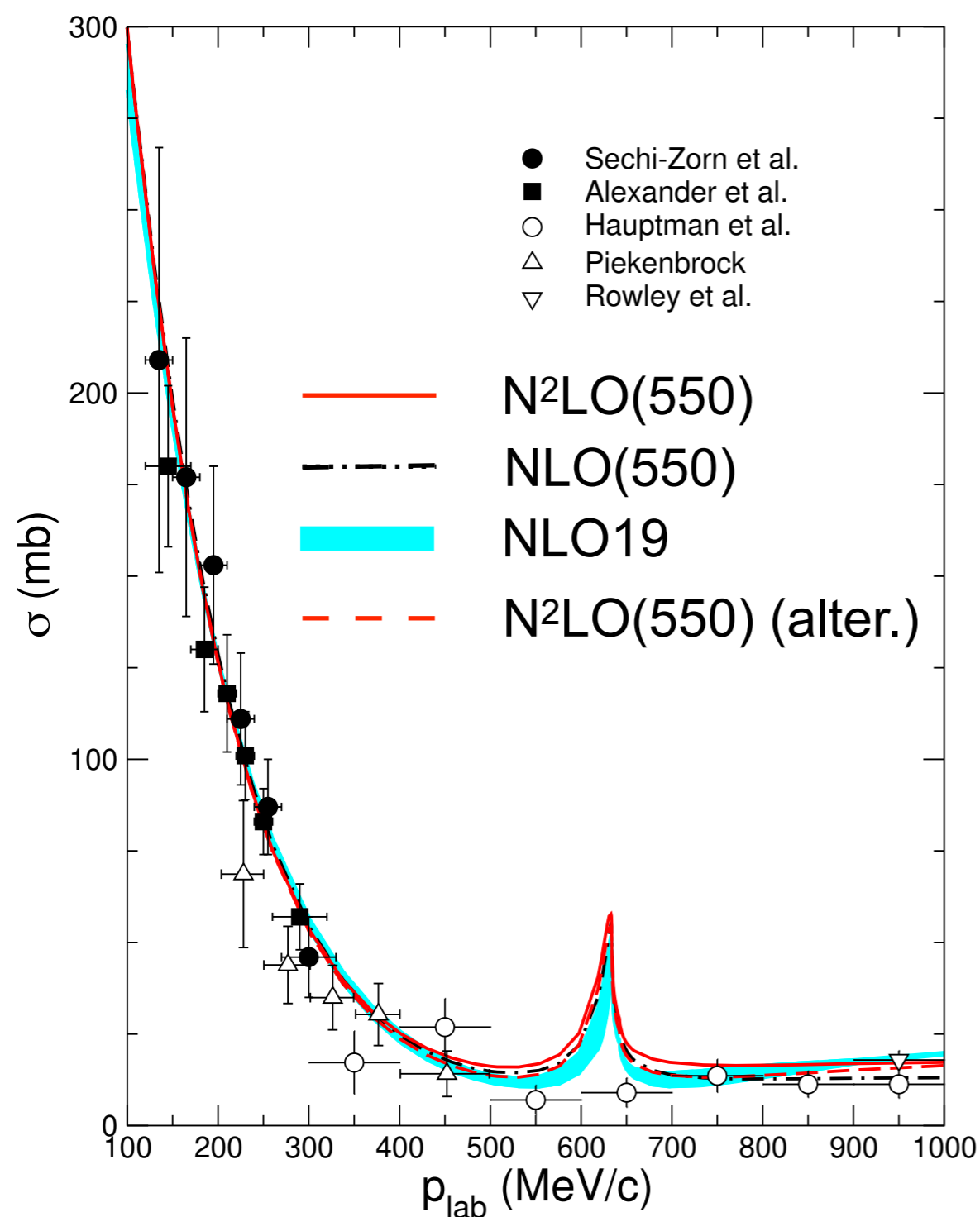
- flavor SU(3) is broken by using physical meson and baryon masses!
- retain only the 2π exchange from the 2 Goldstone boson exchanges
→ should be absorbed in SU(3)-breaking counter terms
- use **36 data** at low energy to determine **s-wave counter** terms
- hypertriton is required to be bound (binding energy **roughly** correct)
→ $a_s(\Lambda p) = -2.8$ fm in NLO and N²LO
- include **SU(3) breaking in LO** counter terms (necessary to avoid bound states in YN)
- assume **SU(3) symmetry for p-waves** counter terms in NLO
values for p-wave counter terms of NN
- fit to differential cross section in N²LO
→ two versions for N²LO(550) differ for differential cross sections
none is clearly preferred

SMS NLO/N²LO interaction



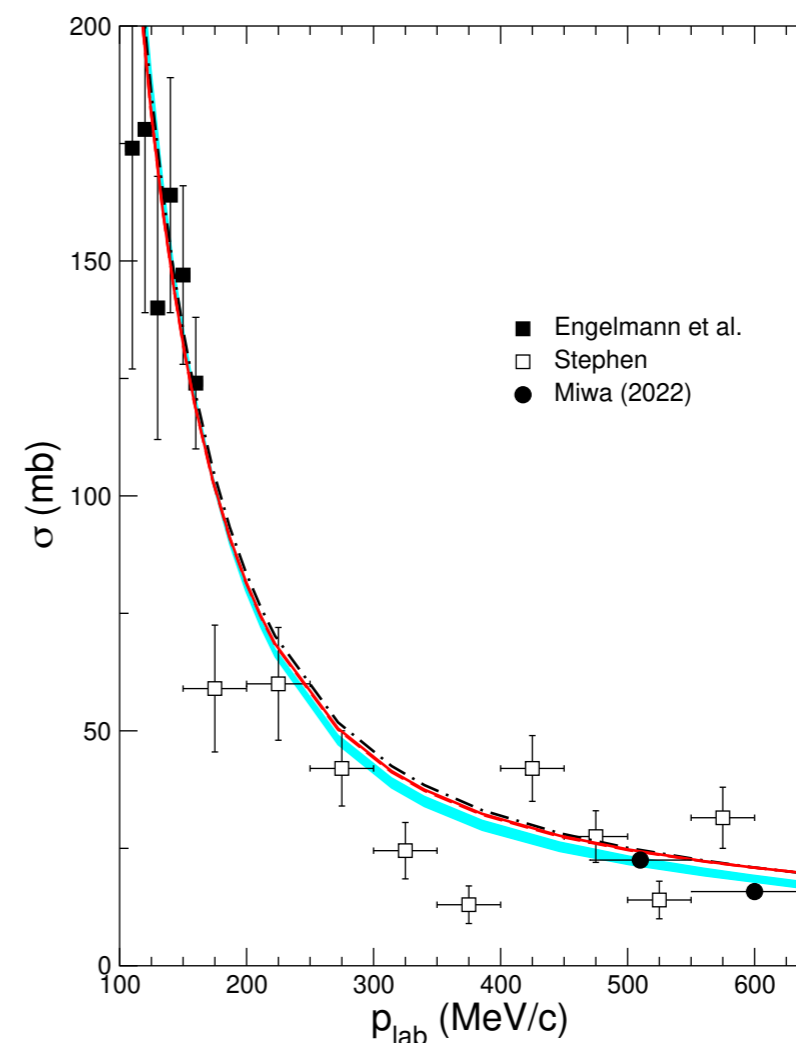
Selected results (show $\Lambda = 550$ MeV, others are very similar in quality)

$\Lambda p \rightarrow \Lambda p$

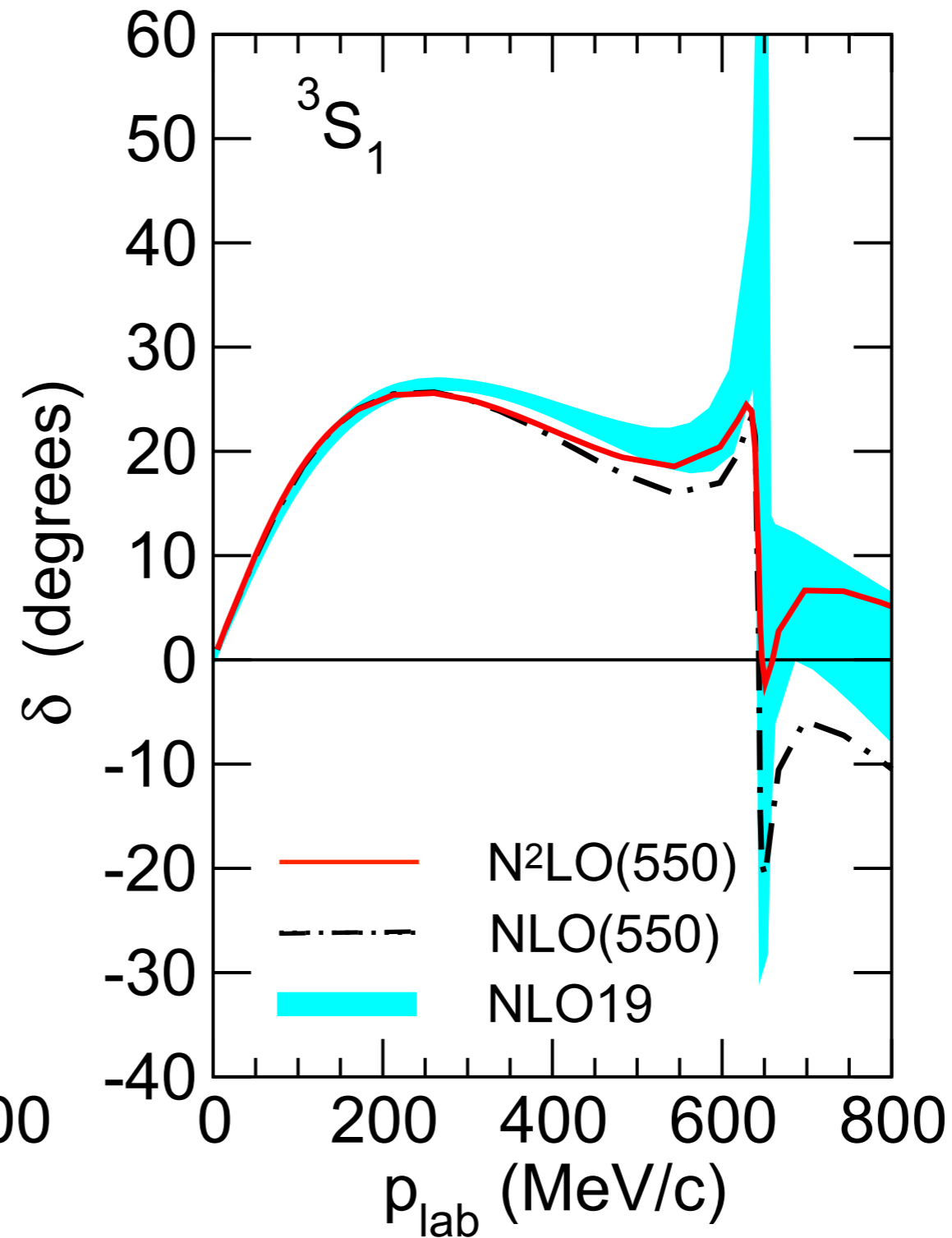
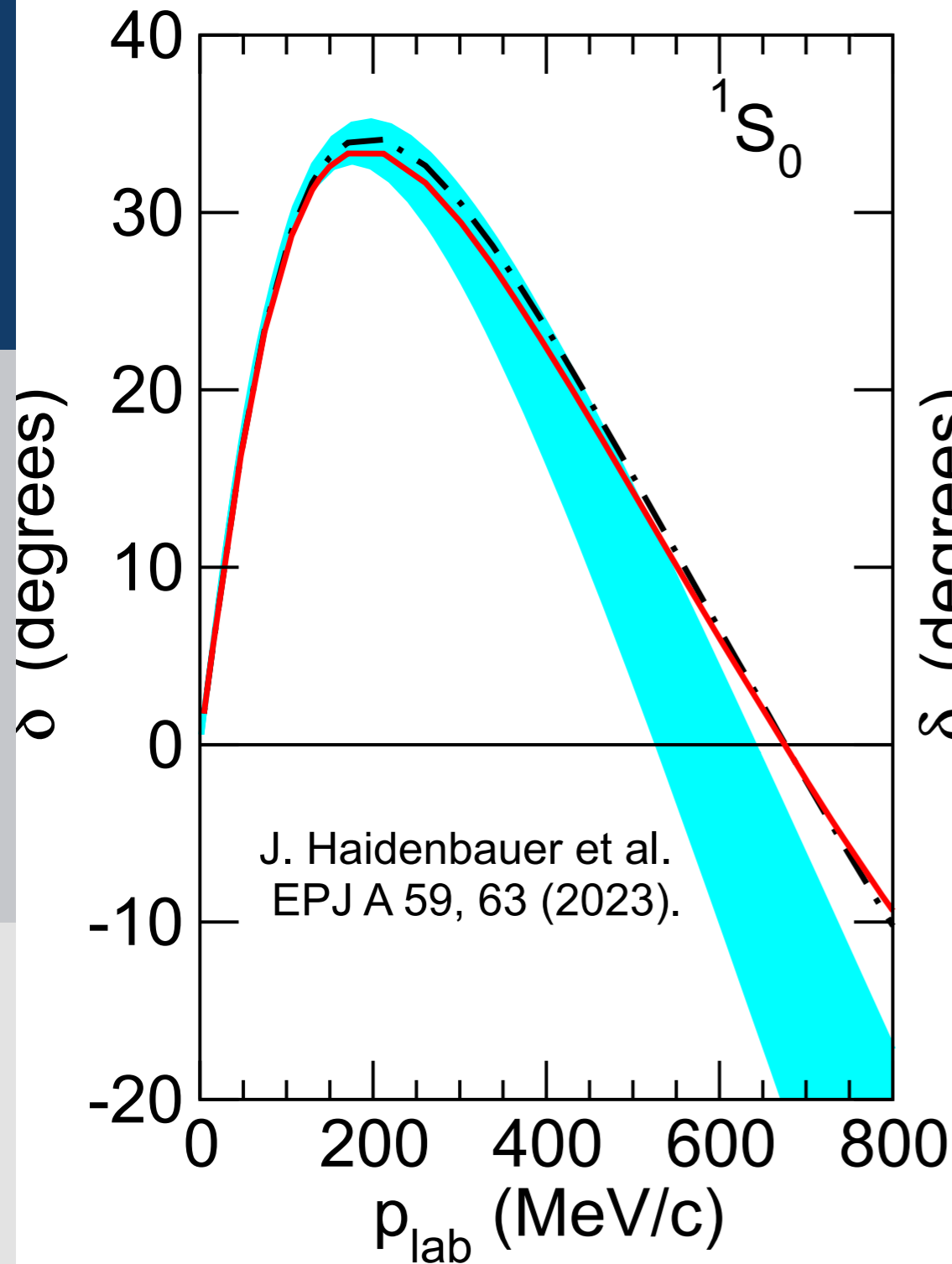
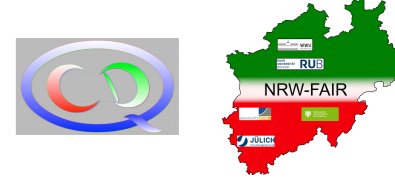


- most relevant cross sections very similar in NLO and N²LO
- similar to NLO19
- alternative fit (see later)

$\Sigma^- p \rightarrow \Lambda n$



Selected phase shifts



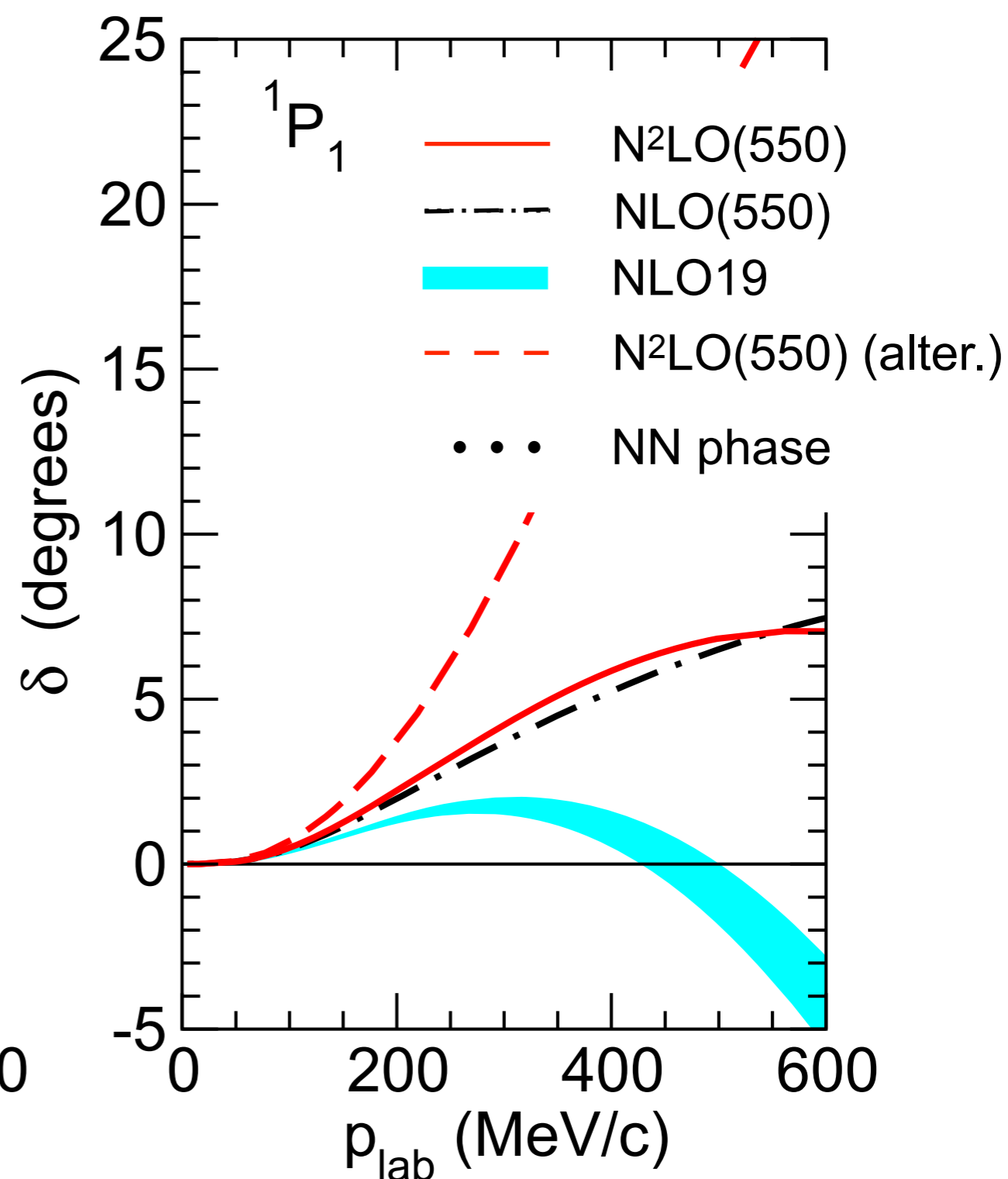
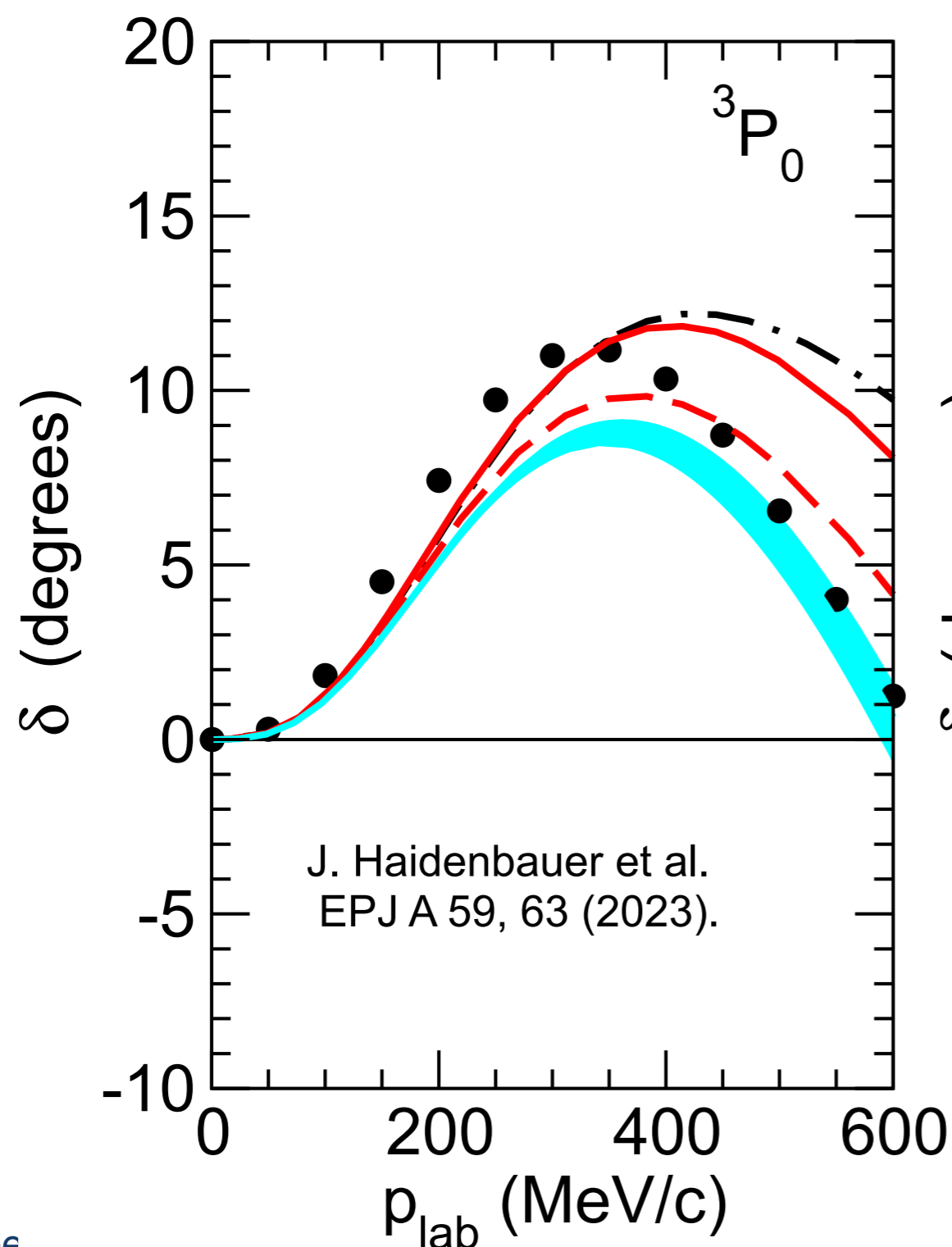
- s-phase shifts very similar in NLO and N²LO
- similar to NLO19 although different tail at high momenta

SMS NLO/N²LO interaction

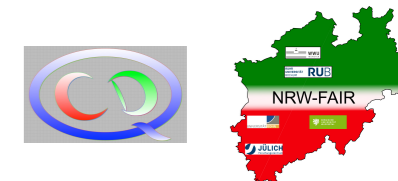


$\Sigma^+ - p$ scattering can be strongly related to NN in $^3P_0, ^3P_1, ^3P_2$ (at least in NLO)

In 1P_1 , in NLO, setting the counter term to zero is OK.

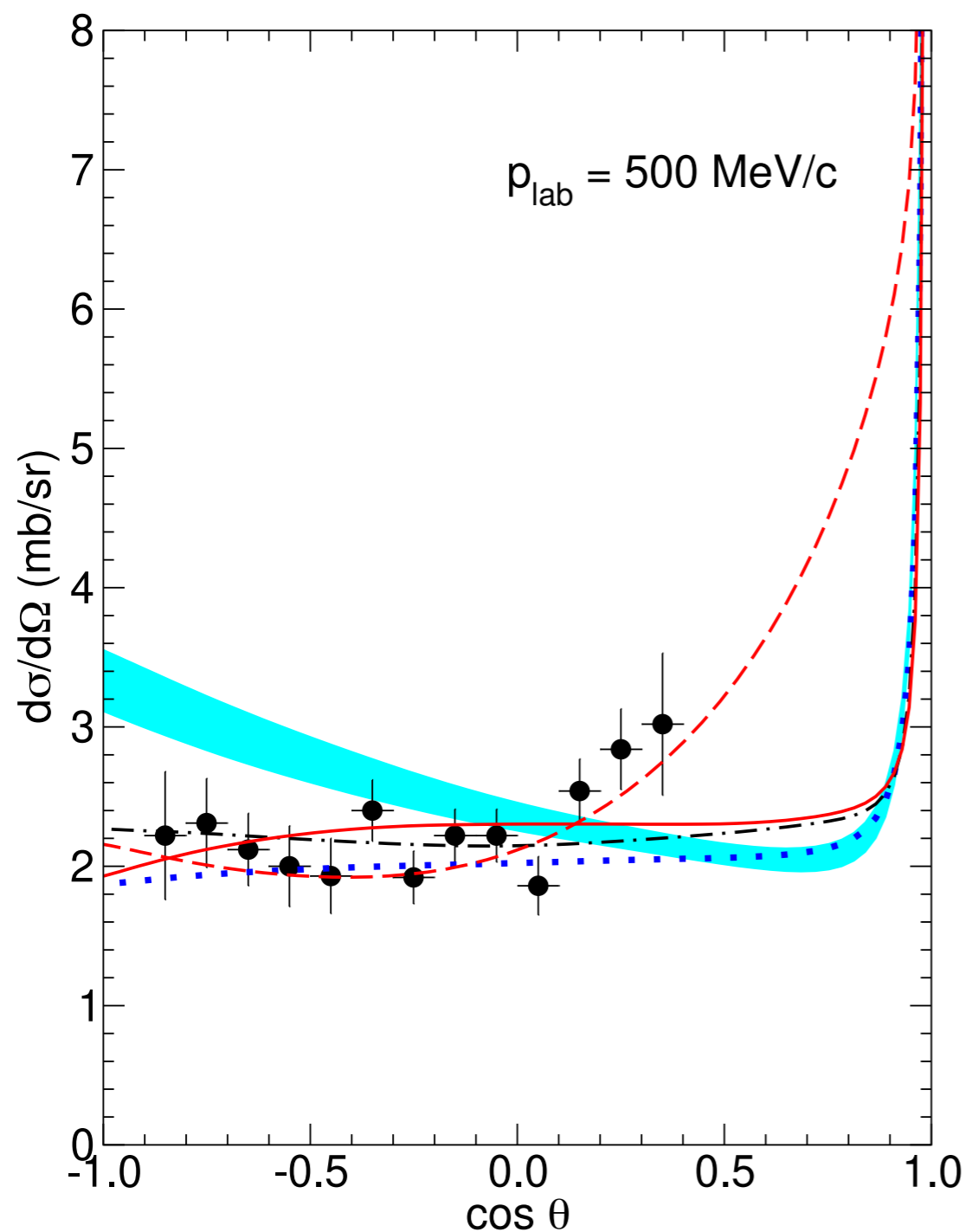


SMS NLO/N²LO interaction

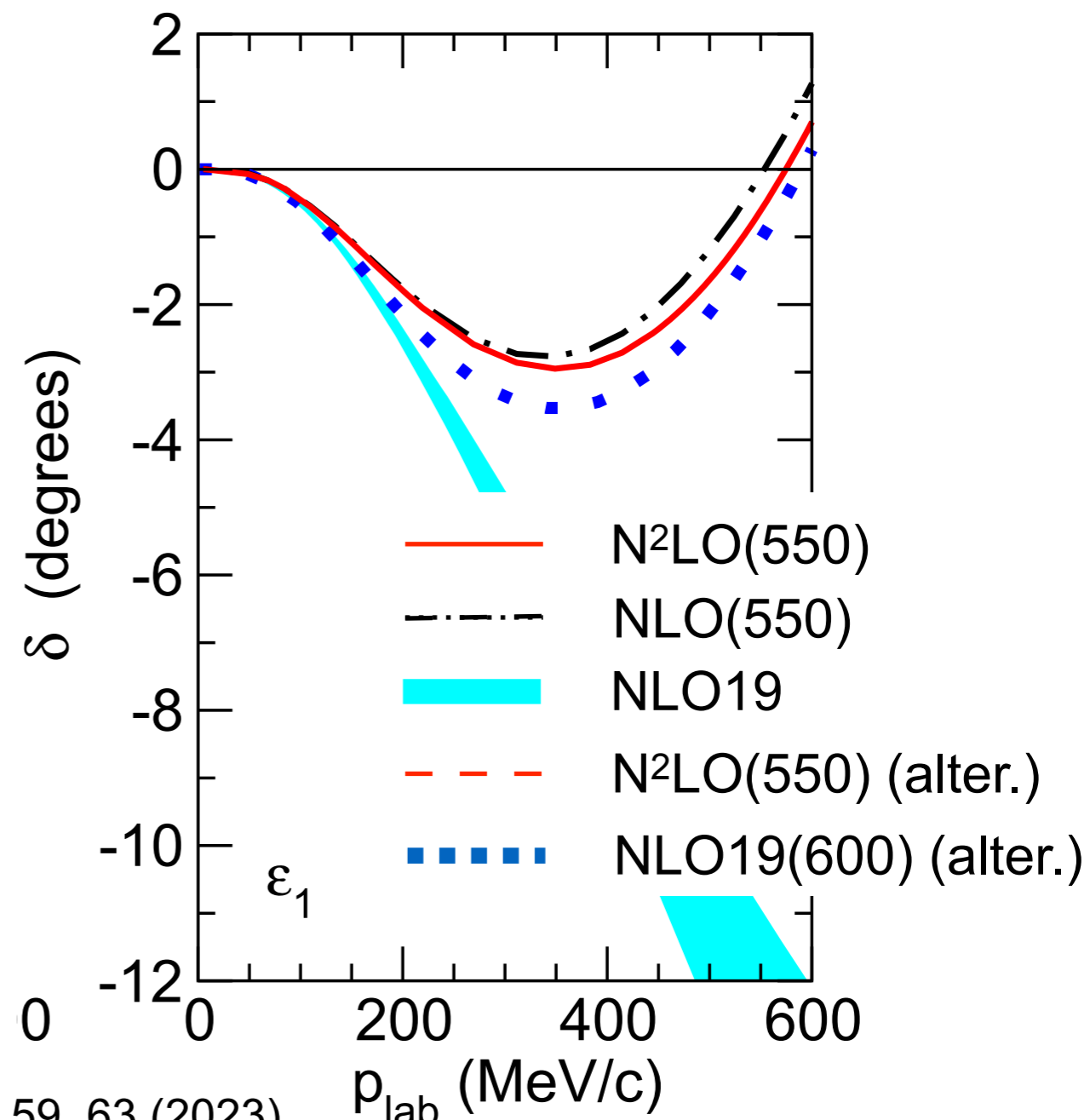


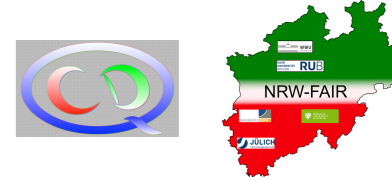
new data (Miwa(2022)) at higher energies provides new constraints!

$$\Sigma^+ p \rightarrow \Sigma^+ p$$



J. Haidenbauer et al. EPJ A 59, 63 (2023).





Notes on the current status of the SMS interactions

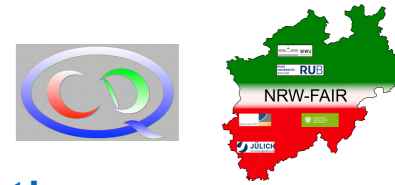
- so far **no YNN forces** (N²LO) have been used in $A > 2$ calculations (see in a moment)
- p-waves are **not uniquely** determined
more accurate hypernuclear calculations and/or additional differential observables (polarizations, more cross differential cross sections, ...)
- data from additional channels will be helpful ($\Lambda p, \Sigma^- p \rightarrow \Sigma^0 n, \dots$)
- calculation for single particle energies in nuclear matter yields results similar to NLO19
— dependence on order and cutoff indicates need to include **YNN forces**
- even ratio of spin singlet/spin triplet strength requires ${}^3_{\Lambda}\text{H}$

Is assumption of negligible YNN force valid for this hypernucleus?

→ uncertainty analysis to

1. pin down dependence on NN force
(motivated by recent work of Gazda et al 2022, Htun et al. 2021)
2. estimate N²LO contribution which quantifies the expected YNN force contribution

Uncertainty analysis to $A = 3$ to 5



Order N²LO requires combination of chiral NN, YN, 3N and **YNN** interaction

Need calculation of separation energies (use Faddeev, Yakubovsky eq. or J-NCSM) and use **different orders** for uncertainty estimate.

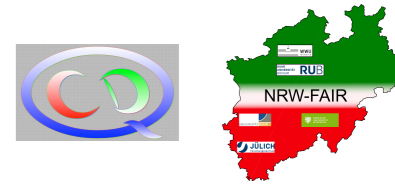
Assuming a negligible numerical uncertainty and the following ansatz for the order by order convergence

$$X_K = X_{ref} \sum_{k=0}^K c_k Q^k \quad \text{where} \quad Q = M_{\pi}^{eff} / \Lambda_b \quad (X_{ref} \text{ LO, exp., max, ...})$$

a **Bayesian analysis** of the uncertainty is possible (see Melendez et al. 2017,2019)

Extracting c_k for $k \leq K$ **from calculations** and assuming identical probability distributions for c_k for $k > K$ the uncertainty is given by the distribution of

$$\delta X_K = X_{ref} \sum_{k=K+1}^{\infty} c_k Q^k$$



Uncertainty analysis to $A = 3$ to 5

How to obtain the distribution for c_k ?

EFT expectation: c_k are natural-sized, i.e. of order 1.

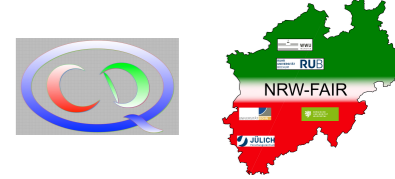
→ defines prior distribution (usually normal distribution with width \bar{c})
 \bar{c} is distributed using an inverse- χ^2 distribution (parameters ν_0, τ_0)

For this choice, the posterior then follows the same distribution (conjugate prior)
 with shifted parameters given the data:

$$\nu = \nu_0 + n_c \quad \nu\tau^2 = \nu_0\tau_0^2 + \vec{c}_k^2 \quad (\vec{c}_k^2 = \sum c_k^2 \text{ for } n_c \text{ values extracted})$$

→ uncertainty follows so-called student t distribution (analytically known)
 allows to extract degree of believe intervals (DoB)

dependence on choice of prior will be less for large n_c !



Uncertainty analysis to $A = 3$ to 5

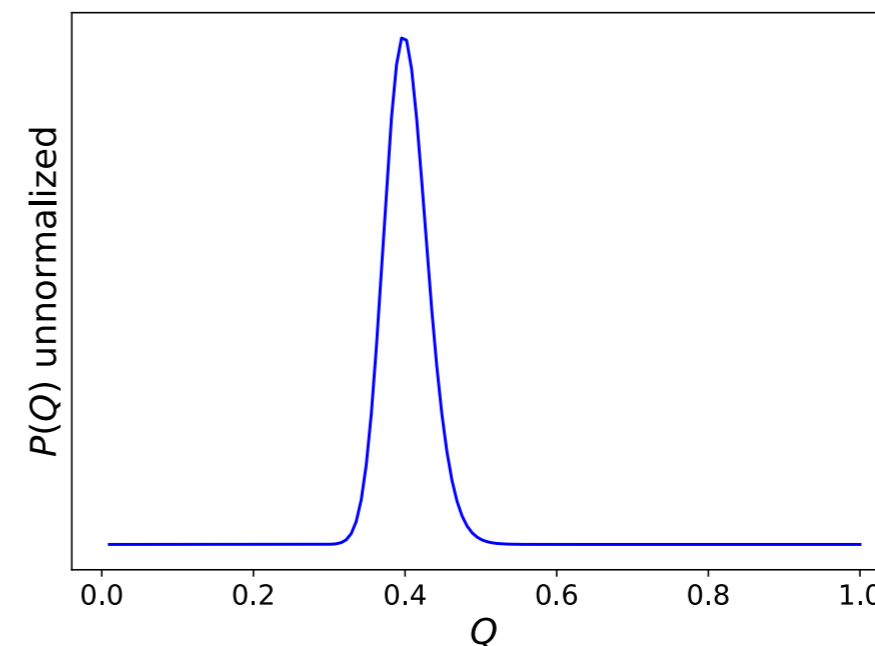
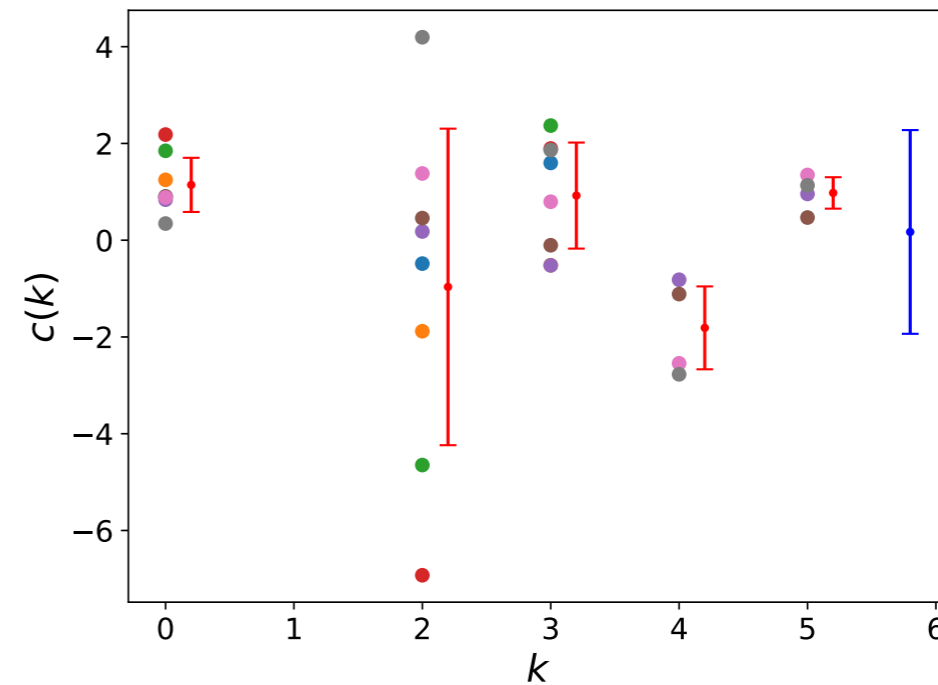
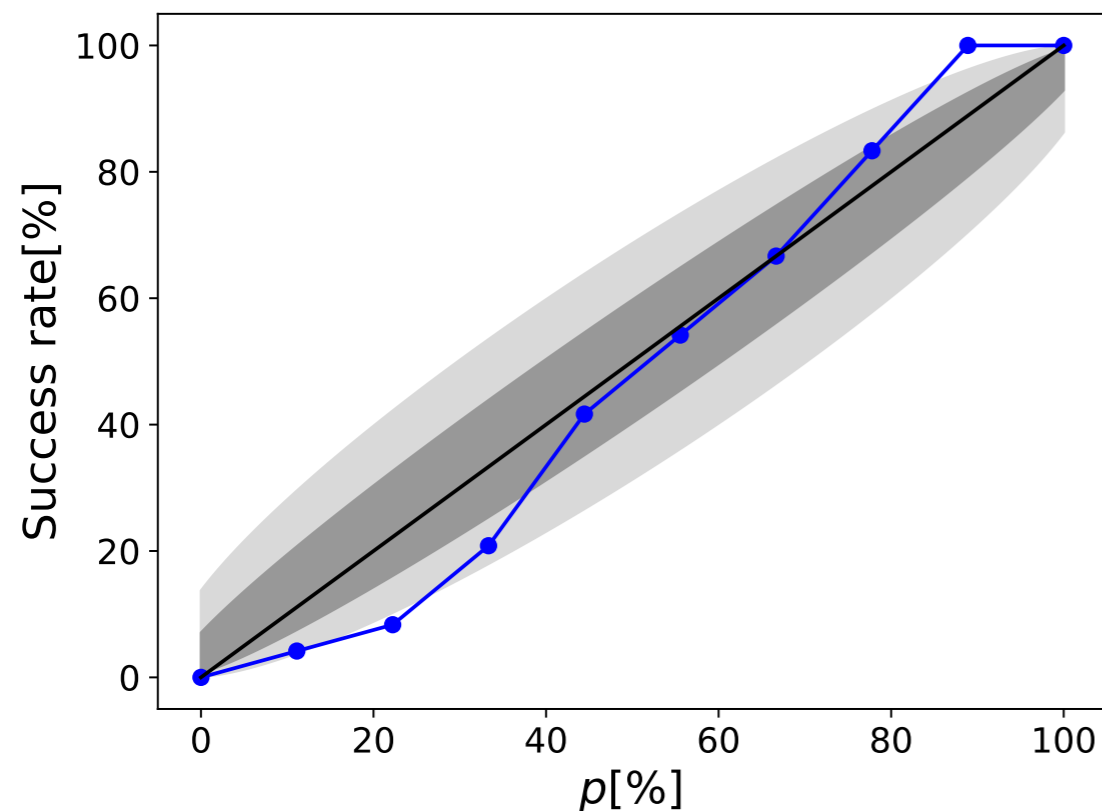
- expansion parameter Q should be consistent with assumption of k independent distribution of c_k
- distribution of prior should be consistent with observed pattern for c_k
- few orders used cannot entirely remove prior dependence

$$Q = 0.4$$

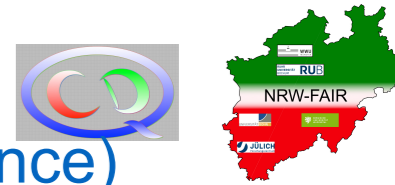
$$\tau_0^2 = 2.25$$

$$\nu_0 = 1.5$$

(see also Maris et al. 2022)

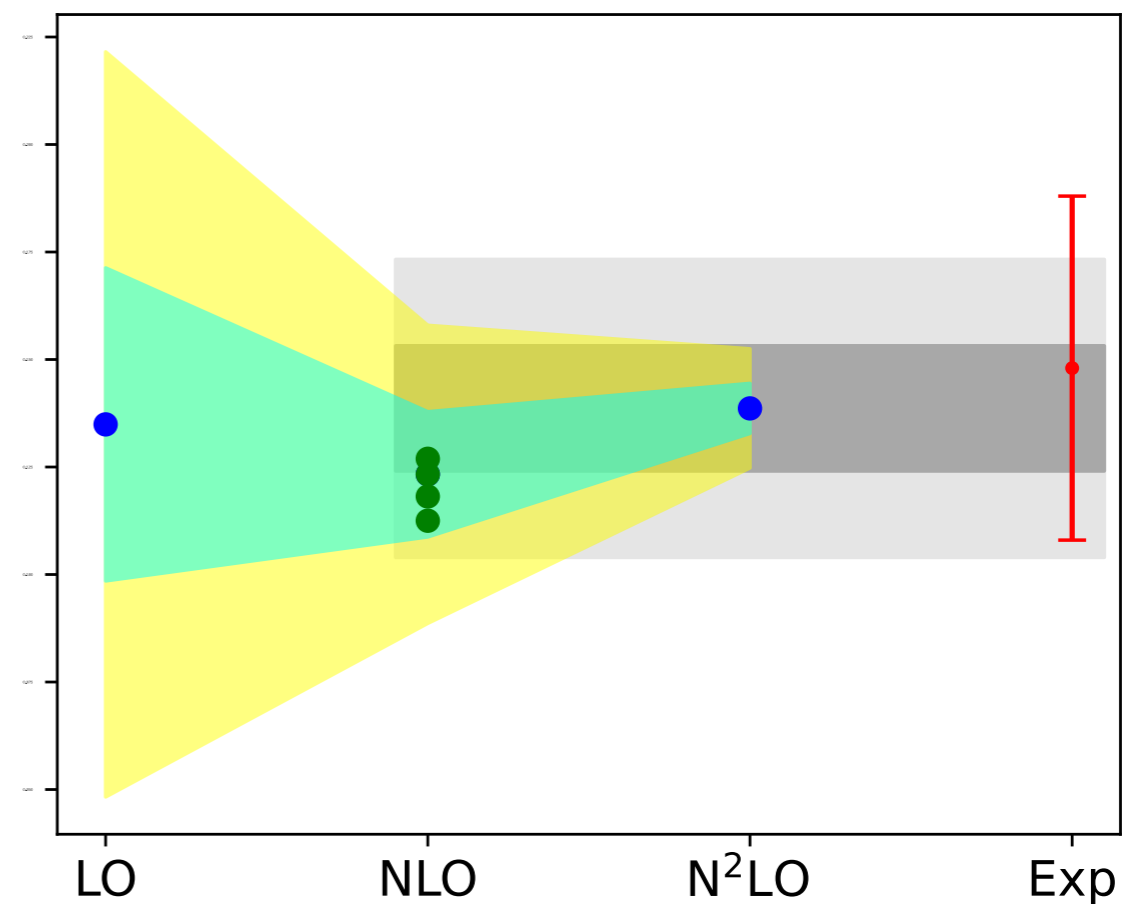
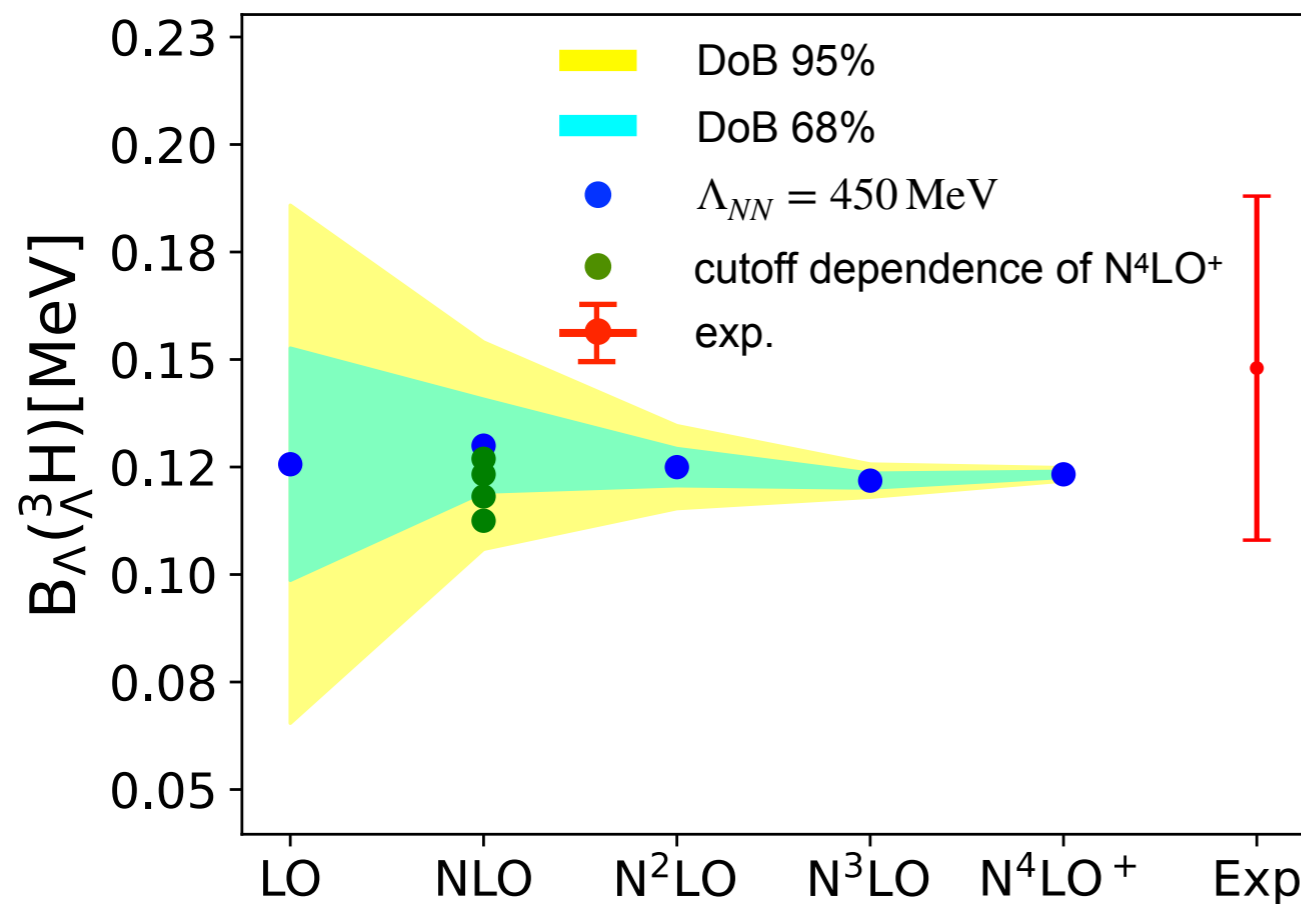


Application to ${}^3_{\Lambda}\text{H}$

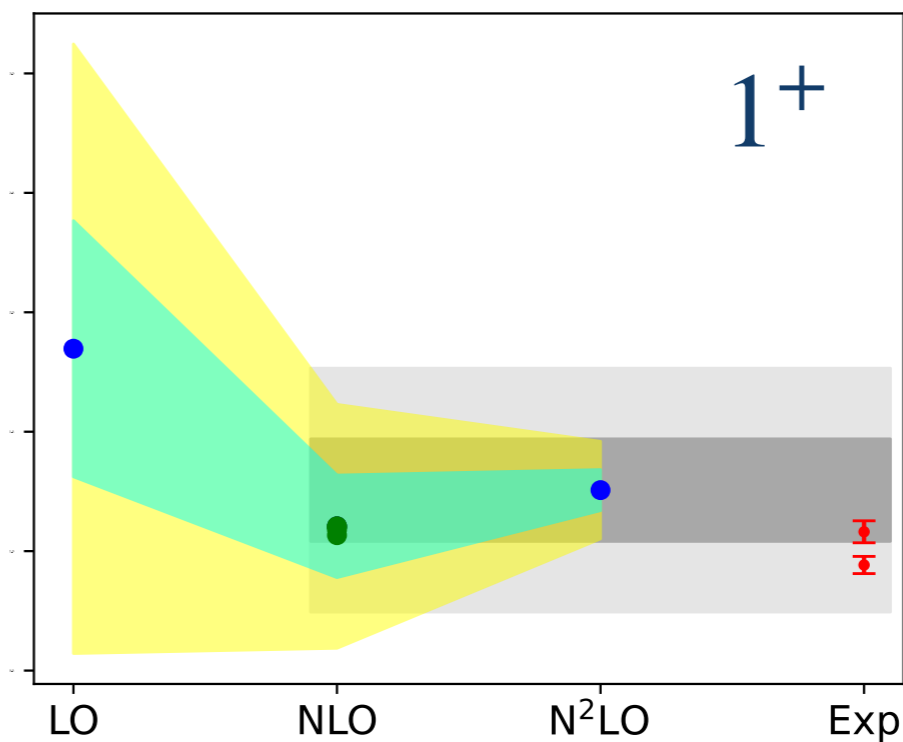
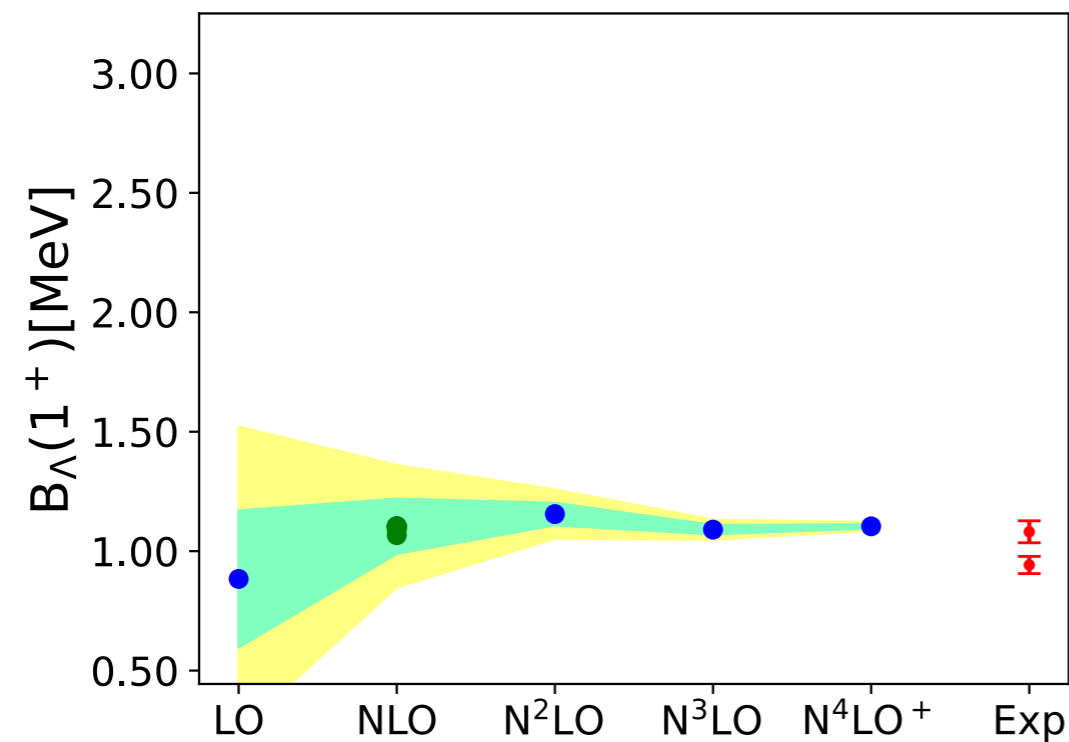
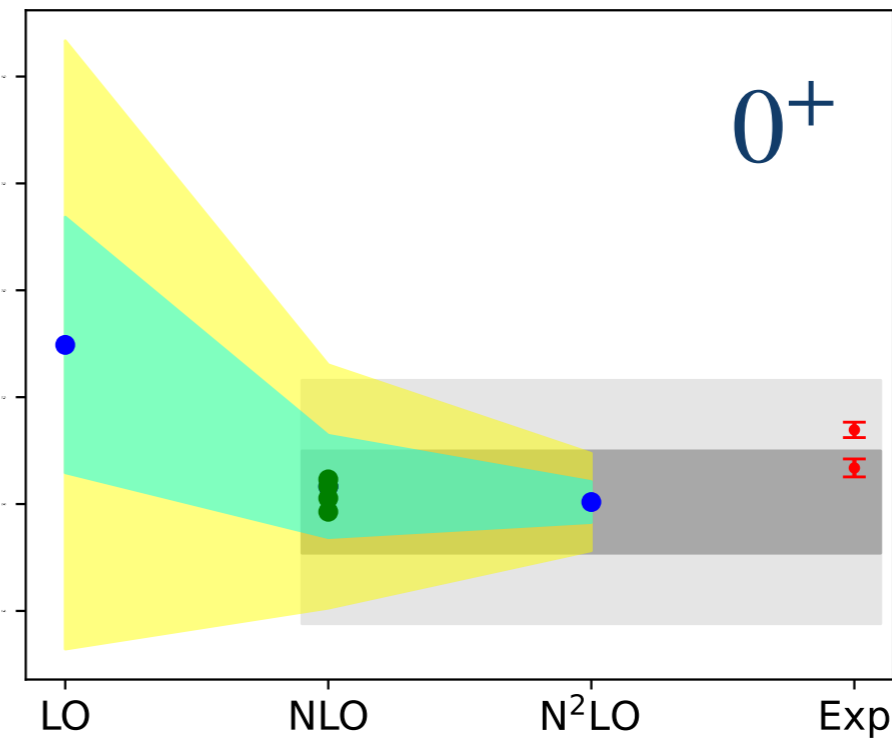
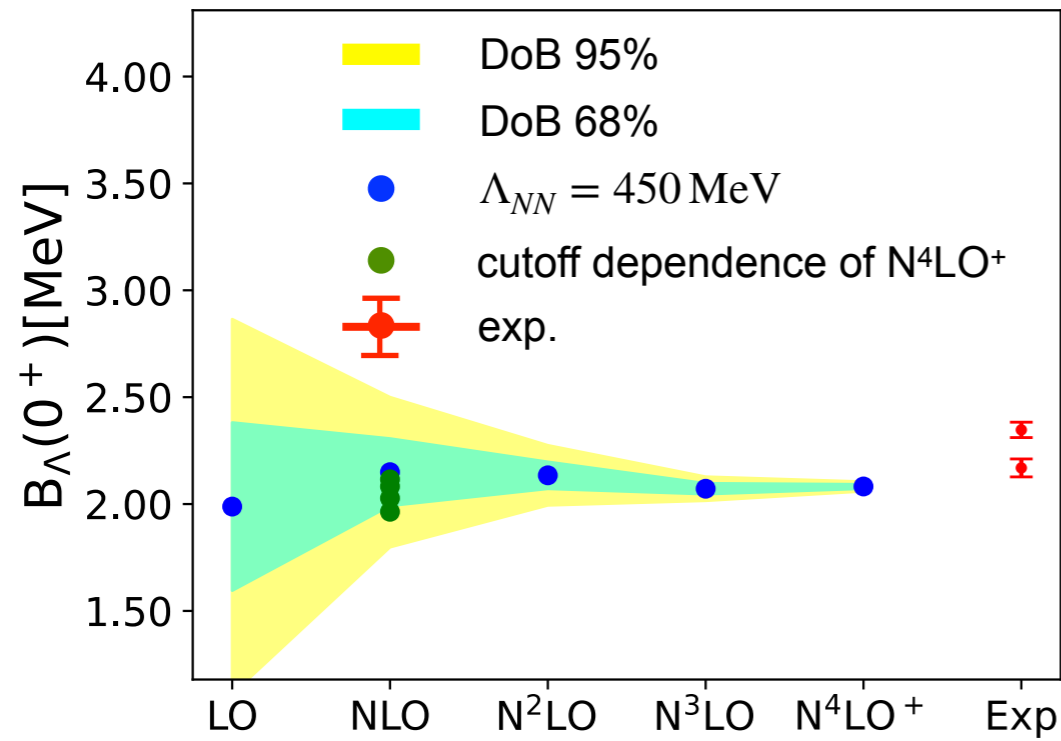
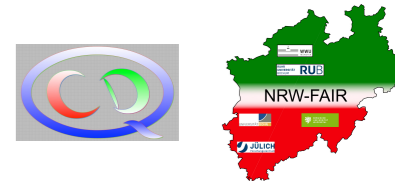


- Q , ν_0 and τ_0 are chosen using all available data (NN and YN convergence)
- uncertainties are extracted using c_k for NN or YN convergence
- use c_k of individual hypernuclei

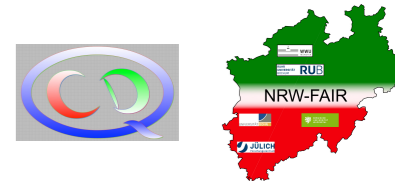
➔ individual uncertainties for NN and YN convergence for each separation energy
consistent with experimental data
cutoff dependence always at least NLO (YNN missing!)



Application to ${}^4_{\Lambda}\text{He}$

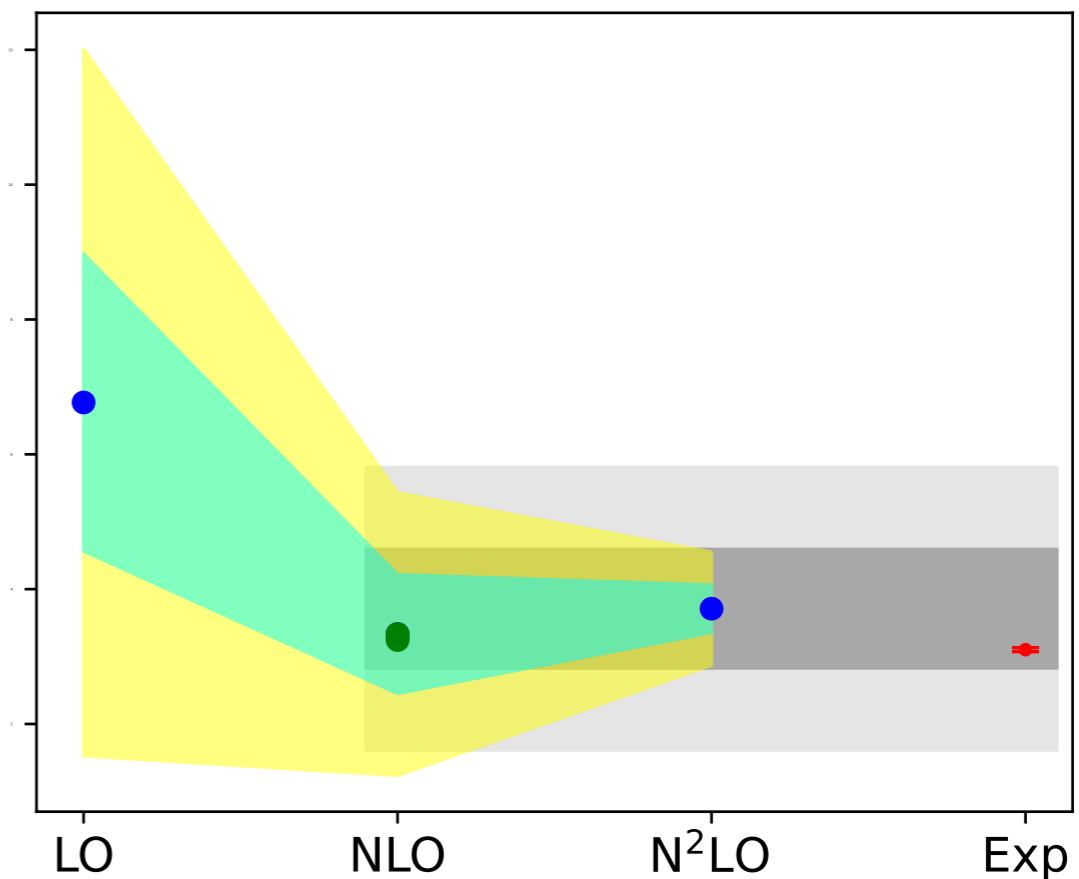
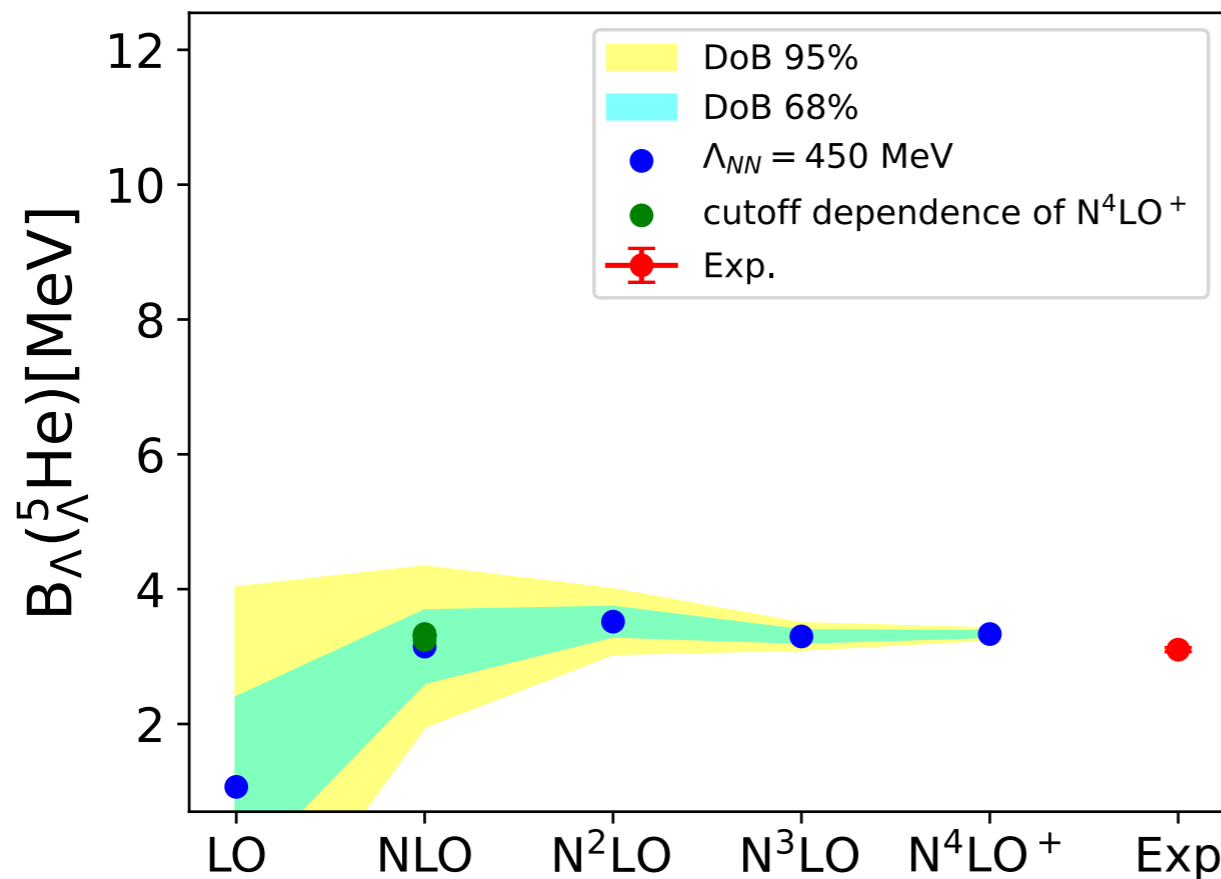


Application to ${}^5_{\Lambda}\text{He}$ and summary



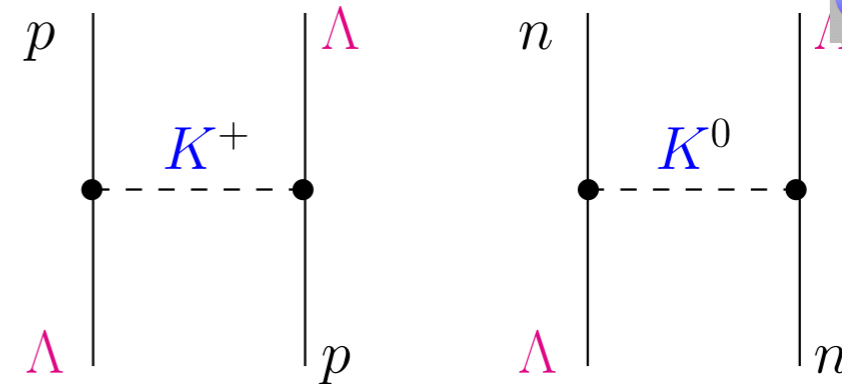
- without YNN: sizable uncertainties at $A = 4$ and 5
- $A = 3$ sufficiently accurate
- NN/YN dependence small at least for $A = 3$

nucleus	$\Delta_{68}(NN)$	$\Delta_{68}(YN)$
${}^3_{\Lambda}\text{H}$	0.011	0.015
${}^4_{\Lambda}\text{He} (0^+)$	0.157	0.239
${}^4_{\Lambda}\text{He} (1^+)$	0.114	0.214
${}^5_{\Lambda}\text{He}$	0.529	0.881



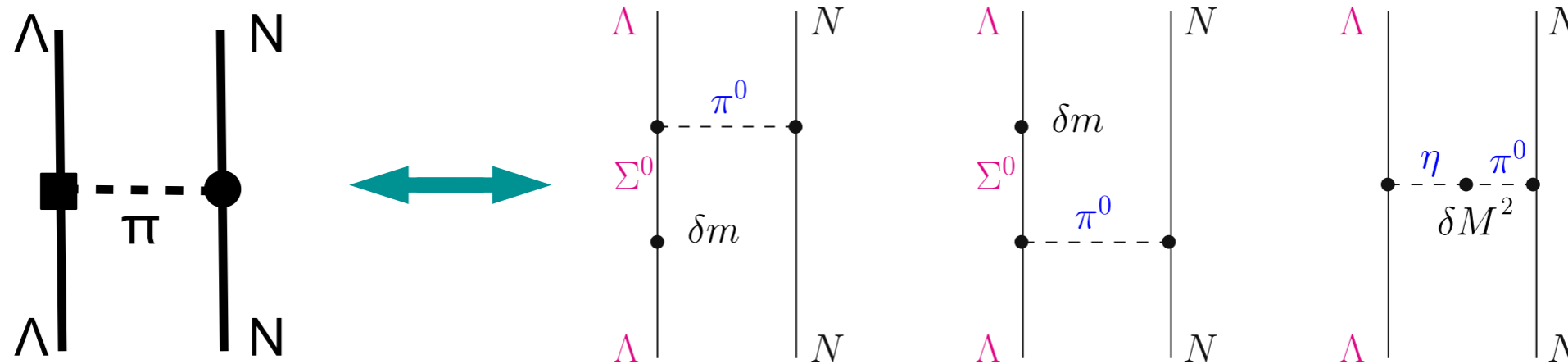
CSB contributions to ΛN interactions

- **formally leading** contributions:
Goldstone boson mass difference
 - very small due to the small relative difference of kaon masses



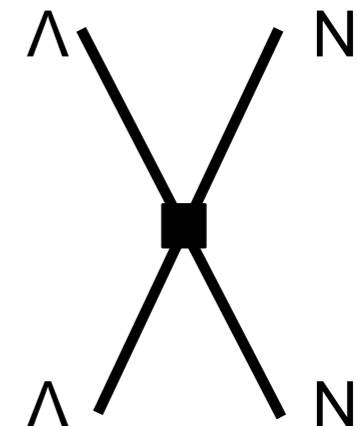
- **subleading but most important**
 - effective CSB $\Lambda\Lambda\pi$ coupling constant (Dalitz, van Hippel, 1964)

$$f_{\Lambda\Lambda\pi} = \left[-2 \frac{\langle \Sigma^0 | \delta m | \Lambda \rangle}{m_{\Sigma^0} - m_{\Lambda}} + \frac{\langle \pi^0 | \delta M^2 | \eta \rangle}{M_{\eta}^2 - M_{\pi^0}^2} \right] f_{\Lambda\Sigma\pi} \approx (-0.0297 - 0.0106) f_{\Lambda\Sigma\pi}$$

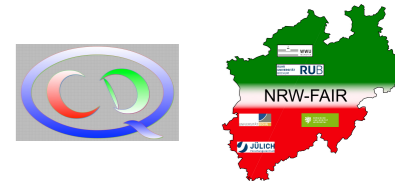


- **so far less considered, but equally important**
 - CSB contact interactions (for singlet and triplet)

Aim: use $A=4$ hypernuclei to determine the two unknown CSB LECs and predict Λn scattering



Fit of contact interactions



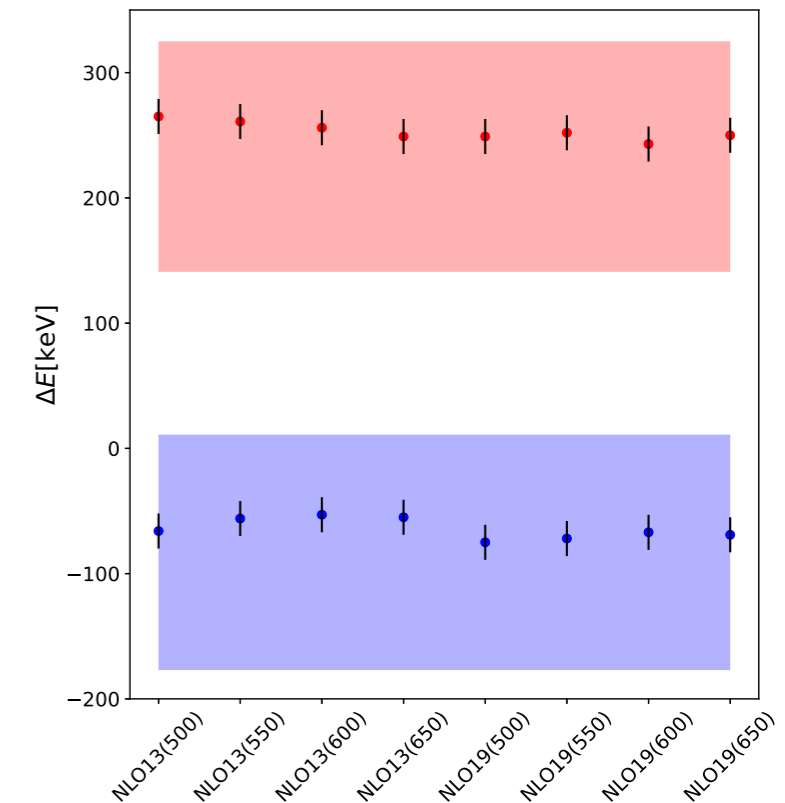
- Adjust the two CSB contact interactions to one main scenario (**CSB1**)
- Size of LECs as expected by power counting

$$\frac{m_d - m_u}{m_u + m_d} \left(\frac{M_\pi}{\Lambda} \right)^2 C_{S,T} \approx 0.3 \cdot 0.04 \cdot 0.5 \cdot 10^4 \text{ GeV} \propto 6 \cdot 10^{-3} \cdot 10^4 \text{ GeV}$$

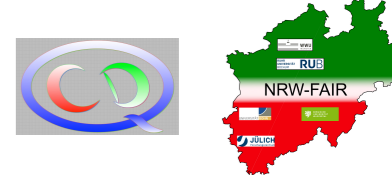
Λ	NLO13		NLO19	
	C_s^{CSB}	C_t^{CSB}	C_s^{CSB}	C_t^{CSB}
500	4.691×10^{-3}	-9.294×10^{-4}	5.590×10^{-3}	-9.505×10^{-4}
550	6.724×10^{-3}	-8.625×10^{-4}	6.863×10^{-3}	-1.260×10^{-3}
600	9.960×10^{-3}	-9.870×10^{-4}	9.217×10^{-3}	-1.305×10^{-3}
650	1.500×10^{-2}	-1.142×10^{-3}	1.240×10^{-2}	-1.395×10^{-3}

The values of the LECs are in 10^4 GeV^{-2}

- Problem: large experimental uncertainty of experiment
- here only **fit to central values** to test theoretical uncertainties



Prediction for Λn scattering



- assuming the current experimental situation for ${}^4_{\Lambda}\text{H}$ / ${}^4_{\Lambda}\text{He}$
- **without CSB:** $a_s^{\Lambda n} \approx 2.9 \text{ fm}$ **with CSB1:** $a_s^{\Lambda n} \approx 3.3 \text{ fm}$
- improved description of Λp data
- almost independent of cutoff & NLO variant
- CSB of triplet is smaller than of singlet

for "**CSB1**": currently accepted experimental values

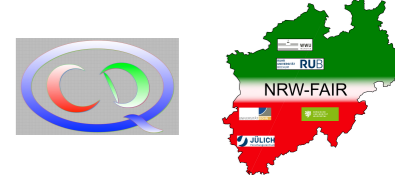
	$a_s^{\Lambda p}$	$a_t^{\Lambda p}$	$a_s^{\Lambda n}$	$a_t^{\Lambda n}$	$\chi^2(\Lambda p)$	$\chi^2(\Sigma N)$	$\chi^2(\text{total})$
NLO13(500)	-2.604	-1.647	-3.267	-1.561	4.47	12.13	16.60
NLO13(550)	-2.586	-1.551	-3.291	-1.469	3.46	12.03	15.49
NLO13(600)	-2.588	-1.573	-3.291	-1.487	3.43	12.38	15.81
NLO13(650)	-2.592	-1.538	-3.271	-1.452	3.70	12.57	16.27
NLO19(500)	-2.649	-1.580	-3.202	-1.467	3.51	14.69	18.20
NLO19(550)	-2.640	-1.524	-3.205	-1.407	3.23	14.19	17.42
NLO19(600)	-2.632	-1.473	-3.227	-1.362	3.45	12.68	16.13
NLO19(650)	-2.620	-1.464	-3.225	-1.365	3.28	12.76	16.04

An accurate prediction for the Λn interaction is possible using hypernuclei!

remeasurement of ${}^4_{\Lambda}\text{H}$ excitation energy to match accuracy for ${}^4_{\Lambda}\text{He}$?

measurement of ${}^4_{\Lambda}\text{He}$ ground state at J-PARC

Conclusions & Outlook



- **YN interactions not well understood**
 - *scarce YN data*
 - *more information necessary to solve "hyperon puzzle"*
- **Hypernuclei provide important constraints**
 - *CSB of ΛN scattering & ${}^4_{\Lambda}\text{He}$ / ${}^4_{\Lambda}\text{H}$*
 - *${}^3_{\Lambda}\text{H}$ is used to constrain the spin dependence*
 - *new experiments planned at J-PARC, MAMI, J-Lab, FAIR,...*
- **New SMS YN interactions**
 - *give an accurate description low energy YN data*
 - *order LO, NLO and N²LO allow uncertainty quantification*
 - *have a non-unique determination of contact interactions (data necessary)*
- **Chiral 3BF need to be included**
 - *NLO uncertainty is sizable in $A = 4$ and 5*
 - *chiral 3BFs are formulated (Petschauer et al., (2016)) and the implementation is currently checked*