Status of the hyperon-nucleon interaction in chiral effective field theory

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(Hoai Le, Ulf-G. Meißner, Andreas Nogga)

Outline

BB interaction in chiral effective field theory

Baryon-baryon interaction in SU(3) χEFT à la Weinberg (1990)

Advantages:

- Power counting
  systematic improvement by going to higher order
- Possibility to derive two- and three-baryon forces and external current operators in a consistent way

- degrees of freedom: octet baryons (N, Λ, Σ, Ξ), pseudoscalar mesons (π, K, η)

- pseudoscalar-meson exchanges
- contact terms – represent unresolved short-distance dynamics
  involve low-energy constants (LECs) that need to be fixed
  by a fit to data

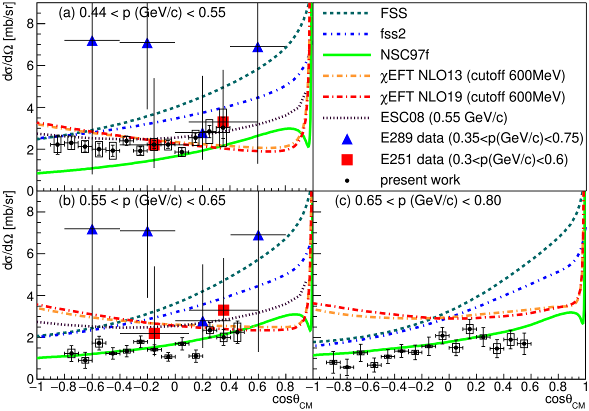
ΛN-ΣN interaction

LO: H. Polinder, J.H., U.-G. Meißner, NPA 779 (2006) 244
NLO13: J.H., S. Petschauer, N. Kaiser, U.-G. Meißner, A. Nogga, W. Weise, NPA 915 (2013) 24
NLO19: J.H., U.-G. Meißner, A. Nogga, EPJA 56 (2020) 91

(*BB* systems with strangeness $S = -1$ to -6)

New ΣN data from E40 Collaboration at J-PARC

$\Sigma^+ p$: T. Nanamura et al., arXiv:2203:08393



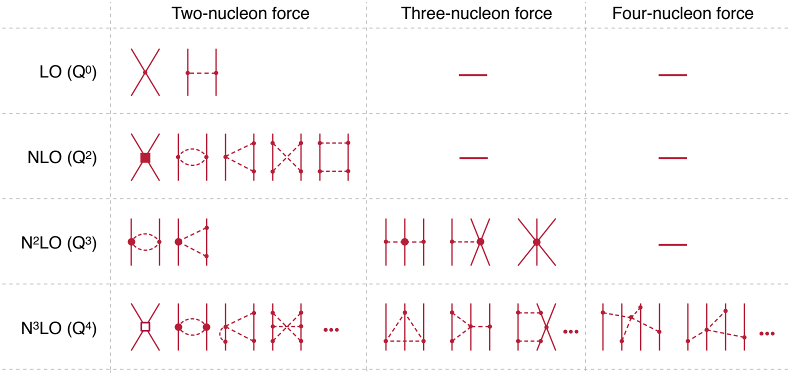
($\Sigma^- p$: K. Miwa et al., PRC 104 (2021) 045204; $\Sigma^- p \to \Lambda n$: K. Miwa et al., PRL 128 (2021) 072501)

$p_{lab} = 500$ MeV/c ($E_{lab} = 100.7$ MeV); $p_{lab} = 600$ MeV/c ($E_{lab} = 142.8$ MeV)

beyond of validity of NLO interaction?; role of higher partial waves?
($\Lambda p \pi^+$ threshold is at $p_{lab} \approx 600$ MeV/c)

Extension of chiral EFT interaction up to NNLO

(Nucleon-nucleon forces in chiral EFT (E. Epelbaum))

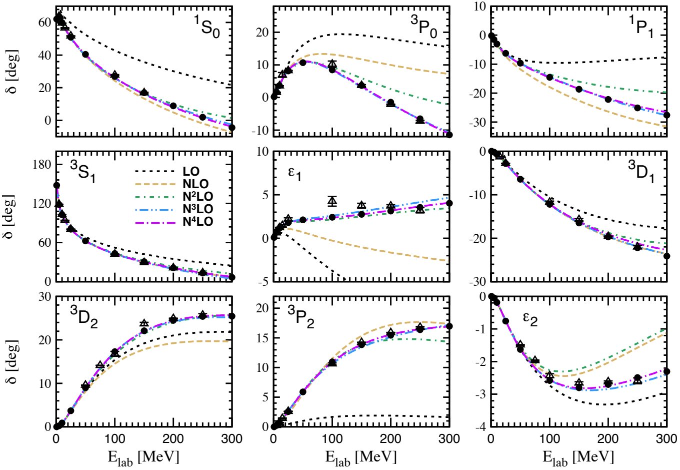


N^{2}LO: no new (additional) LECs in the two-body sector

leading-order three-body forces (3BFs)

NN interaction in chiral EFT

Semilocal momentum-space (SMS) regularized chiral *NN* potential



(Reinert, Krebs, Epelbaum, EPJA 54 (2018) 86) [up to N^{4}LO (N^{4}LO$^+$) !!]

LO to NLO: drastic change in all partial waves

NLO to N^{2}LO: changes mostly in *P*-waves and higher partial waves

chiral YN potential up to NNLO

adopt the framework of Reinert, Krebs, Epelbaum, EPJA 54 (2018) 86:

"Semilocal momentum-space regularized (SMS) chiral NN potentials"

- employ a regulator that minimizes artifacts from cutoff Λ

nonlocal cutoff $\quad (\vec{q} = \vec{p}\,' - \vec{p})$

$$V_{1\pi}^{\rm reg} \propto \frac{e^{-\frac{p'^4+p^4}{\Lambda^4}}}{\vec{q}^2+m_\pi^2} \to \frac{1}{\vec{q}^2+m_\pi^2}\left[1-\frac{p'^4+p^4}{\Lambda^4}+\mathcal{O}(\Lambda^{-8})\right]$$

local cutoff:

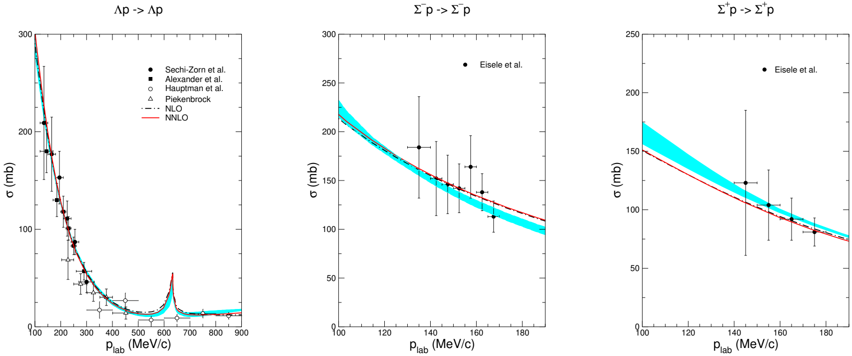
$$V_{1\pi}^{\rm reg} \propto \frac{e^{-\frac{\vec{q}^2+m_\pi^2}{\Lambda^2}}}{\vec{q}^2+m_\pi^2} \to \frac{1}{\vec{q}^2+m_\pi^2} - \frac{1}{\Lambda^2} + \frac{\vec{q}^2+m_\pi^2}{\Lambda^4} + \ldots$$

does not affect long-range physics at any order in the $1/\Lambda^2$ expansion

applicable to 2π exchange too:

$$V_{2\pi} = \frac{2}{\pi}\int_{2m_\pi}^{\infty}\mu d\mu \frac{\rho(\mu)}{\vec{q}^2+\mu^2} \to V_{2\pi}^{\rm reg} = e^{-\frac{\vec{q}^2}{2\Lambda^2}}\frac{2}{\pi}\int_{2m_\pi}^{\infty}\mu d\mu \frac{\rho(\mu)}{\vec{q}^2+\mu^2}e^{-\frac{\mu^2}{2\Lambda^2}} \quad + \ldots$$

Preliminary results for SMS chiral YN interactions

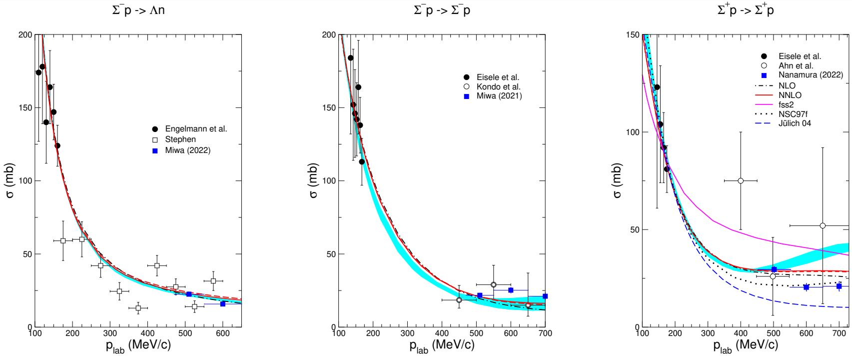


SMS YN potentials up to NLO, NNLO (with $\Lambda = 550$ MeV)
NLO19: J.H., U.-G. Meißner, A. Nogga, EPJA 56 (2020) 91

quality of the fit – total χ^2 (36 data points):
NLO19(600): 16.0 SMS NLO: 15.2 SMS NNLO: 15.6

cross sections dominated by S-waves (are already well described at NLO)
$\rightarrow$ (as expected) practically no change when going to NNLO

Preliminary results for SMS YN interactions



integrated cross sections at higher energies not included in the fitting process!

$\Sigma^+ p \to \Sigma^+ p$ and $\Sigma^- p \to \Sigma^- p$ cross sections:

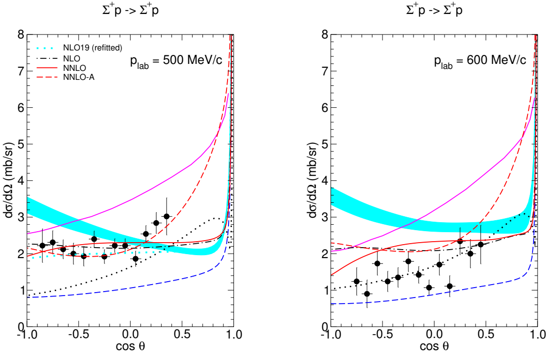
$$\sigma = \frac{2}{\cos\theta_{\max} - \cos\theta_{\min}} \int_{\cos\theta_{\min}}^{\cos\theta_{\max}} \frac{d\sigma(\theta)}{d\cos\theta} d\cos\theta$$

$\cos\theta_{\min} = -0.5; \ \cos\theta_{\max} = 0.5$

fss2 ... Fujiwara et al. (constitutent quark model)   Jülich 04, Nijmegen NSC97f ... meson-exchange potentials

Preliminary results for SMS YN interactions

$\Sigma^+ p$ (T. Nanamura et al., arXiv:2203:08393)



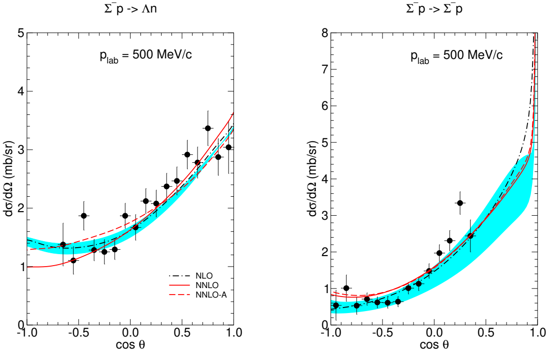
LECs in the $^1S_0$, $^3S_1$-$^3D_1$ fixed from low-energy cross sections

SMS NLO: LECs in ^3P-waves taken over from NN fit (RKE)
(strict SU(3) symmetry: $V_{NN} \equiv V_{\Sigma^+ p}$ in the $^1S_0$, $^3P_{0,1,2}$ partial waves!)

SMS NNLO: LECs in P-waves fitted to the E40 data (two examples)!

data for $(550 \leq p \leq 650)$ MeV/c are overestimated (influence of $\Lambda p \pi^+$ threshold?)

Preliminary results for SMS YN interactions

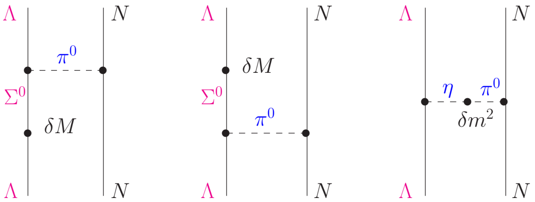


$\Sigma^- p \to \Lambda n$: quite well reproduced by NLO19 (NLO13) and SMS YN potentials
$\Sigma^- p \to \Sigma^- p$: behavior at forward angles remains unclear

$\Sigma^- p$ and $\Sigma^- p \to \Lambda n$ data for $(550 \le p \le 650)$ MeV/c are reproduced with comparable quality

- no unique determination of all P-wave LECs possible
- one needs data from additional channels (Λp, $\Sigma^- p \to \Sigma^0 n$, ...)
- one needs additional differential observables (polarizations, ...)

Charge symmetry breaking in the ΛN interaction



CSB due to $\Lambda - \Sigma^0$ mixing leads to a long-ranged contribution to the ΛN interaction
(R.H. Dalitz & F. von Hippel, PL 10 (1964) 153)

Strength can be estimated from the electromagnetic mass matrix:

$$\langle\Sigma^0|\delta M|\Lambda\rangle = [M_{\Sigma^0} - M_{\Sigma^+} + M_p - M_n]/\sqrt{3}$$

$$\langle\pi^0|\delta m^2|\eta\rangle = [m^2_{\pi^0} - m^2_{\pi^+} + m^2_{K^+} - m^2_{K^0}]/\sqrt{3}$$

$$f_{\Lambda\Lambda\pi} = [-2\frac{\langle\Sigma^0|\delta M|\Lambda\rangle}{M_{\Sigma^0} - M_\Lambda} + \frac{\langle\pi^0|\delta m^2|\eta\rangle}{m^2_\eta - m^2_{\pi^0}}]\, f_{\Lambda\Sigma\pi}$$

latest PDG mass values $\Rightarrow$

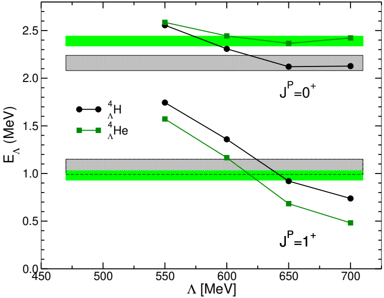
$$f_{\Lambda\Lambda\pi} \approx (-0.0297 - 0.0106)\, f_{\Lambda\Sigma\pi} \approx -0.0403\, f_{\Lambda\Sigma\pi}$$

CSB in ${}^4_\Lambda$H, ${}^4_\Lambda$He by Gazda and Gal

D. Gazda and A. Gal, NPA 954 (2016) 161: assume that

$$V^{CSB}{}_{\Lambda N\to\Lambda N} = -2\frac{\langle\Sigma^0|\delta M|\Lambda\rangle}{M_{\Sigma^0}-M_\Lambda}\tau_{N_z}\frac{1}{\sqrt{3}}\,V_{\Lambda N\to\Sigma N} \qquad \tau_{N_z} = 1(\mathrm{p});\ -1(\mathrm{n})$$

use our LO YN interaction (calculations in the no-core shell model)



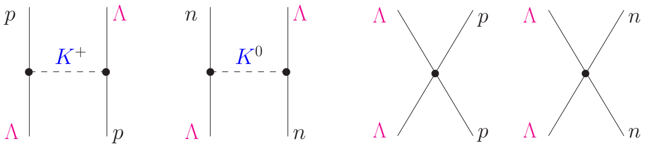
- splitting for the 1^+ state somewhat too large
- fairly strong cutoff dependence

$\Rightarrow$ EFT: the latter signals that something is missing!

CSB in chiral EFT

CSB (CIB) in χEFT: worked out for *pp*, *nn* (and *np*) scattering

Walzl, Meißner, Epelbaum, NPA 693 (2001) 663; Epelbaum, Glöckle, Meißner, NPA 747 (2005) 362
J. Friar et al., PRC 68 (2003) 024003
LØ: Coulomb interaction, m_{π^0}-$m_{\pi^\pm}$ in OPE   NLØ: isospin breaking in $f_{NN\pi}$, leading-order contact terms



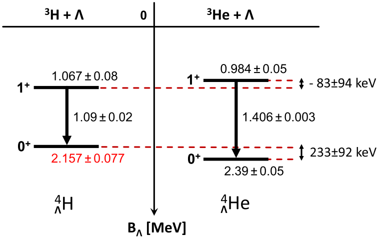
Gazda/Gal results: short-distance dynamics is relevant
→ one has to account for that by appropriate contact terms
(in line with the power counting)

$NN\ ^1S_0$: $a_{pp} - a_{nn} \approx 1.5$ fm
mostly due to short-range forces (ρ^0-ω mixing, a_1^0-f_1 mixing)

Faddeev-Yakubovsky calculation for NLO13 and NLO19 interactions
with CSB forces including contact terms:
(J.H., U.-G. Meißner, A. Nogga, FBS 62 (2021) 105)

Charge symmetry breaking in ${}^4_\Lambda$H-${}^4_\Lambda$He

- $\Delta E(0^+) = E^{0^+}_\Lambda({}^4_\Lambda\mathrm{He}) - E^{0^+}_\Lambda({}^4_\Lambda\mathrm{H})$
  $= 233 \pm 92$ keV
- $\Delta E(1^+) = E^{1^+}_\Lambda({}^4_\Lambda\mathrm{He}) - E^{1^+}_\Lambda({}^4_\Lambda\mathrm{H})$
  $= -83 \pm 94$ keV



adjust CSB contact terms to ΔE's

(Schulz et al., 2016; Yamamoto et al., 2015)

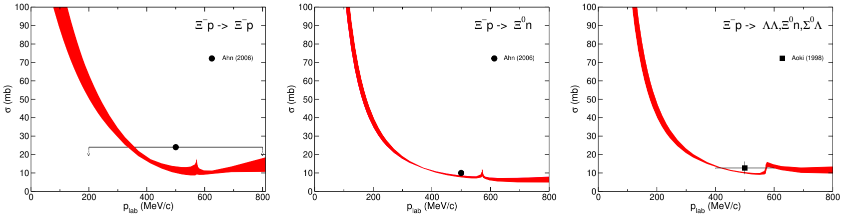
(fm // keV)	$a_s^{\Lambda p}$	$a_s^{\Lambda n}$	$a_t^{\Lambda p}$	$a_t^{\Lambda n}$	$\Delta E(0^+)$	$\Delta E(1^+)$
NLO19(500)	-2.649	-3.202	-1.580	-1.467	249	-75
NLO19(550)	-2.640	-3.205	-1.524	-1.407	252	-72
NLO19(600)	-2.632	-3.227	-1.473	-1.362	243	-67
NLO19(650)	-2.620	-3.225	-1.464	-1.365	250	-69

CSB in singlet (${}^1S_0$) much larger than in triplet (${}^3S_1$)
practically independent of cutoff; same results for NLO13
without CSB: $a_s^{\Lambda p} \approx a_s^{\Lambda n} \approx -2.9$ fm

- CSB in $A = 7, 8$ Λ-hypernuclei, see talk of Hoai Le

Selected results for the ΞN system

(J.K. Ahn et al., PLB 633 (2006) 214; S. Aoki et al., NPA 644 (1998) 365)



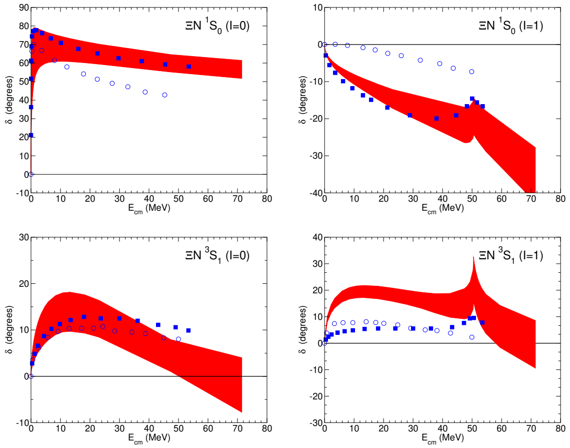
ΞN scattering lengths [in fm]:

	$I=0,\ ^1S_0$	$I=1,\ ^1S_0$		$I=0,\ ^3S_1$		$I=1,\ ^3S_1$	
potential	a_s	a_s	r_s	a_t	r_t	a_t	r_t
NLO (500)	-7.71-i2.03	0.37	-4.80	-0.33	-6.86	-1.17	3.44
NLO (550)	-7.24-i20.79	0.39	-4.95	-0.39	-1.77	-1.15	3.80
NLO (600)	-10.89-i14.91	0.34	-7.20	-0.62	1.00	-1.13	3.95
NLO (650)	-8.14-i2.43	0.31	-9.16	-0.85	1.42	-0.90	4.27

- scattering lengths $|a| \lesssim 1$ fm, except for $I=0,\ ^1S_0$
- ΞN interaction is fairly weak

J.H., U.-G. Meißner, EPJA 58 (2019) 23

ΞN: Comparison with HAL QCD results



HAL QCD Collaboration (almost at physical point, $m_\pi \approx 145$ MeV):
open circles from E. Hiyama et al., PRL 124 (2020) 092501 (no $\Lambda\Sigma$, $\Sigma\Sigma$)
filled squares from M. Kohno & K. Miyagawa, PTEP 2021 (2021) 103D04

Nuclear matter properties

$U_\Xi(p_\Xi = 0)$ [in MeV] at saturation density, $k_F = 1.35\ \mathrm{fm}^{-1}$ ($\rho_0 = 0.166\ \mathrm{fm}^{-3}$)

potential	I	${}^1S_0$	${}^3S_1$	S-waves	P-waves	total
NLO (500)	0	-2.6	-3.3			
	1	12.7	-11.8	-5.0	-0.4	-5.5
NLO (550)	0	-2.9	-3.1			
	1	12.4	-9.5	-3.1	-0.7	-3.8
NLO (550)*	0	-3.15	-3.24			
	1	9.64	-11.0	-7.7	-1.1	-8.8
HAL QCD	0	-3.15	-5.36			
	1	7.12	-2.41	-4.11	-	-4.11
Ehime	0	-0.80	0.47			
(1.82)	1	-1.5	-8.6	-10.43	-11.4	-21.8

"traditional" value for the depth of the Ξ single-particle potential: ≈ -15 MeV

E. Friedman & A. Gal (optical potential, PLB 820 (2021) 136555): $U_\Xi \leq -20$ MeV

Y. Tanimura et al. (relativistic mean field, PRC 105 (2022) 044324): $U_\Xi \approx -12$ MeV

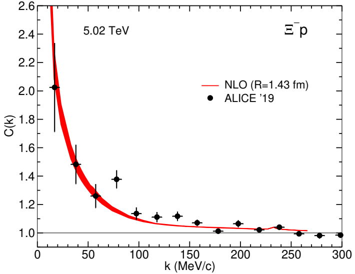
(from analyzing ${}^{15}_{\Xi}$C and ${}^{12}_{\Xi}$Be events)

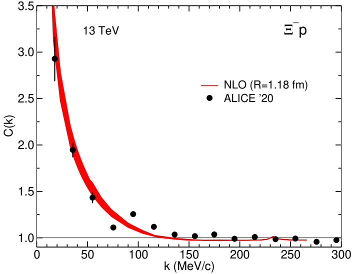
NLO (550)*: M. Kohno, PRC 100 (2019) 024313 (continuous prescription)

HAL QCD: T. Inoue, AIP Conf. Proc. 2130 (2019) 020002

Ehime: M. Yamaguchi et al., PTP 105 (2001) 627

ΞN: two-particle momentum correlation functions





$C_{\text{th}}(k) = \frac{1}{4} C_{^1S_0}(k) + \frac{3}{4} C_{^3S_1}(k)$; $\quad C_\alpha(k) \simeq 1 + \int_0^\infty 4\pi r^2\, dr\, S_{12}(\mathbf{r}) \left[|\psi(k,r)|^2 - |j_0(kr)|^2 \right]$

$C(k) = (a + b\,k)\,(1 + \lambda\,(C_{\text{th}}(k) - 1))$; $\quad S_{12}(\mathbf{r}) = \exp(-r^2/4R^2)/(2\sqrt{\pi}R)^3$

a, b, λ, R ... additional parameters that need to be determined ($\rightarrow$ talk of Yuki Kamiya)

ALICE Collaboration: p-Pb at 5.02 TeV (PRL 123 (2019) 112002) $\quad$ pp at 13 TeV (Nature 588 (2020) 232)

$R = 1.427$ fm; $\lambda = 0.513$ $\quad$ $R = 1.02$ fm; $\lambda = 1$

we adopt $R = 1.427$ fm & 1.18 fm, respectively
(same source radii as found in corresponding fits to pp correlation functions)

(J.H., U.-G. Meißner, arXiv:2201.08238)

Y. Kamiya et al., PRC 105 (2022) 014915, using HAL QCD potential: $R = 1.27$ fm & 1.05 fm
Z.-W. Liu et al., arXiv:2201.04997, cov. χEFT mimicking the HAL QCD potential: $R = 1.427$ fm & 1.182 fm

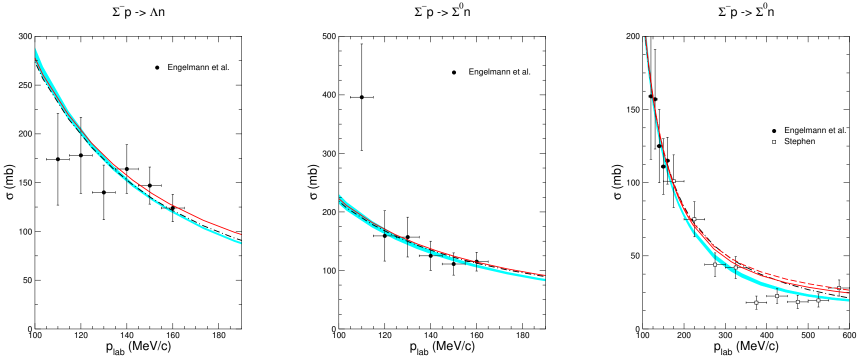
Summary

Hyperon-nucleon interaction within chiral EFT

- ΛN-ΣN interaction within semilocal momentum-space regularized chiral EFT
  confirm our previous YN results (up to NLO) based on a nonlocal regulator
  successful extension to NNLO
  new $\Sigma^{\pm}p$ differential cross sections around $p_{lab} \approx 500$ MeV/c can be described
  unique determination of the P-waves is not yet possible
- Charge symmetry breaking within chiral EFT
  regulator independent results require pertinent contact terms
  CSB splittings in ${}^4_\Lambda\mathrm{He}$-${}^4_\Lambda\mathrm{H}$ ($\Delta E(0^+) = 233 \pm 92$ keV; $\Delta E(1^+) = -83 \pm 94$ keV)
  imply $a_{\Lambda p} - a_{\Lambda n} = 0.62 \pm 0.08$ fm for ${}^1S_0$ state
  however, hypernuclei.kph.uni-mainz.de: 178 ± 55 keV; -139 ± 58 keV
  Elena Botta, HYP2018: 140 ± 120 keV
- ΞN interaction should be fairly weak as suggested by
  the few existing experimental constraints on the ΞN cross sections
  measurements of ΞN two-particle momentum correlations
  lattice QCD simulations close to the physical point
  light Ξ-hypernuclei ($A \geq 4$) could still exist $\rightarrow$ see talk of Hoai Le
- next step: calculate ${}^3_\Lambda\mathrm{H}$, ${}^4_\Lambda\mathrm{He}$, ${}^4_\Lambda\mathrm{H}$, ... with inclusion of three-body forces

backup slides

Preliminary results for SMS YN interactions



SMS YN potentials up to NLO, NNLO (with $\Lambda = 550$ MeV)
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quality of the fit – total χ^2 (36 data points):
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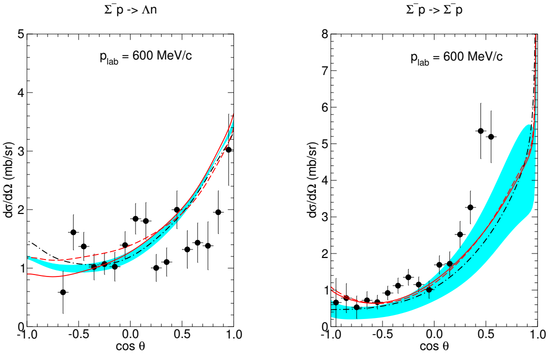
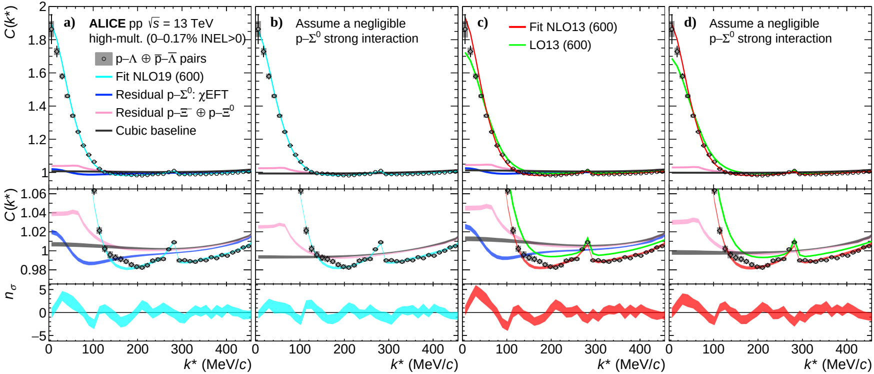


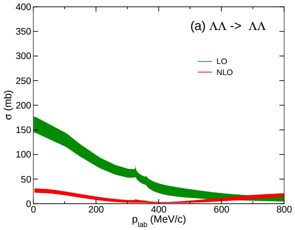
$\Sigma^- p \to \Sigma^- p$: behavior at forward angles remains unclear

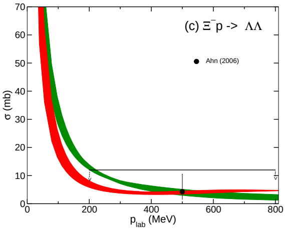
Λp momentum correlation function measured in pp collisions at $\sqrt{s} = 13$ TeV



ALICE Collaboration (Shreyasi Acharya et al.), arXiv:2104.04427

$\Rightarrow$ prediction of NLO19 is fairly well in line with data
sensitive to the assumption about the contribution of the $\Sigma^0 p$ feed-down
"true" Λp amplitude could have slightly weaker energy dependence
(a_t could be about $10 - 15$ % smaller; $\simeq -1.3$ fm instead of $\simeq -1.5$ fm)





ΛΛ effective range parameters

	NLO				LO			
Λ	500	550	600	650	550	600	650	700
$a_{^1S0}$	−0.62	−0.61	−0.66	−0.70	−1.52	−1.52	−1.54	−1.67
$r_{^1S0}$	7.00	6.06	5.05	4.56	0.82	0.59	0.31	0.34

empirical: $a_{\Lambda\Lambda}$ = -1.2 ± 0.6 fm (Gasparyan et al.) $-1.92 < a_{\Lambda\Lambda} < -0.50$ fm (A. Ohnishi et al.)

J.H., U.-G. Meißner, S. Petschauer, NPA 954 (2016) 273

$SU(3)$ structure of contact terms for BB

$SU(3)$ structure for scattering of two octet baryons $\rightarrow$
$8 \otimes 8 = 1 \oplus 8_a \oplus 8_s \oplus 10^* \oplus 10 \oplus 27$

BB interaction can be given in terms of LECs corresponding to the $SU(3)_f$ irreducible representations: C^1, C^{8a}, C^{8s}, C^{10^*}, C^{10}, C^{27}

	Channel	I	V_α	V_β	$V_{\beta\to\alpha}$
$S = 0$	$NN \rightarrow NN$	0	$-$	$C^{10^*}_\beta$	$-$
	$NN \rightarrow NN$	1	C^{27}_α	$-$	$-$
$S = -1$	$\Lambda N \rightarrow \Lambda N$	$\frac{1}{2}$	$\frac{1}{10}\left(9C^{27}_\alpha + C^{8_s}_\alpha\right)$	$\frac{1}{2}\left(C^{8_a}_\beta + C^{10^*}_\beta\right)$	$-C^{8sa}$
	$\Lambda N \rightarrow \Sigma N$	$\frac{1}{2}$	$\frac{3}{10}\left(-C^{27}_\alpha + C^{8_s}_\alpha\right)$	$\frac{1}{2}\left(-C^{8_a}_\beta + C^{10^*}_\beta\right)$	$-3C^{8sa}$
					C^{8sa}
	$\Sigma N \rightarrow \Sigma N$	$\frac{1}{2}$	$\frac{1}{10}\left(C^{27}_\alpha + 9C^{8_s}_\alpha\right)$	$\frac{1}{2}\left(C^{8_a}_\beta + C^{10^*}_\beta\right)$	$3C^{8sa}$
	$\Sigma N \rightarrow \Sigma N$	$\frac{3}{2}$	C^{27}_α	C^{10}_β	$-$

$\alpha = {}^1S_0, {}^3P_0, {}^3P_1, {}^3P_2, \quad \beta = {}^3S_1, {}^3S_1 - {}^3D_1, {}^1P_1$

No. of contact terms: LO: 2 (NN) + 3 (YN) + 1 (YY)
NLO: 7 (NN) + 11 (YN) + 4 (YY)

(No. of spin-isospin channels in NN+YN: 10 $\quad S = -2, -3, -4$: 27)

Contact terms for YN – partial-wave projected

spin-momentum structure up to NLO

$$
\begin{aligned}
V(^1S_0) &= \tilde{C}_{^1S_0} + C_{^1S_0}(p^2 + p'^2) \\
V(^3S_1) &= \tilde{C}_{^3S_1} + C_{^3S_1}(p^2 + p'^2) \\
V(\alpha) &= C_\alpha \, p \, p' \qquad \alpha \hat{=} \, ^1P_1, \; ^3P_0, \; ^3P_1, \; ^3P_2 \\
V(^3D_1 - \, ^3S_1) &= C_{^3S_1 - \, ^3D_1} \, p'^2
\end{aligned}
$$

$$
\begin{aligned}
V(^1P_1 - \, ^3P_1) &= C_{^1P_1 - \, ^3P_1} \, p \, p' \\
V(^3P_1 - \, ^1P_1) &= C_{^3P_1 - \, ^1P_1} \, p \, p'
\end{aligned}
$$

(antisymmetric spin-orbit force: $(\vec{\sigma}_1 - \vec{\sigma}_2) \cdot (\vec{q} \times \vec{k})$)

- $\tilde{C}_\alpha$, C_α ... low-energy constants (LECs)
- need to be fixed by a fit to (NN, YN, ...) data

chiral YN potential up to NNLO

Reinert, Krebs, Epelbaum, EPJA 54 (2018) 86

- Λ: 350 – 550 MeV ... 450 MeV give best results

YN interaction: approximate SU(3) flavor symmetry
$m_\pi = 138$ MeV, $m_K = 495$ MeV, $m_\eta = 547$ MeV

want to keep effects from SU(3) symmetry breaking generated by the single-meson exchange contributions
$\Rightarrow \Lambda$: 500 – 600 MeV

two-meson exchange contributions: πK, $\pi\eta$, ... are represented by contact terms

$\Rightarrow$ some SU(3) symmetry breaking in the YN LECs

(S. Petschauer, N. Kaiser, NPA 916 (2013) 1)

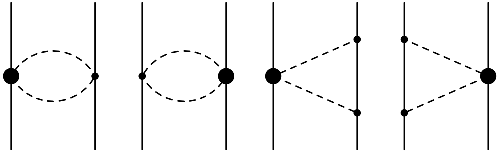
$$V^{cont} = \tilde{C}^\alpha + C^\alpha(p^2 + p'^2) + C^\chi(m_K^2 - m_\pi^2)$$

$\tilde{C}^\alpha$, C^α, $\alpha = \{27\}, \{10^*\}, \{10\}, \{8_s\}, \{8_a\}, \{1\}$, ... “regular” contact terms in SU(3) chiral EFT
C_i^χ: SU(3) symmetry breaking contact terms

(in NLO13 and NLO19 ΛN-ΣN potentials we assumed that $C_i^\chi = 0$)

chiral YN interaction up to NNLO

- no new BB contact terms (no additional LECs) enter
- sub-leading meson-baryon vertices enter at NNLO



πN: fixed from calculating pion-nucleon scattering in chiral perturbation theory
sub-leading (up to Q^2) πN LECs: $c_1 = -0.74$; $c_3 = -3.61$; $c_4 = 2.44$
(cf. RKE 2018)

$\pi\Lambda$, $\pi\Sigma$, $\pi\Lambda \leftrightarrow \pi\Sigma$:
involve additional LECs: d_1, d_2, d_3, b_D, b_F, b_0, b_1, b_2, b_3, b_4
fixed from resonance saturation via decuplet baryons ($\Sigma^*(1385)$)
(cf. Petschauer et al., NPA 957 (2017) 347)

Coupled channels Lippmann-Schwinger Equation

$$
\begin{aligned}
T^{\nu'\nu,J}_{\rho'\rho}(p',p) &= V^{\nu'\nu,J}_{\rho'\rho}(p',p) \\
&+ \sum_{\rho'',\nu''}\int_0^\infty \frac{dp''p''^2}{(2\pi)^3}\, V^{\nu'\nu'',J}_{\rho'\rho''}(p',p'')\,\frac{2\mu_{\rho''}}{p^2-p''^2+i\eta}\, T^{\nu''\nu,J}_{\rho''\rho}(p'',p)
\end{aligned}
$$

ρ', ρ = ΛN, ΣN ($\Lambda\Lambda$, ΞN, $\Lambda\Sigma$, $\Sigma\Sigma$)

LS equation is solved for particle channels (in momentum space)

Coulomb interaction is included via the Vincent-Phatak method

SMS: A nonlocal regulator is applied to the contact terms

$$V^{\nu'\nu,J}_{\rho'\rho}(p',p) \to f^{\Lambda}(p')\,V^{\nu'\nu,J}_{\rho'\rho}(p',p)\,f^{\Lambda}(p); \quad f^{\Lambda}(p) = e^{-(p/\Lambda)^2}$$

consider values $\Lambda = 500$ - 600 MeV [guided by NN, achieved χ^2]

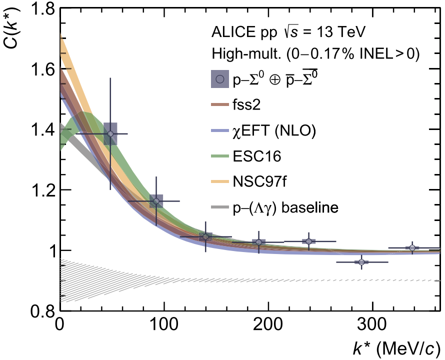
NLO19 (NLO13): A a nonlocal regulator is applied to the whole potential

$$V^{\nu'\nu,J}_{\rho'\rho}(p',p) \to f^{\Lambda}(p')\,V^{\nu'\nu,J}_{\rho'\rho}(p',p)\,f^{\Lambda}(p); \quad f^{\Lambda}(p) = e^{-(p/\Lambda)^4}$$

with values $\Lambda = 500$ - 650 MeV

Femtoscopic studies by ALICE at LHC/CERN

$\Sigma^0 p$ momentum correlation function measured in pp collisions at $\sqrt{s} = 13$ TeV

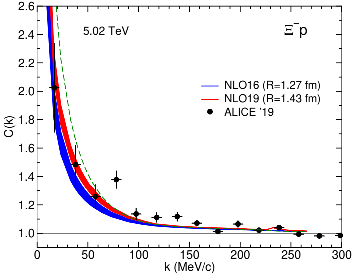


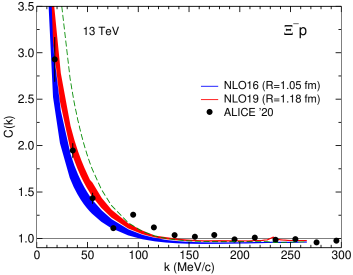
ALICE Collaboration (Shreyasi Acharya et al.), PLB 805 (2020) 135419

open channels ($\Sigma^+ n$, Λp) make theoretical analysis more complicated,

cf. J.H., NPA 981 (2019) 1

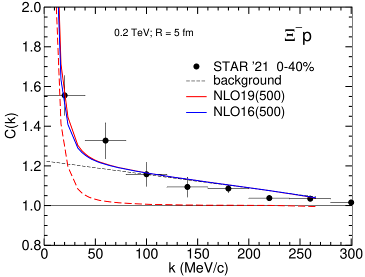
ΞN: two-particle momentum correlation functions

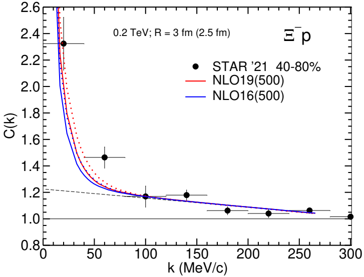




– – – ... result for ΞN interaction that produces $U_\Xi \approx -15$ MeV

Results for $S = -2$: $\Xi^- p$

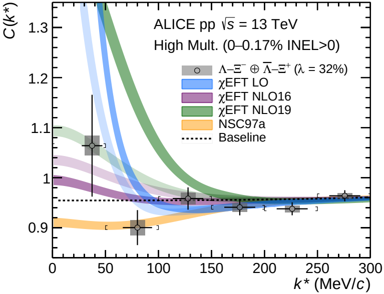


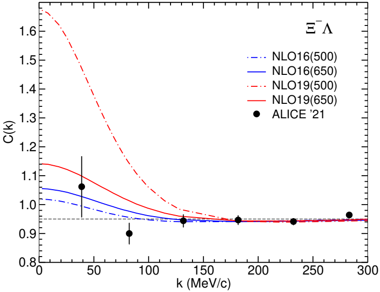


Moe Isshiki (STAR Collaboration) at SQM 2021 (arXiv:2109.10953): Au+Au at 200 GeV

only preliminary results available so far

Results for $S = -3$: $\Lambda\Xi^-$





ALICE Collaboration, arXiv:2204.10258: *pp* at 13 TeV

$R = 1.03$ fm; $\lambda = 0.36$

LO potential (J.H., U.-G. Meißner, PLB 684 (2010) 275):
produces a bound state $\longrightarrow$ not supported by measurement

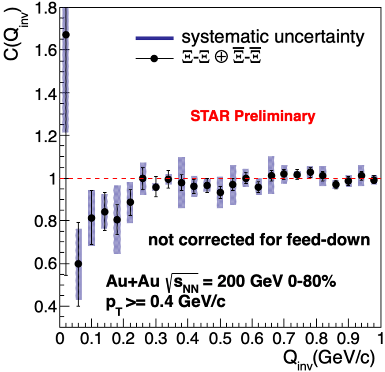
LO rel. χEFT potential (Z.-W. Liu et al., PRC 103 (2021) 025201): likewise too attractive

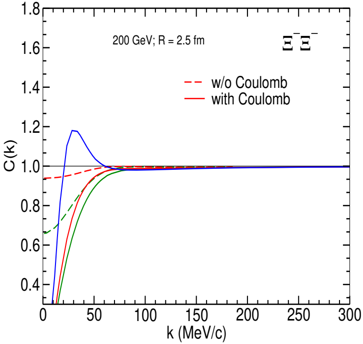
NLO19:
$a_s = -0.99 \cdots - 0.89$ fm, $r_s = 4.63 \cdots 5.77$ fm; $a_t = -0.42 \cdots -1.66$ fm, $r_t = 6.33 \cdots 1.49$ fm

NLO16:
$a_s = -0.99 \cdots - 0.89$ fm, $r_s = 4.63 \cdots 5.77$ fm; $a_t = 0.026 \cdots -0.12$ fm, $r_t = 32.0 \cdots 702$ fm

(J.H., U.-G. Meißner, arXiv:2201.08238)

Results for $S = -4$: $\Xi^-\Xi^-$





Moe Isshiki (STAR Collaboration) at SQM 2021 (arXiv:2109.10953): Au+Au at 200 GeV

only preliminary results: $R = 2.5 - 5$ fm; $\lambda = ??$

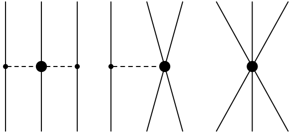
use for calculation: $R = 2.5$ fm; $\lambda = 1$

$a_s = -7.04$ fm (no SU(3) breaking) $\quad -1.71$ fm (moderate SU(3) breaking) $\quad -0.71$ fm (strong SU(3) breaking)

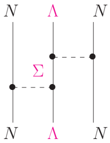
Three-body forces

- SU(3) χEFT 3BFs at NNLO (S. Petschauer et al., PRC 93 (2016) 014001)
- however, 5 LECs for ΛNN 3BF alone! (only 2 LECs for NNN)

SU(3) χEFT :

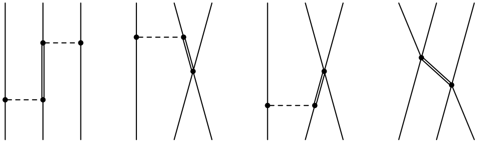


(at LO!) :



solve coupled channel (ΛN-ΣN) Faddeev-Yakubovsky equations:
$\Rightarrow$ ΛNN “3BF” from Σ coupling is automatically included

- 3BFs with inclusion of decuplet baryons (S. Petschauer et al., NPA 957 (2017) 347)



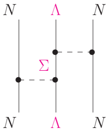
estimate ΛNN 3BF based on the $\Sigma^*(1385)$ excitation (appear at NLO!)

- only 1 LEC for ΛNN (2 LECs for YNN in general)

Estimation of 3BFs based on NLO results

- ${}^{3}_{\Lambda}\mathrm{H}$
  (a) cutoff variation: ΔE_Λ (3BF) $\leq$ 50 keV
  (b) "3BF" from ΛN-ΣN coupling:

  switch off ΛN-ΣN coupling
  in Faddeev-Yakubovsky equations:
  ΔE_Λ (3BF) $\approx$ 10 keV
  expect similar/smaller ΔE_Λ from $\Sigma^*(1385)$ excitation



  (c) ${}^{3}\mathrm{H}$: 3NF $\sim Q^3\, |\langle V_{NN}\rangle|_{{}^3\mathrm{H}} \sim$ 650 keV
  ($|\langle V_{NN}\rangle|_{{}^3\mathrm{H}} \sim$ 50 MeV; $Q \sim m_\pi/\Lambda_b$; $\Lambda_b \simeq$ 600 MeV)
  ${}^{3}_{\Lambda}\mathrm{H}$: $|\langle V_{\Lambda N}\rangle|_{{}^3_\Lambda\mathrm{H}} \sim$ 3 MeV $\rightarrow \Delta E_\Lambda$ (3BF) $\approx Q^3\, |\langle V_{\Lambda N}\rangle|_{{}^3_\Lambda\mathrm{H}} \simeq$ 40 keV

- ${}^{4}_{\Lambda}\mathrm{H}$, ${}^{4}_{\Lambda}\mathrm{He}$
  (a) cutoff variation: ΔE_Λ (3BF) $\approx$ 200 keV (0^+) and $\approx$ 300 keV (1^+)
  (b) "3BF" from ΛN-ΣN coupling:
  ΔE_Λ (3BF) $\approx$ 230 – 340 keV (0^+), $\approx$ 150 – 180 keV (1^+)

${}^{3}_{\Lambda}\mathrm{H}$ and ${}^{4}_{\Lambda}\mathrm{H}$(He) calculations with explicit inclusion of 3BFs utilizing the decuplet saturation are planned for the future