

# Nuclear Lattice Effective Field Theory: Status A.D. 2024

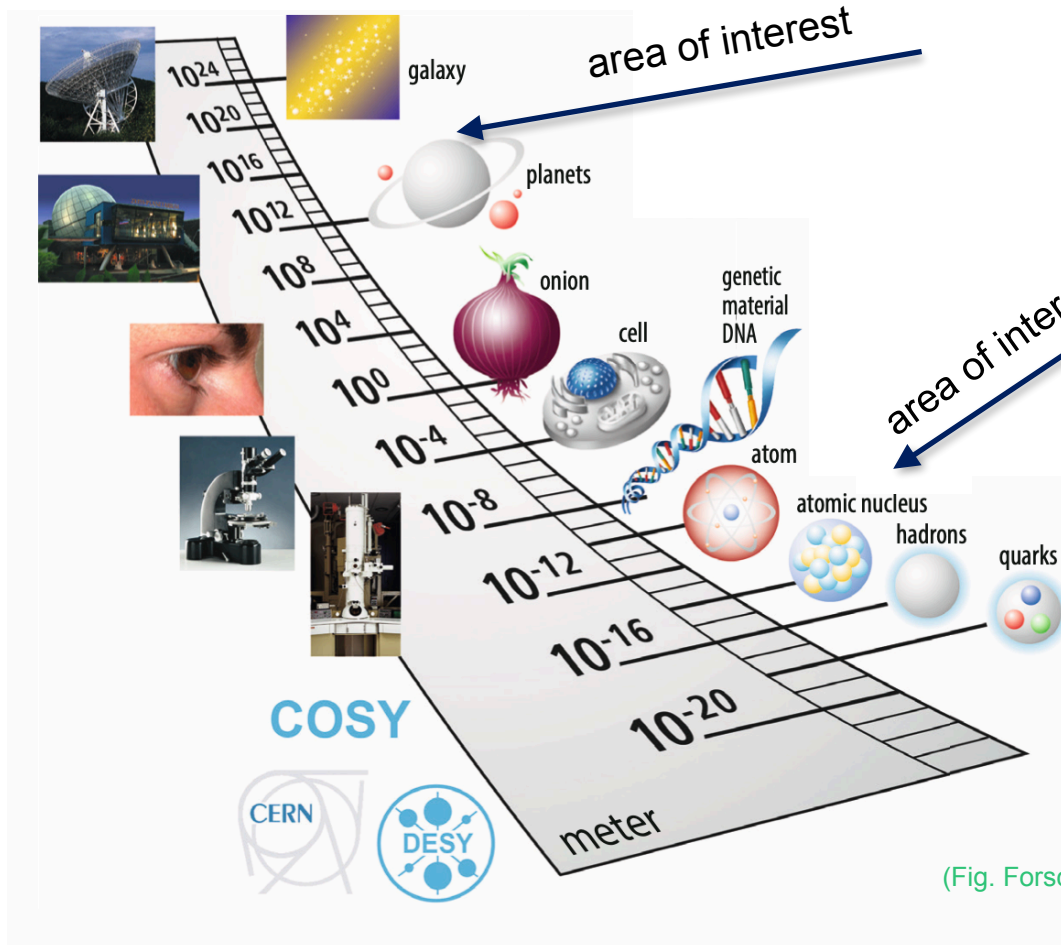
Fabian Hildenbrand, IAS-4, Forschungszentrum Jülich, Germany  
In collaboration with Ulf-G. Meißner and NLEFT

## Outline

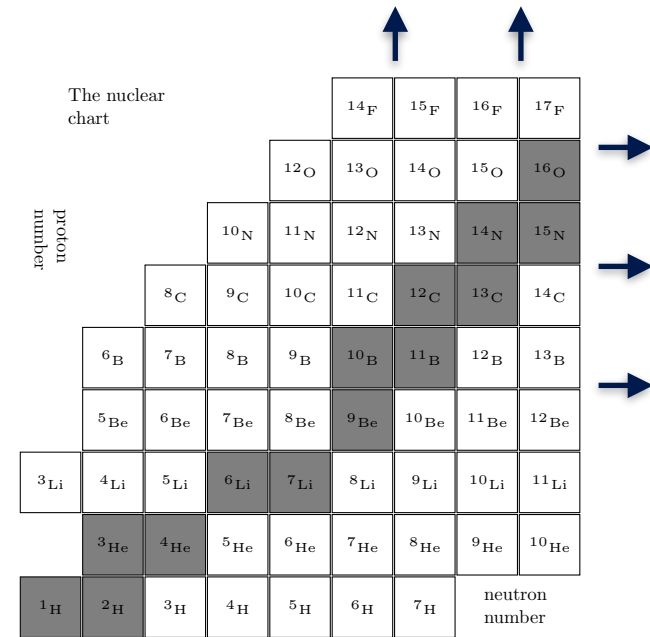
- Motivation
- ▶ Introduction to NLEFT
- ▶ Wavefunction Matching
- ▶ Neutron Stars
- Summary and Outlook

# What is Nuclear Lattice Effective Field Theory?

- Goal: understanding nuclear structure and reactions from small systems to stars



(Fig. Forschungszentrum Jülich)



Open questions:

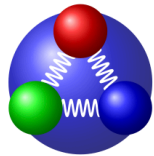
- Limits of stability
- Three-body forces
- Cluster structures
- Neutron stars

- At different scales, there are different relevant degrees of freedom
- Resolution matters: exploit this separation of scales

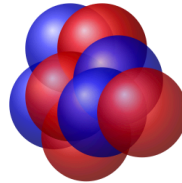
# Dominant interactions at different scales



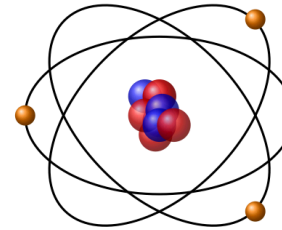
Quarks  
 $< 10^{-16}$



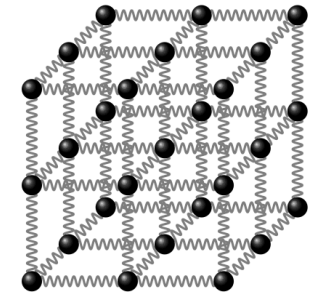
Nucleons  
 $\sim 10^{-13}$



Nucleus  
 $\sim 10^{-12}$



Atoms  
 $\sim 10^{-8}$  cm



Matter

(Fig. S.Elhatisari)



strong force



nuclear force



electromagnetic force  
 sometimes gravity

Utilize this separation! There is no need to consider things that are not relevant at a specific scale

# Chiral Nuclear Effective Field Theory

- Pick the correct scale :  $\sim 10^{-15}$  m = 1 fermi
- Pick the relevant degrees of freedom : protons and neutrons and pions

*This is not Lattice QCD,  
no quarks or gluons!*



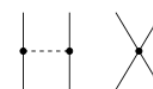
construct effective field theory

- Low energy chiral effective field theory of QCD
  - No model, systematic improvement is possible due to a hierarchy of forces
  - Systematic error analysis possible
  - Typical nuclear systems are far away from breakdown scale

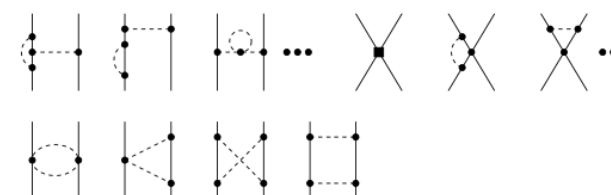


combine with lattice methods

Leading order



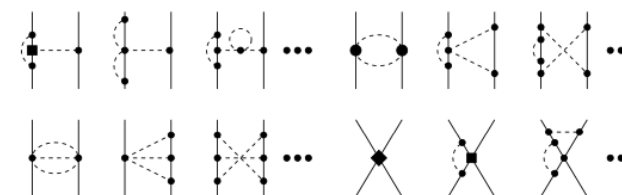
Next-to-leading order



Next-to-next-to-leading order



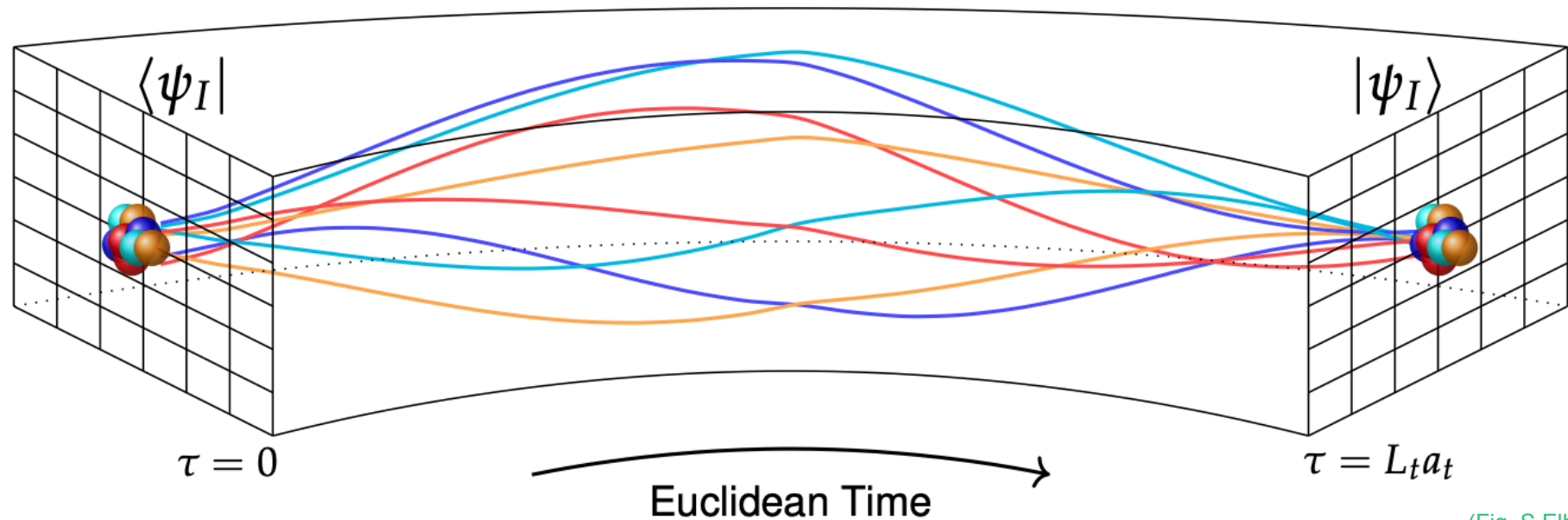
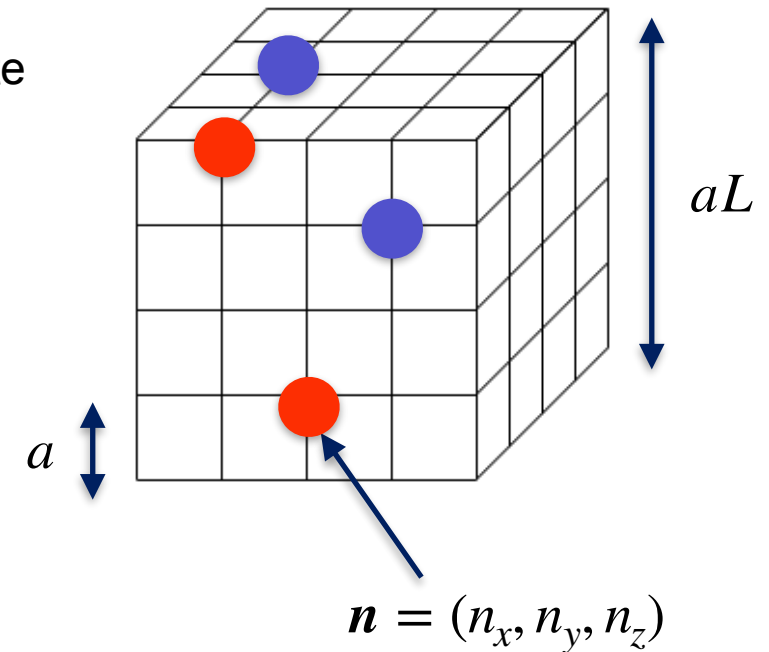
Next-to-next-to-next-to-leading order



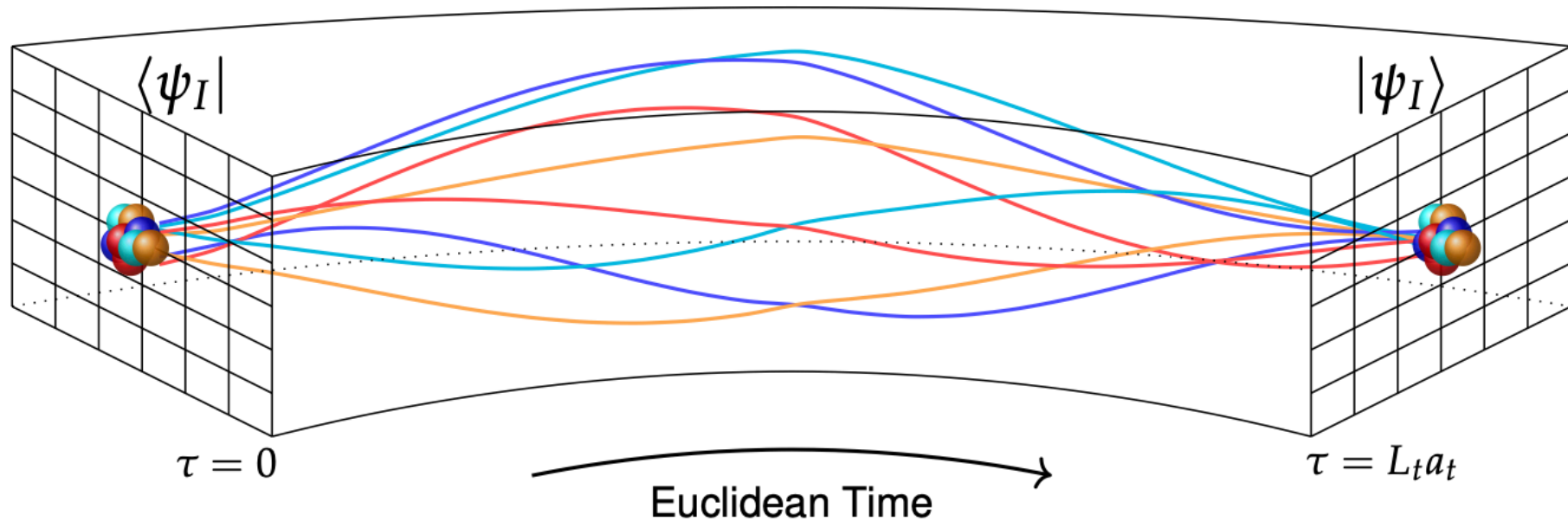
(Epelbaum et al. RevModPhys.81.1773)

# Method: Lattice Monte Carlo

- Lattice  $\Rightarrow$  cubic volume of size  $(La)^3$  with discrete lattice site
- Discretized chiral potentials, contact interactions  
one-pion exchange, Coulomb (Epelbaum et al.)
- Do Euclidean time evolution and extract i.e. energies  
as transient energy  $E = -\frac{d}{d\tau} \ln(Z(\tau))$



(Fig. S.Elhatisari)



(Fig. S.Elhatisari)

- Auxiliary fields to handle many particles efficiently:
- Idea: replace interactions between nucleons with interaction of a nucleon with an auxiliary field

$$\exp\left(-\frac{C}{2}(N^\dagger N)^2\right) = \sqrt{\frac{1}{2}} \int dA \exp\left[-\frac{A^2}{2} + \sqrt{CA}(N^\dagger N)\right]$$

Since nucleons only interact with an auxiliary field  $\Rightarrow$  perfect for parallel computing

# The Challenge : Sign Problem

- Sign problem in a nutshell : makes life hard!

fermionic wave functions change sign when two fermions are interchanged. The systems are strongly interacting



highly oscillatory function, huge cancellation effects

*no reasonable computation at higher orders possible*

- First way out: utilize approximate SU(4) Wigner symmetry

(E. Wigner Phys.Rev 51(1937))

treats protons and neutrons on equal footing



same particle with different quantum number isospin



create SU(4) invariant interactions, calculations are almost sign free



*some aspects of nuclear structure are well described*

(S. Shen et al. Nature Commun.14 (2023))

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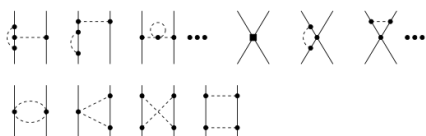
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but we want and need:

Leading order



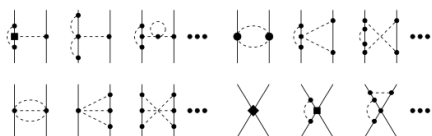
Next-to-leading order



Next-to-next-to-leading order



Next-to-next-to-next-to-leading order



Need improved perturbation theory to go beyond N2LO



create SU(4) invariant interactions, calculations are almost sign free



some aspects of nuclear structure are well described

(S. Shen et al. Nature Commun.14 (2023))



# Wavefunction Matching!

(S.Elhatisari et al. Nature 630(2024))

- New method to solve the quantum many-body problem
- Not only applicable in our field but maybe also in yours!

starting point :  $H_S$  which has acceptable sign problem and can be simulated



How to connect?

target :  $H$  which has a severe sign problem



lowest eigenstate  $|\psi_S^0\rangle$

projected and normalised :

$$|\phi_S^0\rangle = \mathcal{P} |\psi_S^0\rangle / ||\psi_S^0||$$

map  $|\psi_S^0\rangle$  on  $|\psi^0\rangle$



lowest eigenstate  $|\psi^0\rangle$

projected and normalised :

$$|\phi^0\rangle = \mathcal{P} |\psi^0\rangle / ||\psi^0||$$

$$U = \begin{cases} 1 & r > R \\ |\phi^\perp\rangle\langle\phi_S^\perp| & r < R \end{cases}$$

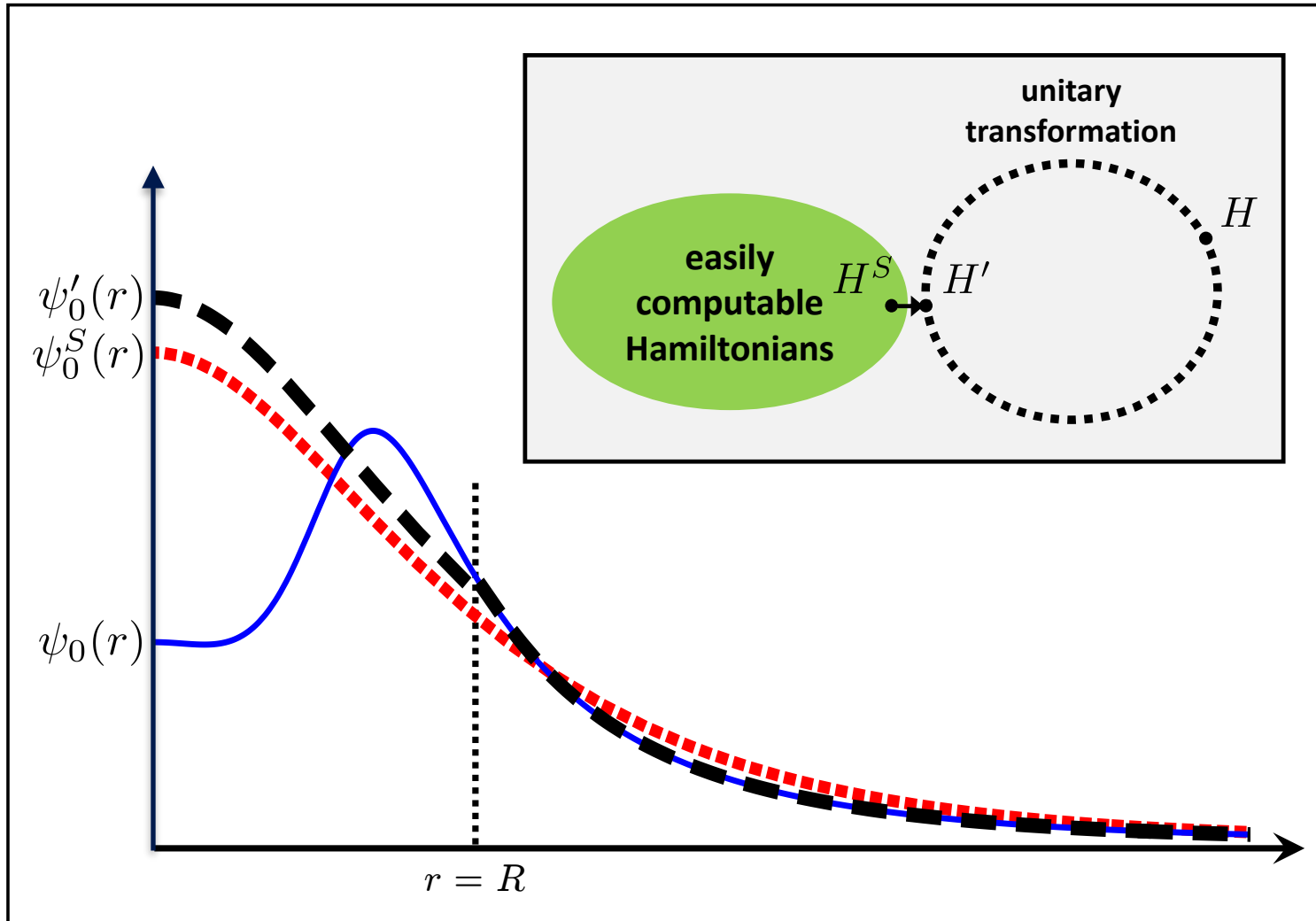
Unitary trafo with finite range



$$H' = U^\dagger H U$$

perturbation theory now works from  $H_S$  to  $H'$

# Wavefunction Matching II !

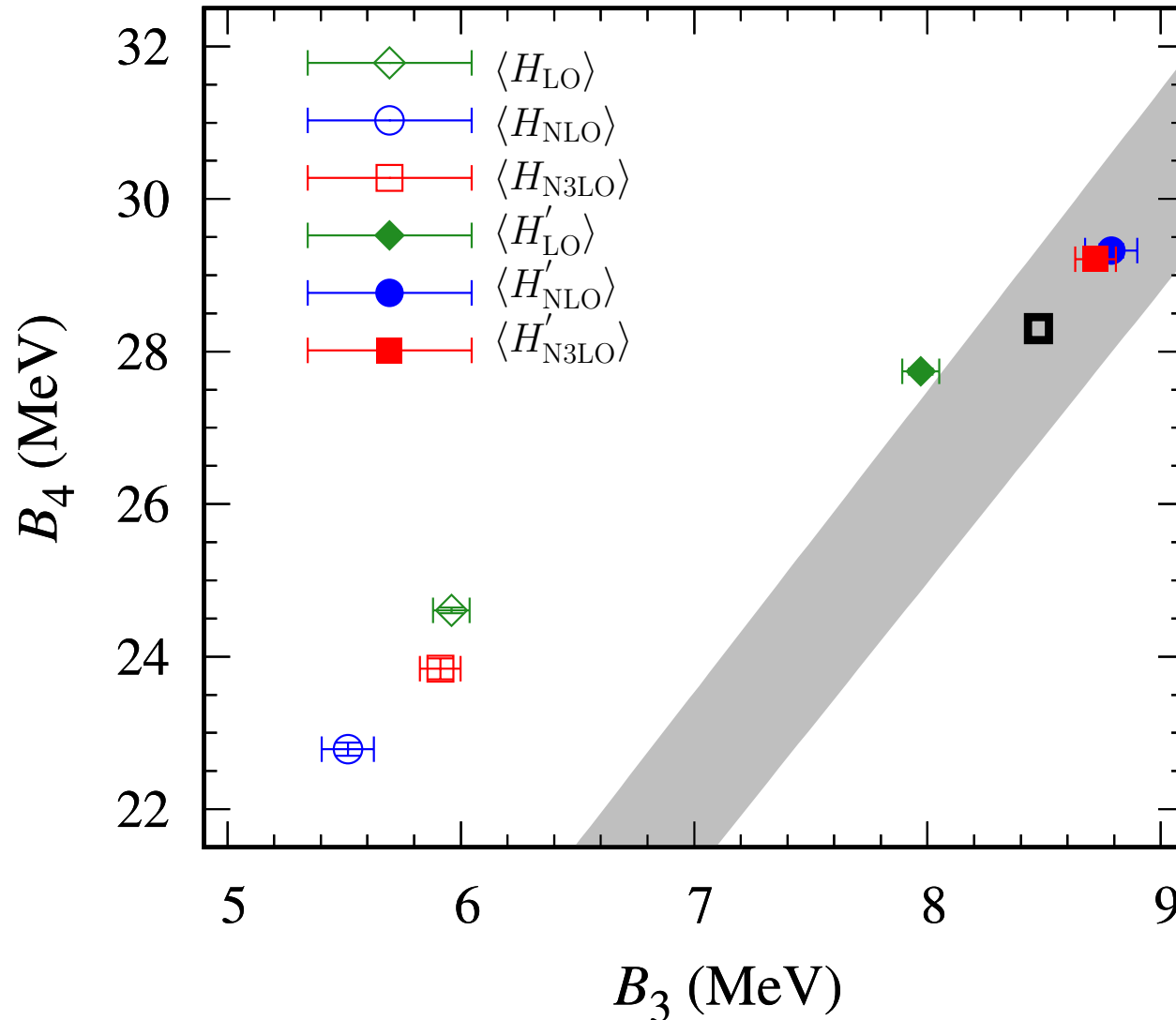


- Use  $H^S$  for non-perturbative part and do perturbation theory to  $H'$
- Direct use of  $H'$ , no reconstruction of higher order forces needed!

This is no renormalization group transformation

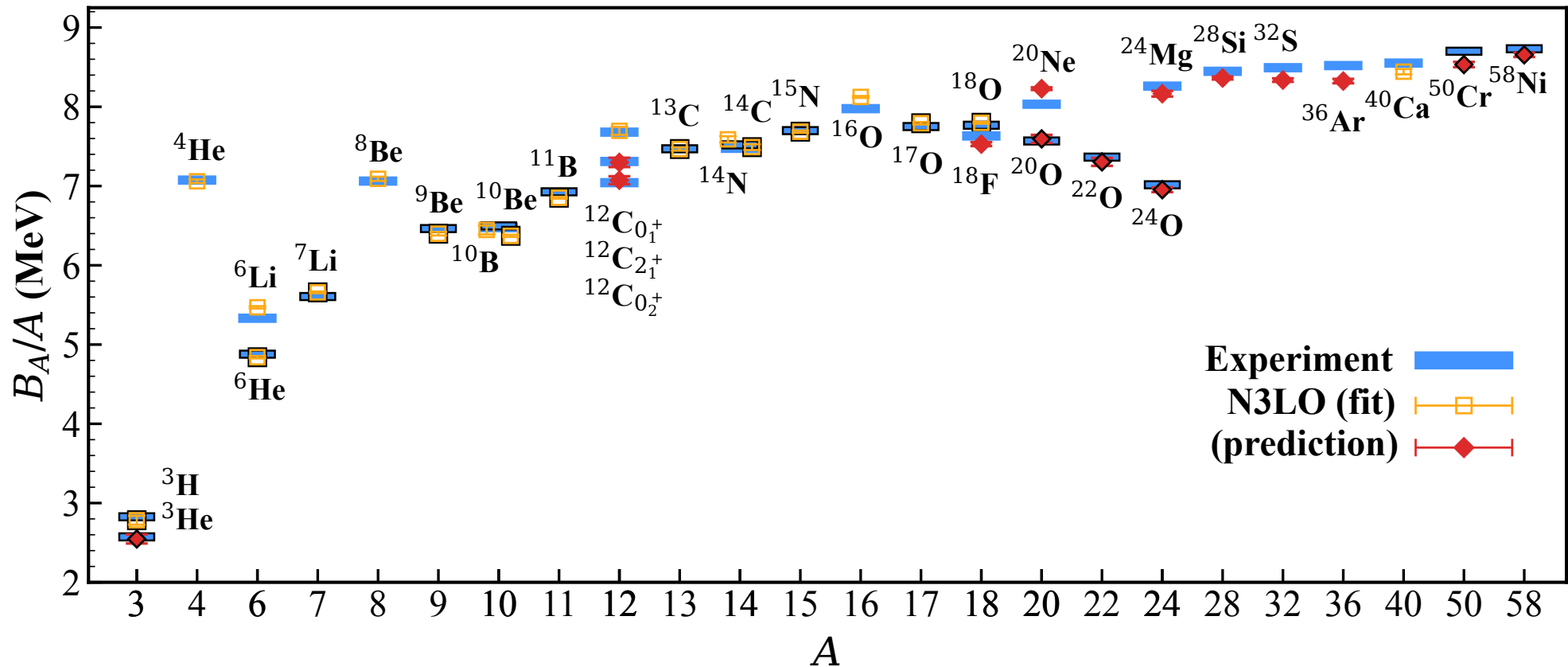
# Does it Work in Practice? (Tjon Band)

- Universal correlation between 3- and 4-body binding energy (L.Platter et al. PLB 607(2005))
- Tjon band is reproduced using Wavefunction Matching



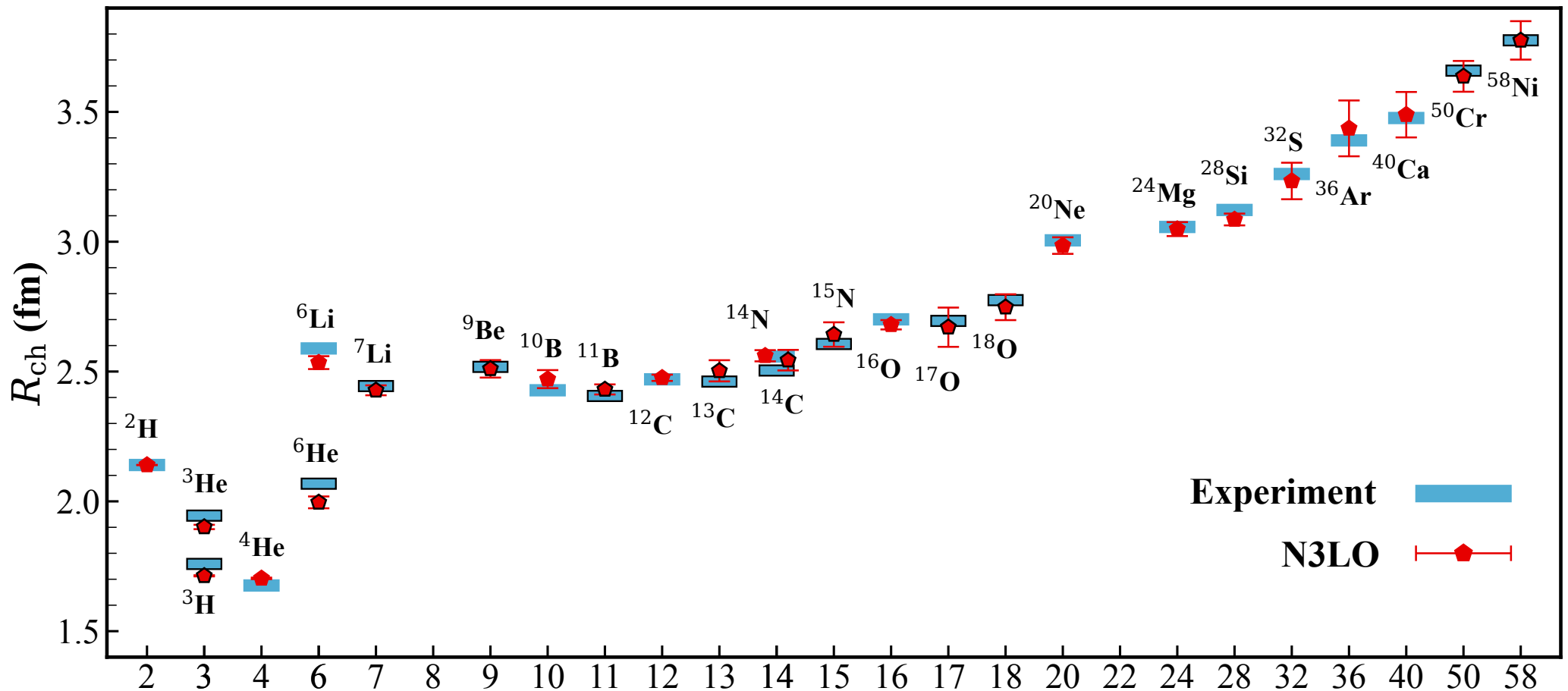
# Binding Energy up to Medium Mass Nuclei

- Fix 3NF forces using history matching, systematic errors accessible
- Ground and excited states are well reproduced over a large excerpt of the nuclear chart
- $a = 1.32$  fm



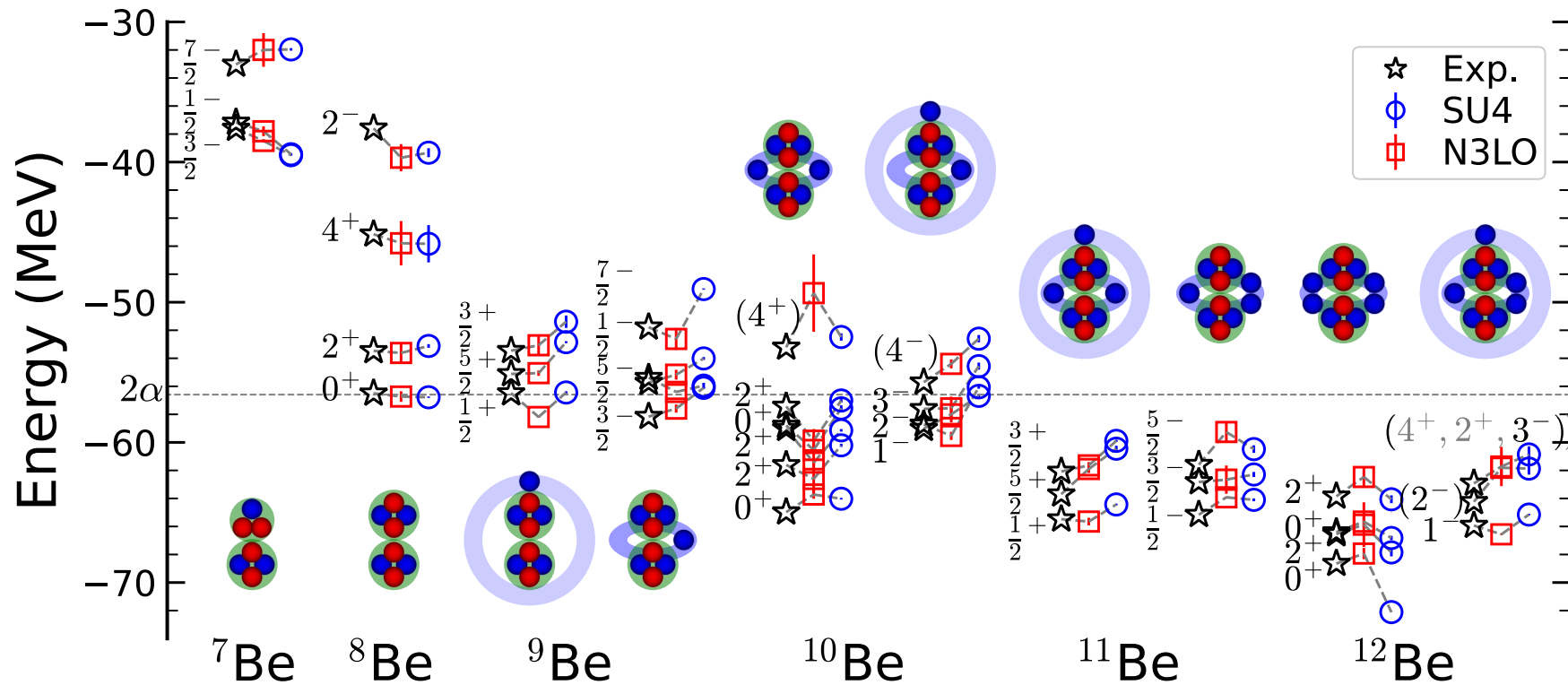
# Predictions for Charge Radii

- Charge radii are spot on!
- No fitting to charge radii, everything is a prediction!
- $a = 1.32$  fm, statistical errors can be reduced



# In Detail Study of Beryllium Isotopes

- Very different structures can be described within one calculation
- SU(4) interaction works quite well
- N3LO improves the result further



(S. Shen et al. arxiv: 2411.14935 (2024))

# A very brief Introduction to Neutron Stars

- Compact remnants of stars at the end of their life
- Typical quantities:

$$r \sim 10 \text{ km}$$

$$M \sim (1 - 2)M_{\odot}$$

limited by Tolman-Oppenheimer-Volkoff equations

(R. Tolman Phys. Rev. 55(1939))

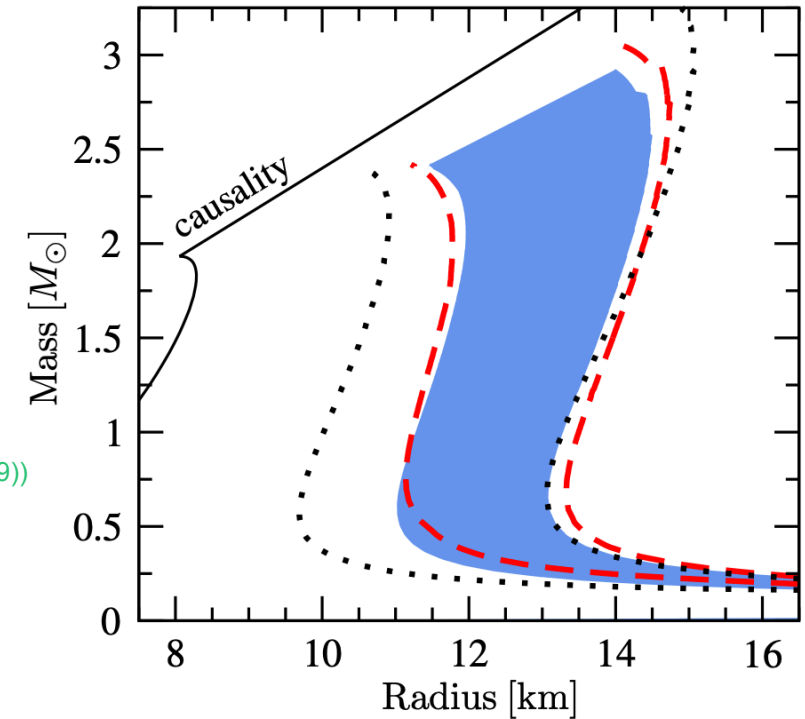
(J. Oppenheimer, G. Volkoff Phys. Rev. 55(1939))

super dense objects,  
 $\rho$  several times saturation  
density in the core

Mass-Radius relation,  
if EoS correct, every  
neutron star falls on this  
line

Open questions:

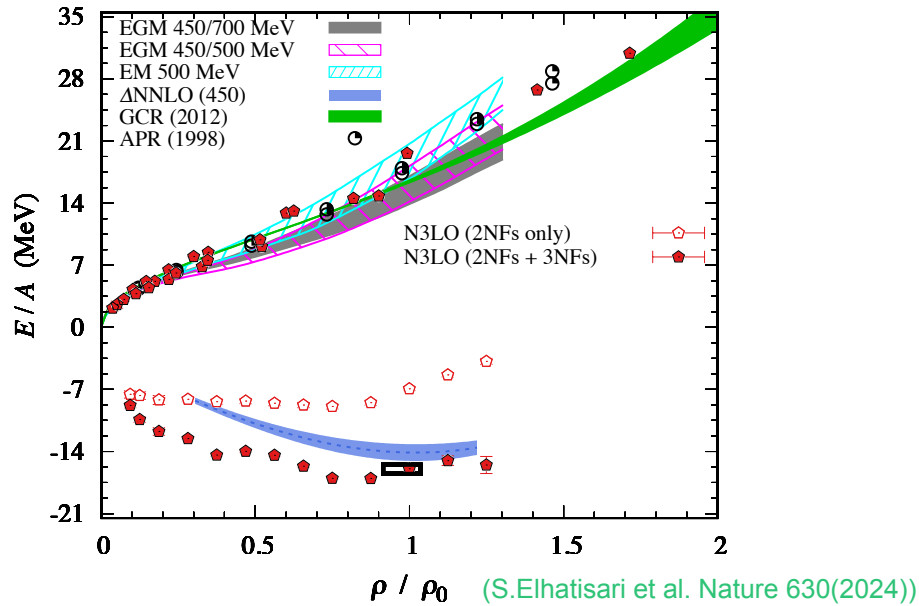
- density in the centre
- equation of state (EoS)
- inclusion of other baryons



(Hebeler et al. Astrophys.J. 773(2013))

extract this from NLEFT,  
check inclusions of other  
baryons as well

# EoS of pure neutron and symmetric nuclear matter



addition of hyperons, typically softens the equation of states



neutrons stars are no longer supported



naiv assumptions make it reasonable for hyperons to appear at such densities



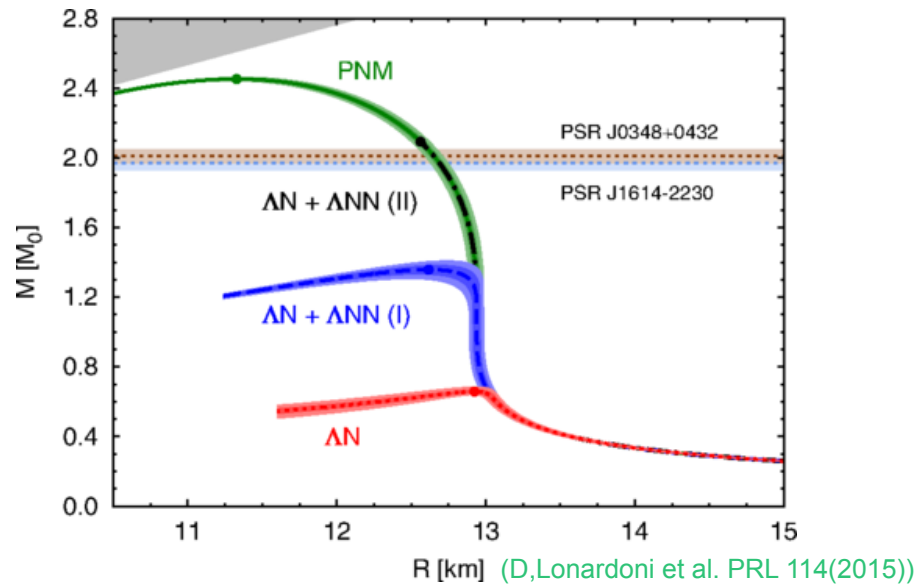
Hyperon Puzzle



resolved by three-body forces

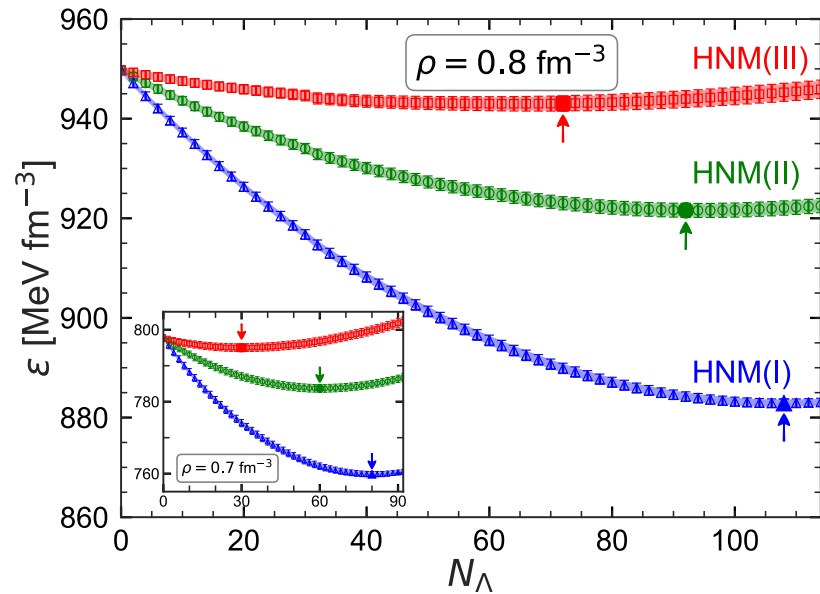


no ab initio calculation with large amount of hyperons



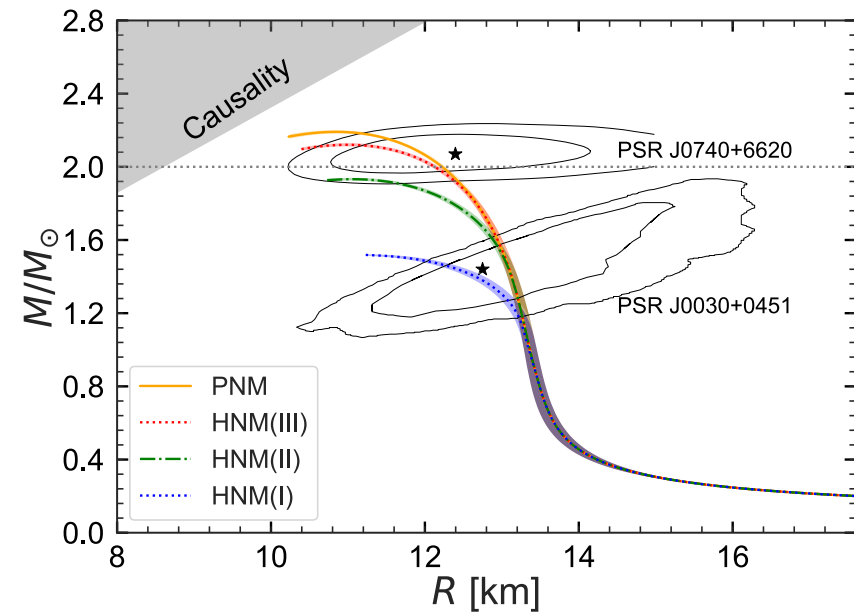
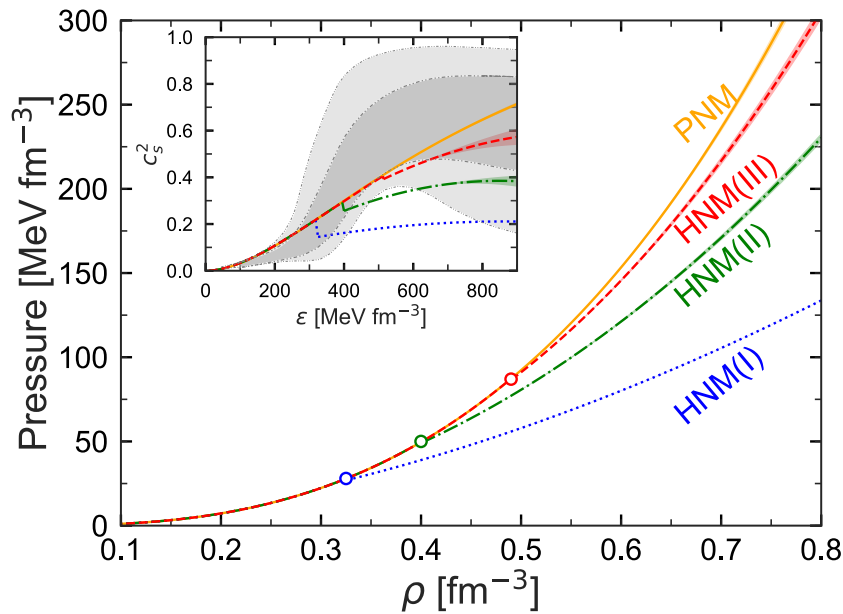


# Hypernuclear matter from NLEFT



calculation with  $\sim 240$  particles and up to 50% hyperons

we explore a large set of different configurations up to 5 times saturation density



(H.Tong et al. Sci. Bull. (in press)  
arxiv:2405.01887(2024))

- Give you a first insight in Nuclear Lattice Effective Field Theory
  - Relevant degrees of freedoms are neutrons and protons
  - Combine effective field theory with lattice methods
- Presented a new approach to solve the quantum many-body problem:
  - Wavefunction Matching
  - Makes a whole set of new calculations available
  - High precision consistent calculation of matter radii and binding energies
- Showed how to connect  $\sim 200$  particles with  $10^{69}$  at the example of neutron stars
  - First ab initio calculation using a significant number of  $\Lambda$ s
  - EoS including  $\Lambda$  particles which is consistent with nature

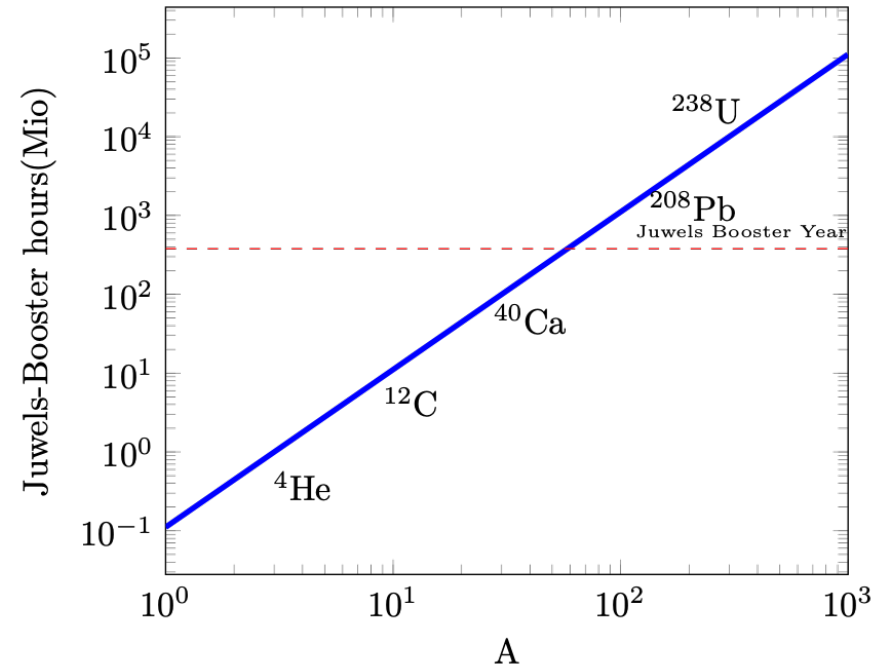
- JUPITER changes the game:

$^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$  is now in reach



critical to understand the carbon/oxygen ratio in the universe

Expected Compute time needed for  $\alpha - A$  scattering



the Sn line:



(Nudat3(2025))

typical benchmark to go to even heavier nuclei



different structures accessible