

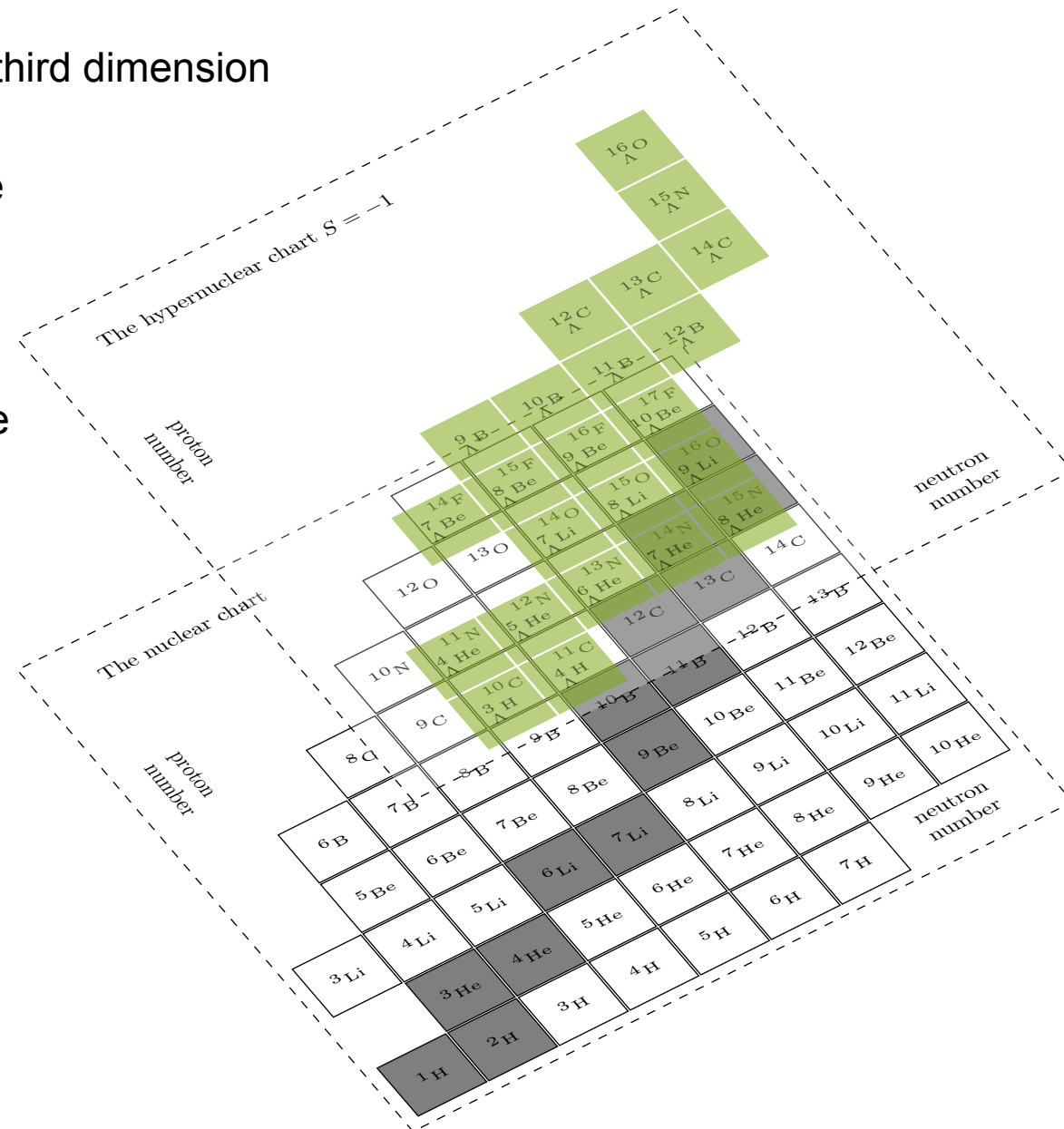
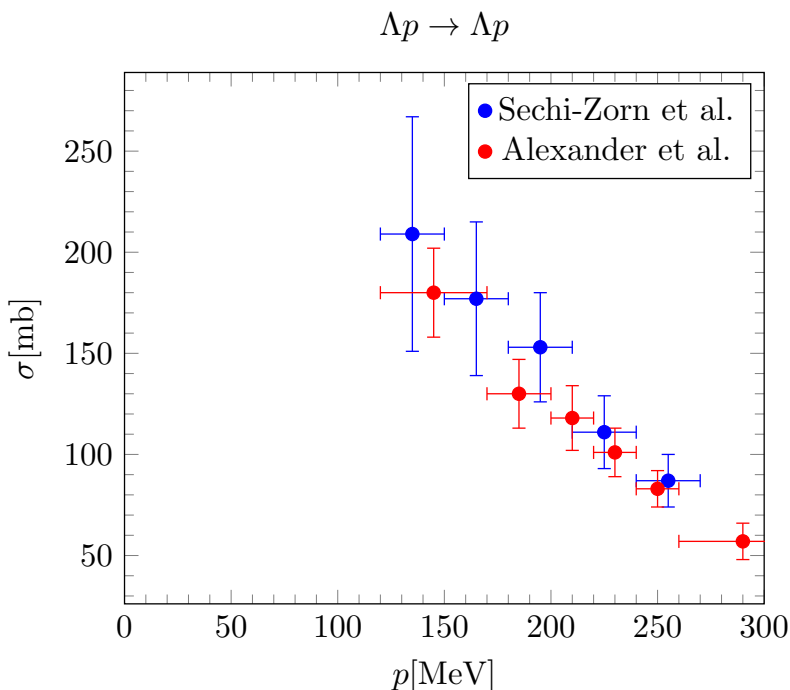
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Outline

- Motivation
- ▶ From NLEFT to (Hyper) NLEFT
 - ▶ Lattice Interaction
 - ▶ Two-body results
 - ▶ Inclusion of three-body forces
- Summary and Outlook

Hypernuclear physics in a nutshell

- Strangeness extends the nuclear chart to a third dimension
- Unique opportunity to study the strong force without the Pauli principle
- Typical approach from nuclear physics does not work since two-body data is sparse



- Gateway : **Three-Body Systems**

Starting point for (Hyper) Nuclear Lattice EFT

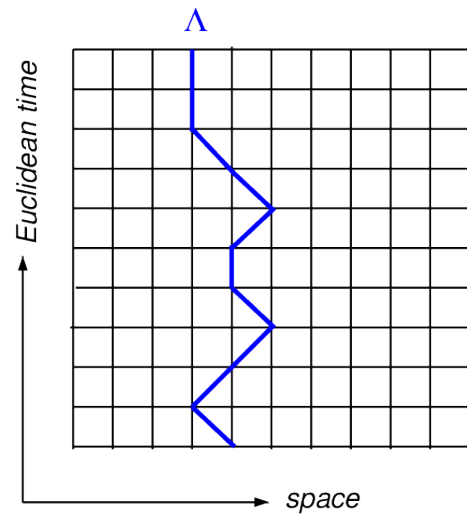
very successful nuclear program:
 using AFMC and shuttle algorithm
 wave function matching to obtain
 precise results for nuclei and
 charge radii

AFMC does not converge as good as in a
 pure nuclear matter simulation

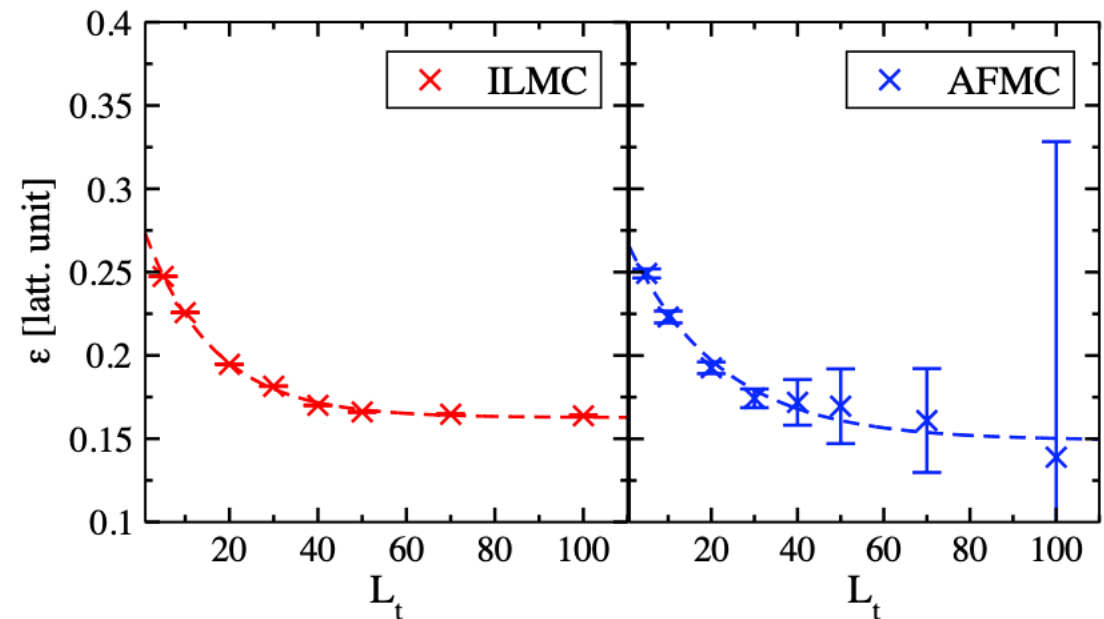
Need to develop a method that treats this
 impurities more efficient

Treat impurity as worldline:

(S.Bour, D.Lee, H.-W. Hammer, U.-G. Meißner)



(D. Frame, T.A. Lähde, D. Lee, U.-G. Meißner)



Starting point for (Hyper) Nuclear Lattice EFT

- Challenge with IFMC, need to collect millions of worldlines

→ Can we still do hypernuclear calculations with AFMC ?

→ Important for possible applications with many hyperons

- Taylor interaction to work non-perturbative with our best NN interaction

→ Evolve together with NN counterparts
 Constraints smearing parameters to the NN ones

$A = 3$ 0.97

$A = 4$ 0.89

$A = 5$ 1 ← α - core

$A = 7$ 0.92

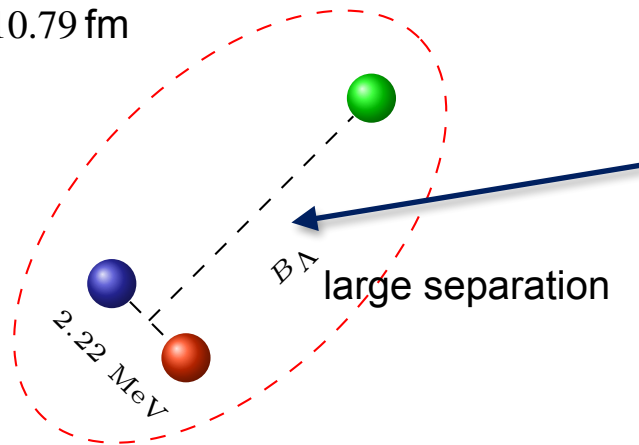
$A = 13$ 0.97

$L = 12$ $Lt = 500$

this is very promising,
for larger hypernuclei

Construction of a first Lattice ΛN interaction

$$\sqrt{\langle r_{\Lambda d}^2 \rangle} \approx 10.79 \text{ fm}$$



Emulsion:

$$B_\Lambda = 0.130 \pm 0.050 \text{ MeV} \text{ Juric 1973}$$

Heavy Ion:

$$B_\Lambda = 0.406 \pm 0.120 \text{ MeV} \text{ Star 2020}$$

$$B_\Lambda = 0.102 \pm 0.063 \text{ MeV} \text{ Alice 2023}$$

World Average:

$$B_\Lambda = 0.164 \pm 0.043 \text{ MeV} \text{ Mainz 2025}$$

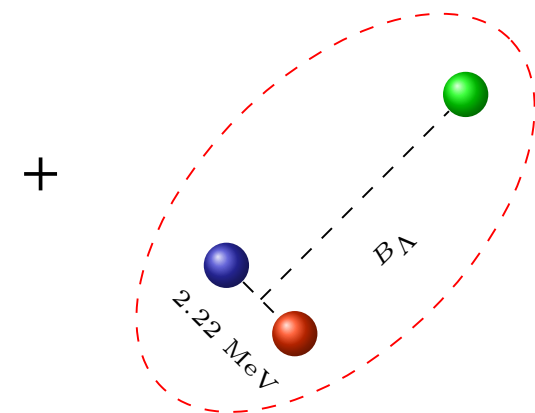
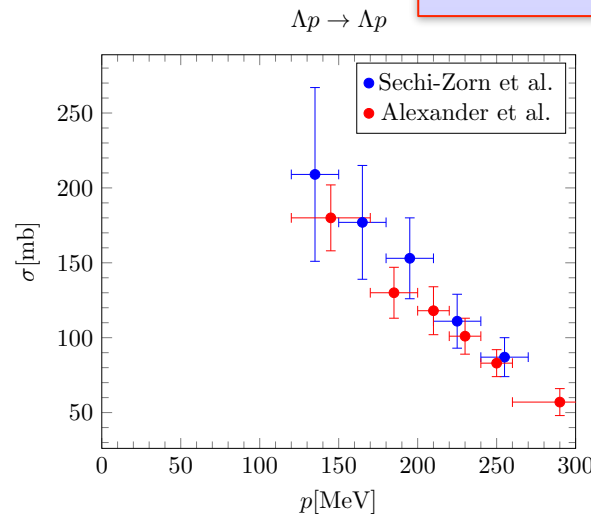
Shallow S-Wave State

$$J^P = \frac{1}{2}^+$$

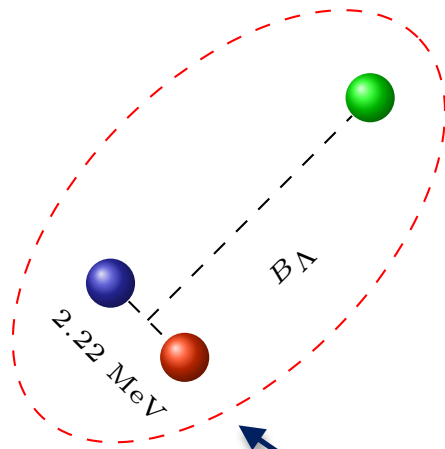
Distinguishable

$$I = 0 \Rightarrow \frac{1}{\sqrt{2}} (pn - np) \Lambda$$

Combine 2-Body data with hypertriton in exact calculation



Construction of a first Lattice Λ N interaction



—————>
large separation

Large box sizes needed

GPU Lanczos Code to fix 2-body forces

—————> Allows $L = 23.76$ fm

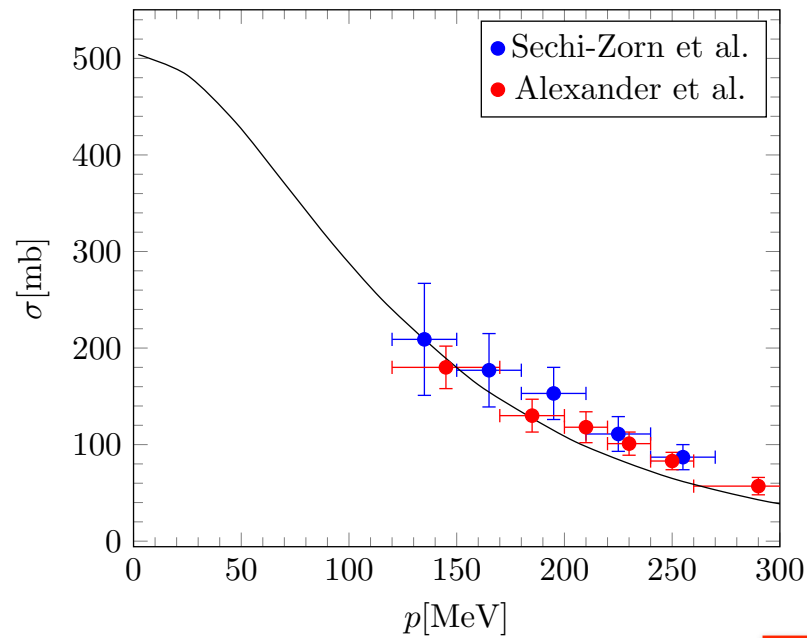
Tightly bound compared to Λ separation

Use simpler nuclear interaction

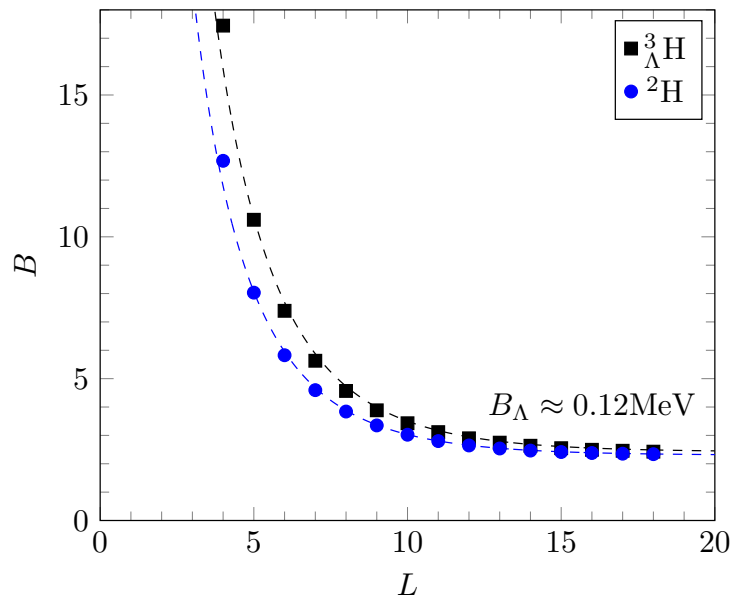
Non-perturbative LO interaction $\begin{cases} \rightarrow {}^3S_1 \\ \rightarrow {}^1S_0 \end{cases}$ interaction

Construction of a first Lattice ΛN interaction

$\Lambda p \rightarrow \Lambda p$



Lanczos Three-body Result



Best SMS N^2LO interaction

(Haidenbauer et al.)

$$a_s = -2.80 \text{ fm} \quad r_s = 2.89 \text{ fm}$$

$$a_t = -1.58 \text{ fm} \quad r_t = 3.09 \text{ fm}$$

This interaction

$$a_s = -2.89 \text{ fm} \quad r_s = 3.28 \text{ fm}$$

$$a_t = -1.60 \text{ fm} \quad r_t = 3.94 \text{ fm}$$

Phase shift similar to $p \sim 60$ MeV

$$E(L) = E_{L \rightarrow \infty} + \frac{A}{L} e^{-\frac{L}{L_0}} \approx \text{Emulsion}$$

$$B_{L \rightarrow \infty}^{\Lambda} = (90 + 30) \text{ keV} \approx 120 \text{ keV}$$

2-Body

GIR corrections

This is different from the previous talk

Results: Two Body interaction (L=12 I.u.) (light nuclei)

During Evolution:

Spin-averaged Interaction:

$$C = \frac{3 \ ^3S_1 + \ ^1S_0}{4}$$

Perturbative part:

Spin-dependent Interaction:

$$C_S = \frac{\ ^3S_1 - \ ^1S_0}{4}$$

Nuclear Interaction:

N^3 LO interaction, same as for WFM results

Results: Two Body

Experiment

$$B_{\Lambda}(\ ^3\text{H}) = 0.38 \pm 0.08 \text{ MeV}$$

$$0.164 \pm 0.43 \text{ MeV}$$

→ Box effect, consistent with exact L=12 result

$$B_{\Lambda}(\ ^4\text{H}^{0+}) = 2.11 \pm 0.18 \text{ MeV}$$

$$2.169 \pm 0.042 \text{ MeV}$$

$$B_{\Lambda}(\ ^4\text{H}^{1+}) = 1.23 \pm 0.18 \text{ MeV}$$

$$1.081 \pm 0.042 \text{ MeV}$$

→ Splitting quite good, missing 0.2 MeV

$$B_{\Lambda}(\ ^5\text{He}) = 3.51 \pm 0.12 \text{ MeV}$$

$$3.102 \pm 0.03 \text{ MeV}$$

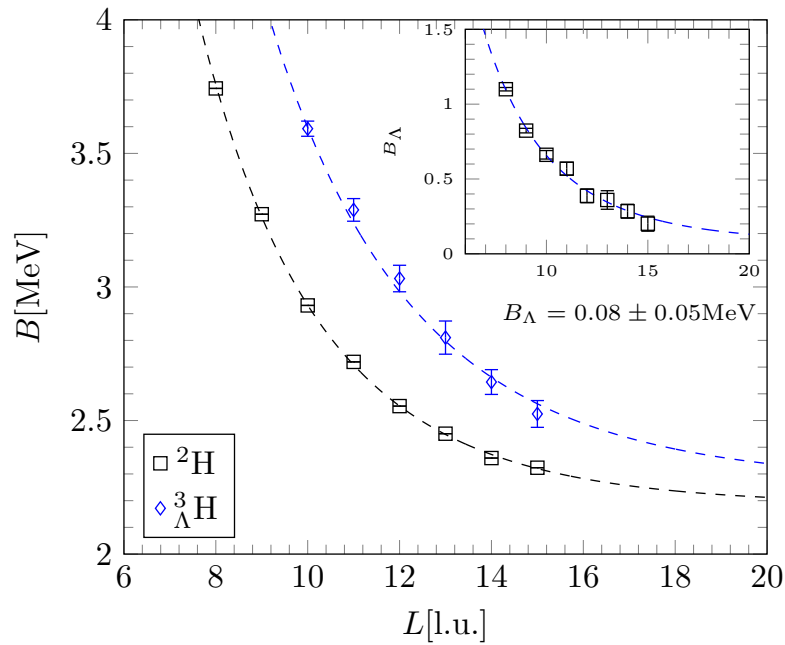
→ Smaller overbinding compared to other LO calculations

$$B_{\Lambda}(\ ^7\text{Li}) = 5.68 \pm 0.96 \text{ MeV}$$

$$5.619 \pm 0.06 \text{ MeV}$$

→ Typically overbound by ~1 MeV in LO calculations

Box Size effects:



Proper extrapolation:

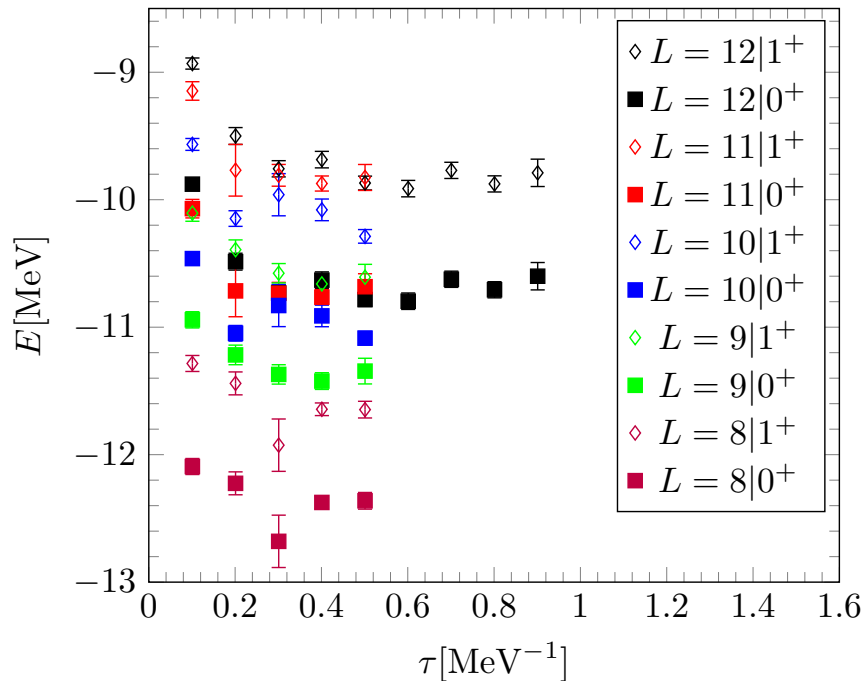
Experiment

$$B_{\Lambda}({}^3_{\Lambda}\text{H}) = 0.08 \pm 0.05 \text{ MeV}$$

$$0.164 \pm 0.43 \text{ MeV}$$

→ Result consistent with newest experimental data

- Hypertriton has large Λ separation of ~ 10.6 fm
- Second largest system is the 4-body system



→ L = 12 shows converged results

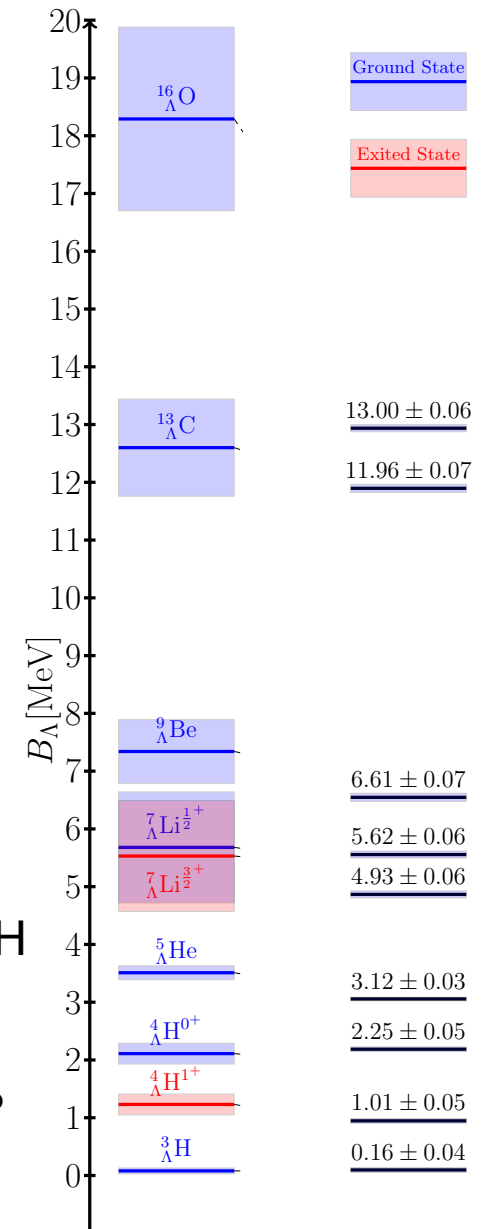
→ L = 12 Box is big enough for all other hypernuclei

→ Finite Box under control

Results: Two Body interaction, further analysis

Missing ~ 0.2 MeV in $A=4$ systems

$A=5$ system only slightly overbound



Less Attraction



Modify C



Fix $^5_\Lambda\text{He}$



Underbinding of $^4_\Lambda\text{H}$



Can three-body forces help us here?

Experiment

Structure of contact three-body forces

(Petschauer et al.)

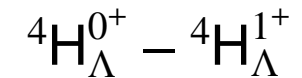
$$V_{ct}^{\Lambda NN} = C_1(1 - \boldsymbol{\sigma}_2 \cdot \boldsymbol{\sigma}_3)(3 + \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3) \quad \leftarrow I = 1$$

$$+ C_2 \boldsymbol{\sigma}_1 \cdot (\boldsymbol{\sigma}_2 + \boldsymbol{\sigma}_3)(1 - \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3) \quad \leftarrow I = 0$$

$$+ C_3(3 + \boldsymbol{\sigma}_2 \cdot \boldsymbol{\sigma}_3)(1 - \boldsymbol{\tau}_2 \cdot \boldsymbol{\tau}_3) \quad \leftarrow I = 0$$

Can influence hypertriton

C_2 only interaction that depends on Λ spin



Splitting

C_1, C_3 are iso/spin interchanged to each other

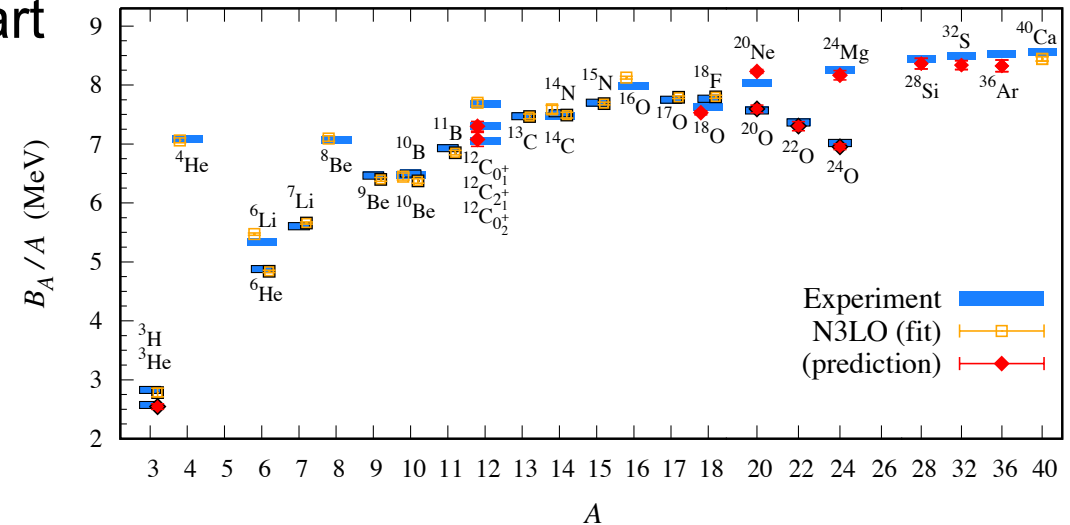
Might not split
for small
Hypernuclei

Effectively N3LO χ EFT(NN) + LO π EFT(YN)

Results: Fitting 3-Body forces

Nuclear 3-Body Forces are fitted as part of the WFM interaction

→ Use similar classes of non-local as well as local smeared YNN Forces



(Elhatisari et al.)

→ Leads to a total of 343 (7 each) combination of YNN forces

→ Use Root Mean Square Deviation to select TBF :

$$\text{RMSD}(S) = \sqrt{\frac{1}{M_S} \sum_{i \in S} \left(\frac{iB_{\Lambda}^c - iB_{\Lambda}^{\text{exp}}}{iB_{\Lambda}^{\text{exp}}} \right)^2}$$

→ Without TBF RMSD(S) = 18.4 %

First approach Decouplet Saturation

$$\begin{aligned}
 V_{ct}^{\Lambda NN} &= C_1(1 - \sigma_2 \cdot \sigma_3)(3 + \tau_2 \cdot \tau_3) \longrightarrow C_1 \\
 &+ C_2 \sigma_1 \cdot (\sigma_2 + \sigma_3)(1 - \tau_2 \cdot \tau_3) \longrightarrow C_2 = 0 \\
 &+ C_3(3 + \sigma_2 \cdot \sigma_3)(1 - \tau_2 \cdot \tau_3) \longrightarrow C_1
 \end{aligned}$$

Enforce same smearing

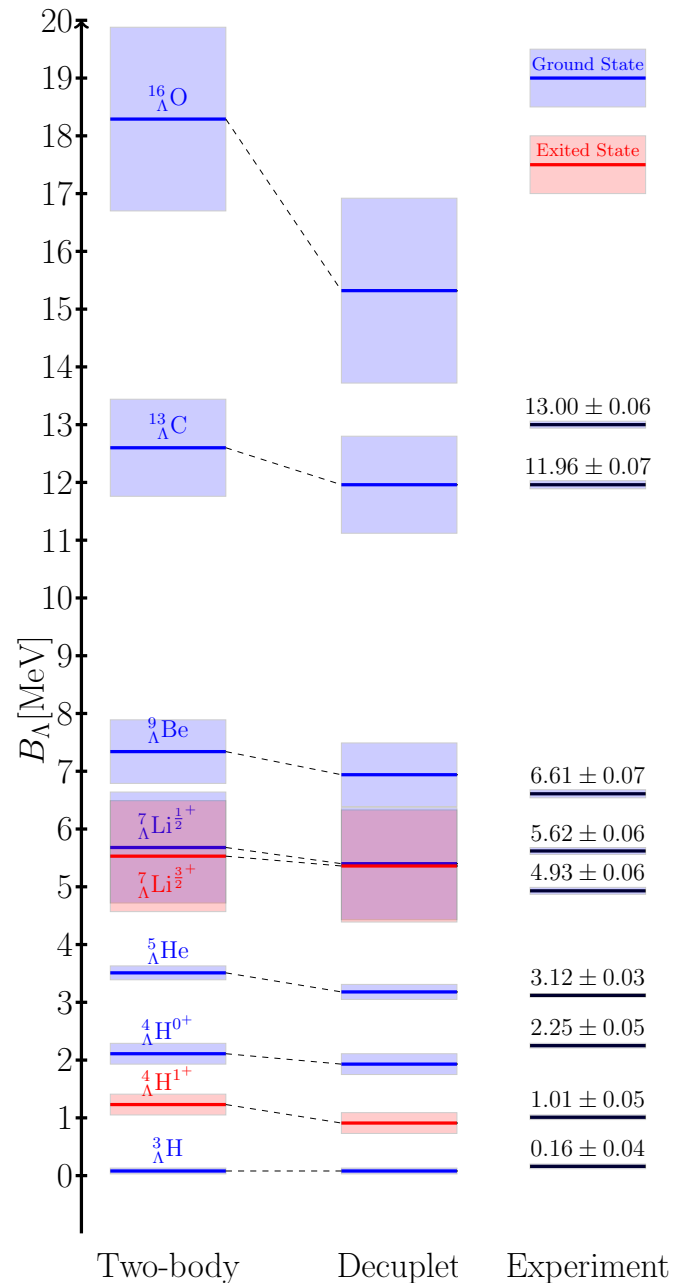
Also include excited state of ${}^7_{\Lambda}\text{Li}$ and ${}^9_{\Lambda}\text{Be}, {}^{13}_{\Lambda}\text{C}, {}^{16}_{\Lambda}\text{O}$, so far all overbound

	Fit to all $A \geq 4$	Fit to only $A=4/5$
RMSD(S)	9.2% – 14.6%	9.3% – 14.7%

Since $C_2 = 0$:
splitting remains untouched

Improvement due to an overall downwards shift

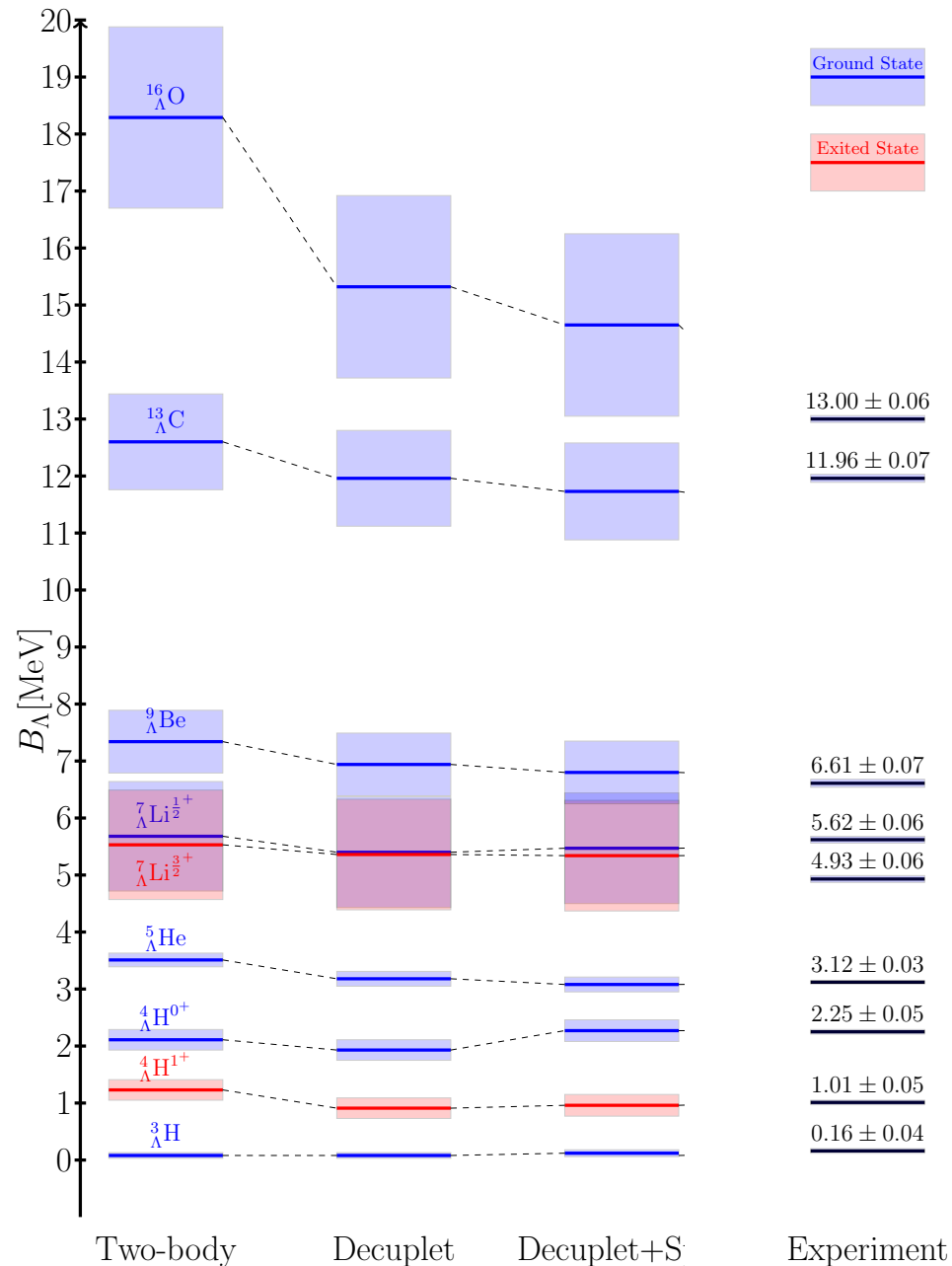
How does this translate to Binding energies?



- Overall downwards shift in energies
- Underbinding of the 4-body systems
- Splitting in 4/7-body system is still the same
- Naiv improvment:

$$\begin{aligned}
 V_{ct}^{\Lambda NN} = & C_1(1 - \sigma_2 \cdot \sigma_3)(3 + \tau_2 \cdot \tau_3) \longrightarrow C_1 \\
 & + C_2 \sigma_1 \cdot (\sigma_2 + \sigma_3)(1 - \tau_2 \cdot \tau_3) \longrightarrow C_2 \neq 0 \\
 & + C_3(3 + \sigma_2 \cdot \sigma_3)(1 - \tau_2 \cdot \tau_3) \longrightarrow C_1
 \end{aligned}$$

Decouplet + Spin dependent force?



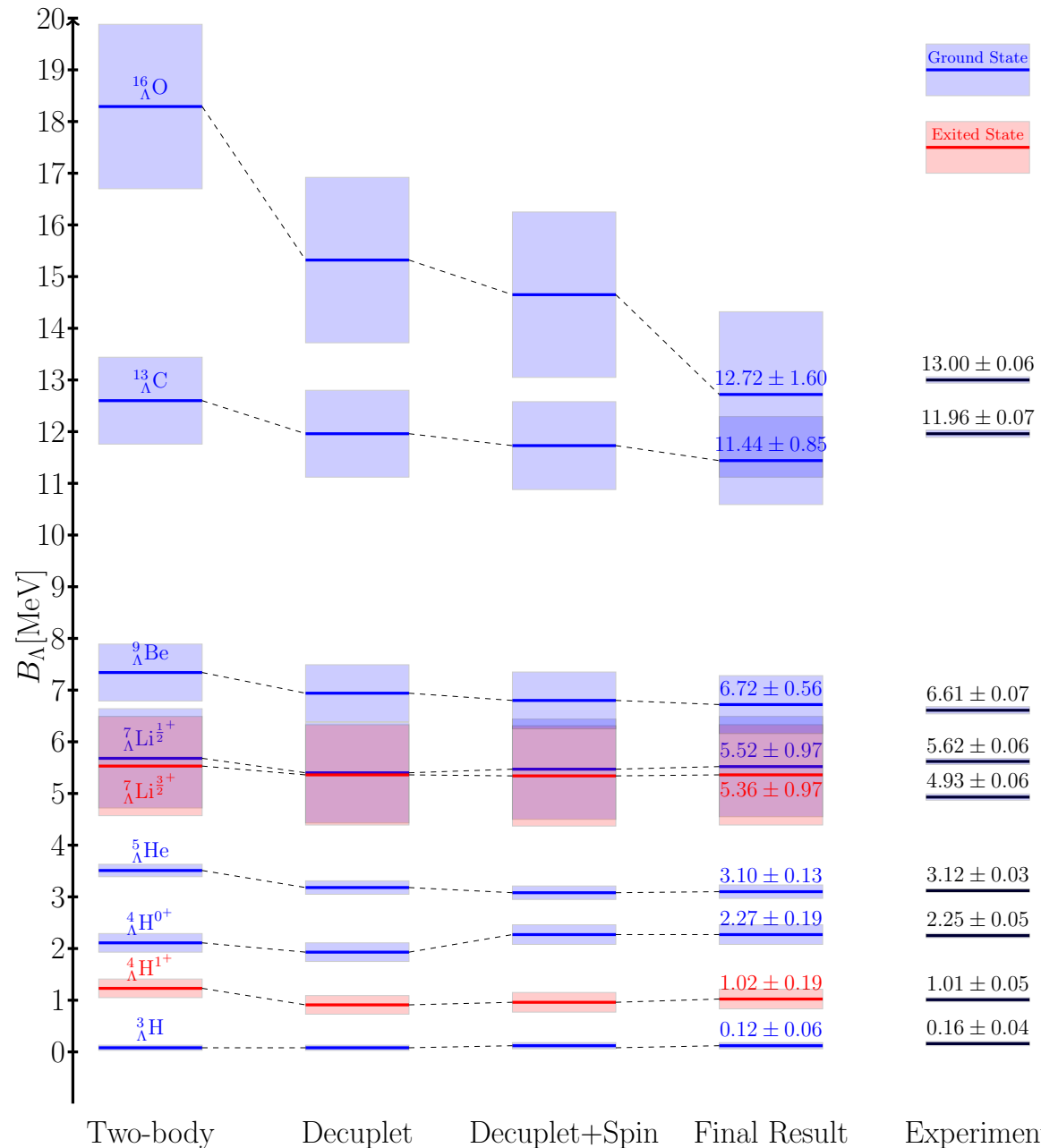
- 49 different combinations of forces
- If fitted to the 4/5 body system 21 improve the overall description

	Fit to all $A \geq 4$	Fit to only $A=4/5$
Best:	5.3%	5.7%

Weakly smeared forces outperform
Stronger smeared forces

Original selection of smearing parameter
is sufficient

Final result



- All 343 different combinations

Fit to all $A \geq 4$ Fit to only $A=4/5$

3.6 %

3.7 %

- All hypernuclei are consistent

- Good splitting in the 4 body system

- Best results have only local smearing, V_2 is always unsmearred

Possible Paths to improvement

- Go to higher orders in the two-body interaction



Typical LO problems
go away in other
methods

- Include two-pion exchange/pion exchange 3B forces



Long-Range behaviour
of the interaction

- fit two-body forces with better nuclear interaction



Removes any
dependence of the NN
Force on the YN Force

- Improve statistics in the NN part of the hypernuclei



Main uncertainty from
sampling of the NN
part of the nucleus

Summary and Outlook

Good Results for light hypernuclei nuclei $A=3-16$
with $N^3LO(NN)$ and $LO(YN)$ interaction

Method scales with A , straightforward application to the whole
hypernuclear chart

Many possible path ways to improve the results

Calculate the hypernuclear chart

Many excited states in $A=7/9$ hypernuclei

