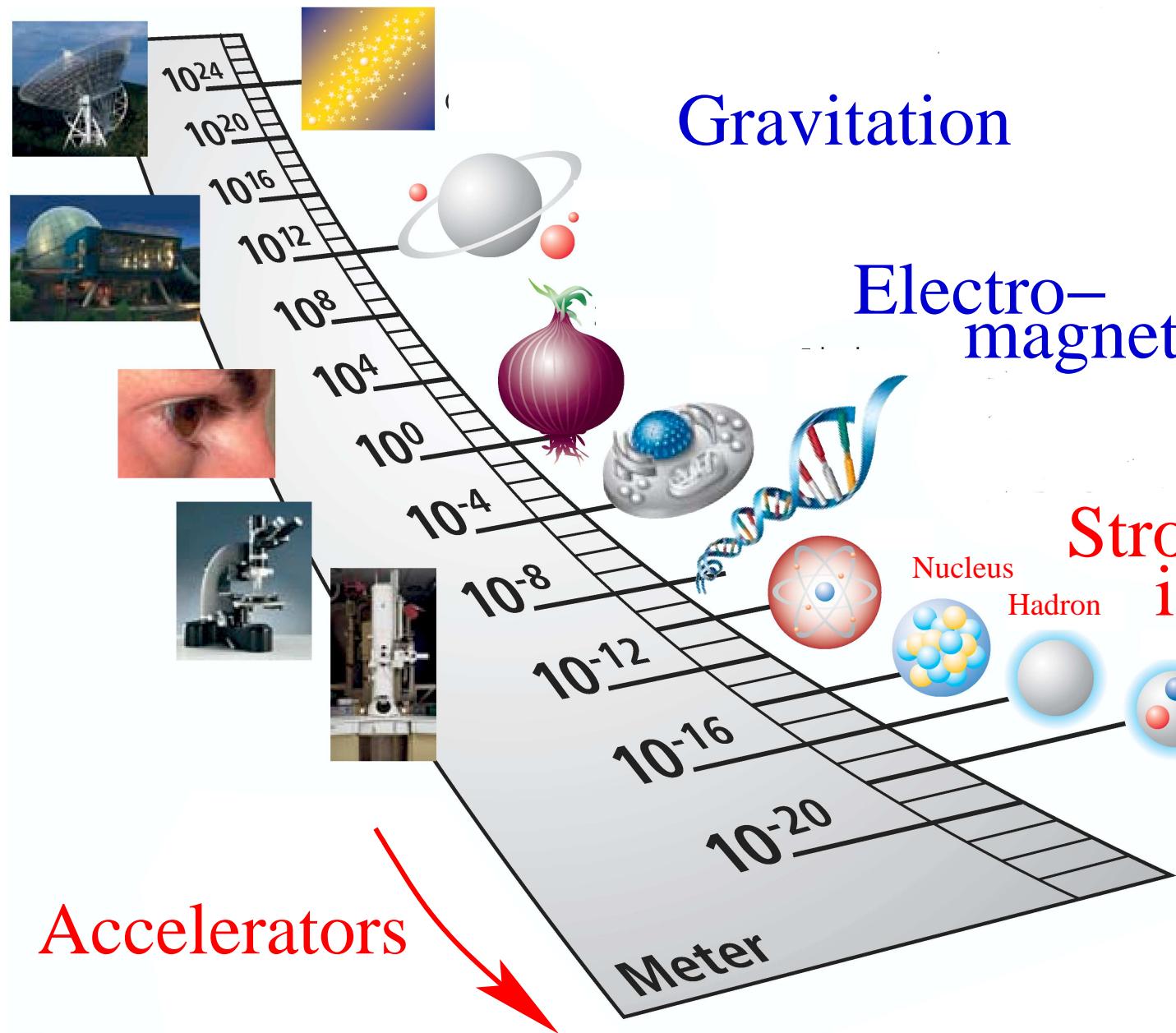


# Highlights and Challenges of Hadron Physics

Christoph Hanhart

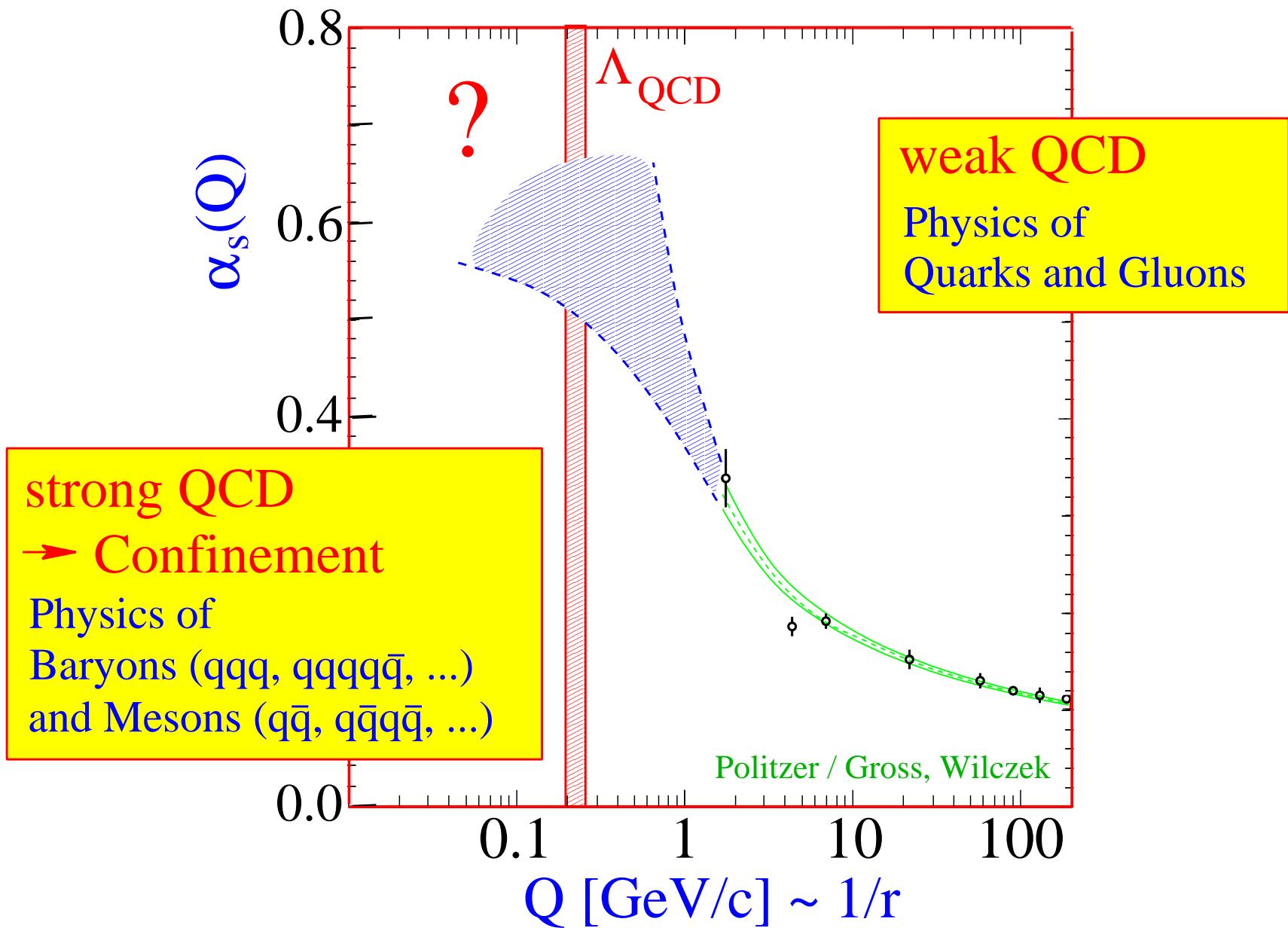
Forschungszentrum Jülich

# The fundamental interactions

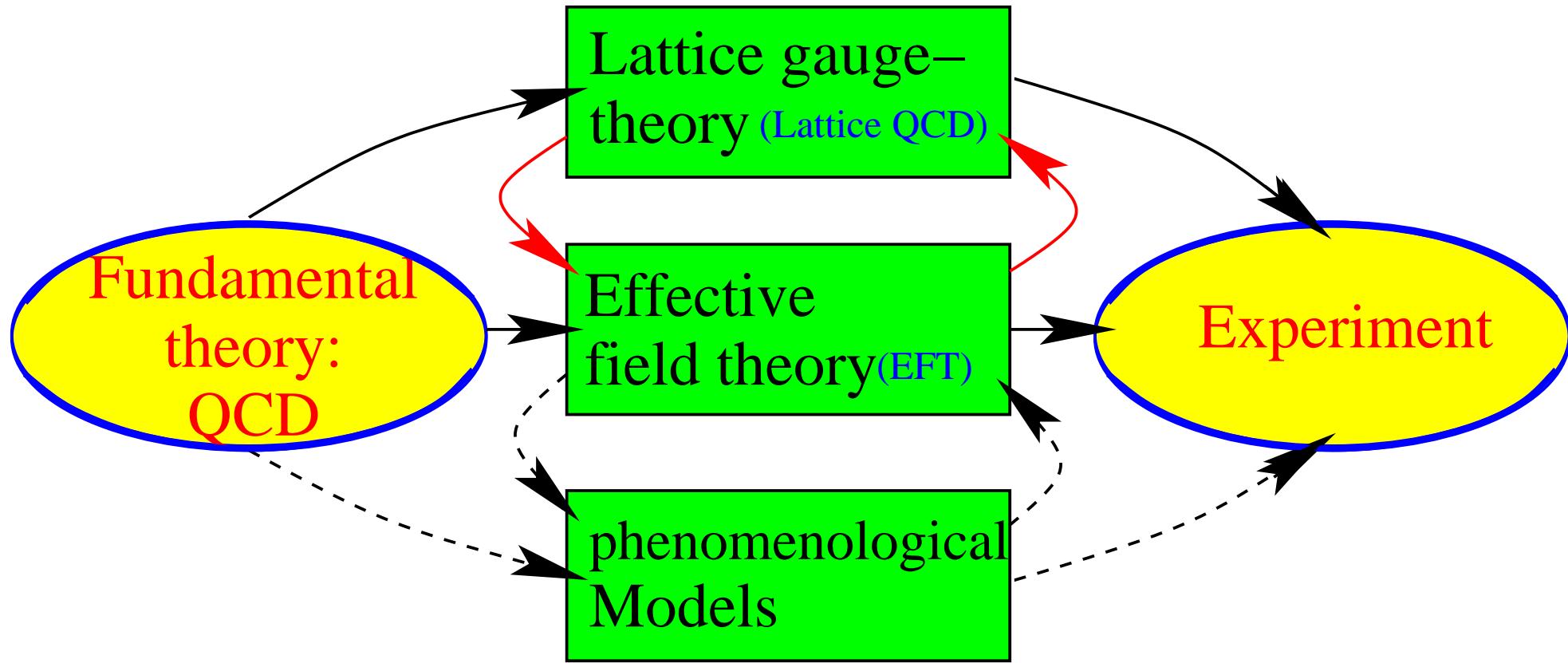


in addition:  
weak  
interaction

# The faces of QCD



# Strategies to treat strong QCD



- **Lattice Gauge Theory**: direct solution  
but requires long computing times → few Observables
- **EFT**: most general structure compatible with QCD  
but with increasing accuracy more free parameters

# Limiting cases of QCD

Weinberg/ Gasser, Leutwyler/ Isgur, Wise

$$\mathcal{L}_{\text{QCD}} = \sum_f \bar{q}_f \{ i\partial - m_f + g A^a t^a \} q_f - \frac{1}{4} G_{\lambda\rho}^a G^{\lambda\rho a} + \dots$$

## Limit of massless Quarks

$$\mathcal{L}_{\text{QCD}} = \bar{q}_L \{ i\partial + g A^a t^a \} q_L + \bar{q}_R \{ i\partial + g A^a t^a \} q_R + \mathcal{O}(m_f/\Lambda_{\text{QCD}})$$

L and R Quarks decouple + spontaneous symmetry breaking

→ Chiral Perturbation Theory (ChPT)

## Limit of infinitely Heavy Quarks

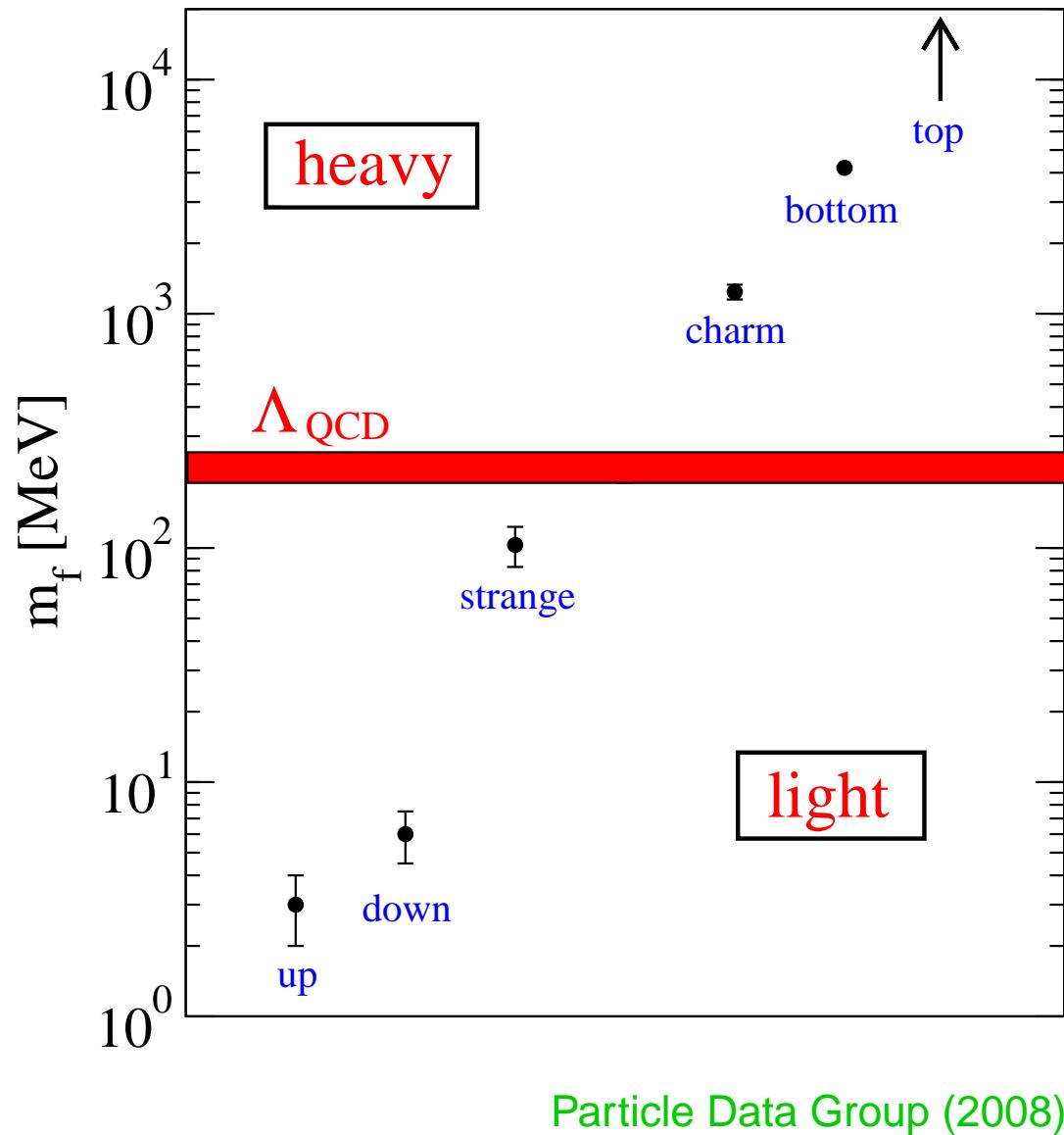
$$\mathcal{L}_{\text{QCD}} = \bar{q}_f \{ i v \cdot \partial + g v \cdot A^a t^a \} q_f + \mathcal{O}(\Lambda_{\text{QCD}}/m_f)$$

Independent of Heavy Quark Spin and Flavour

→ Heavy Quark Effective Field Theory (HQEFT)

# Quark Masses

Quark Masses (in  $\overline{\text{MS}}$  at  $\mu=2 \text{ GeV}$ )

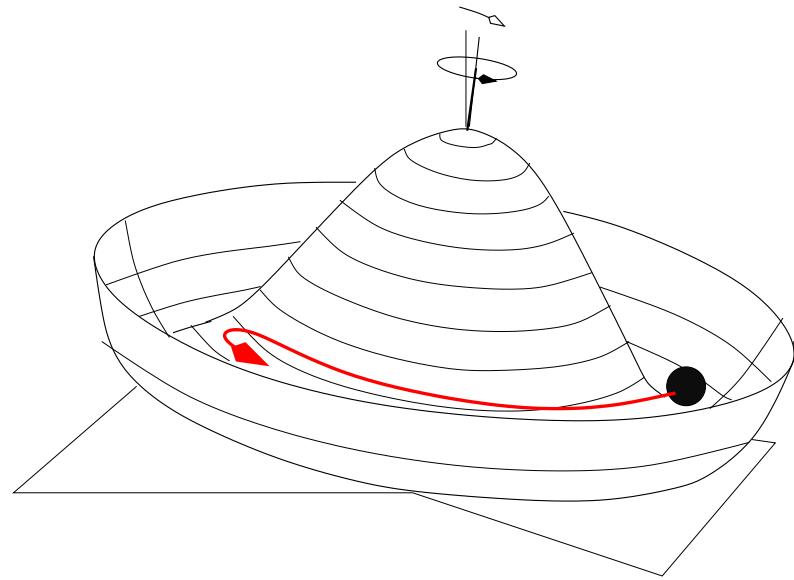
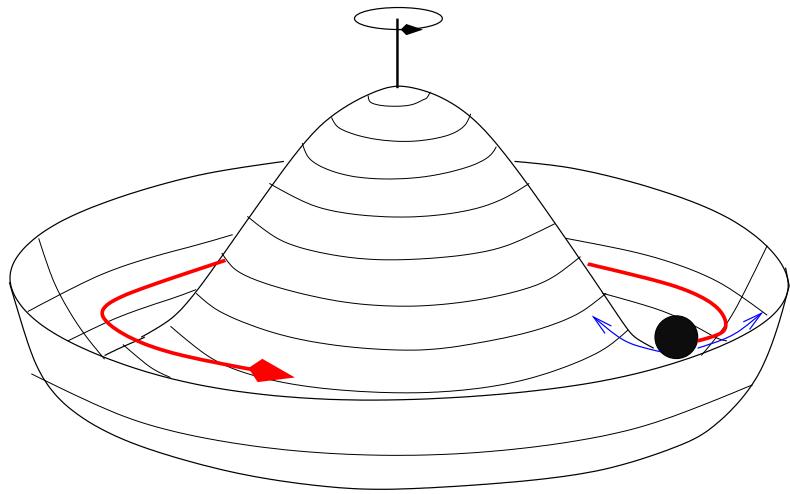
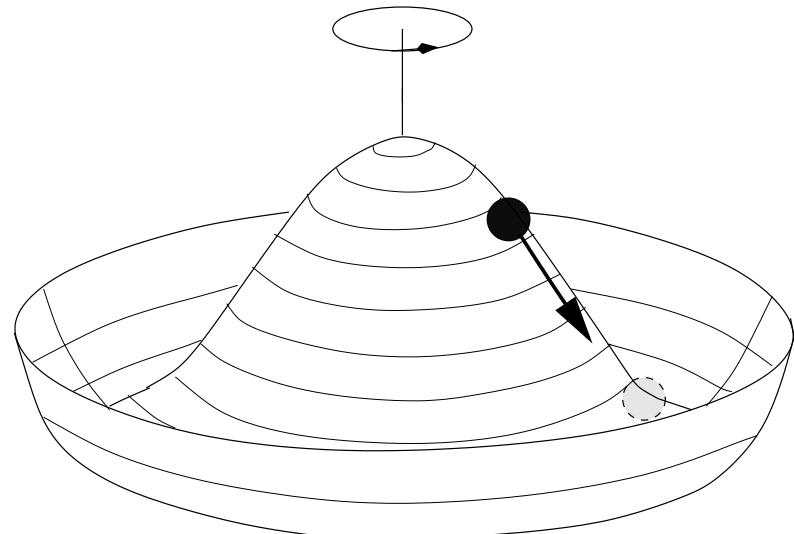
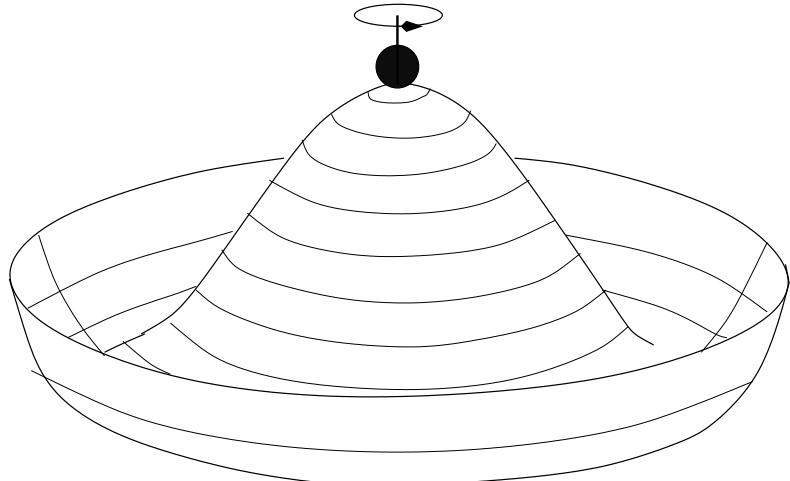


Expect **very different phenomena** for light (u,d,s) and **heavy (c,b)** quarks

- What are the spectra?
- What structures are there?

Study **systematically** particle properties, decays, and interactions!

# Spontaneous Symmetry Breaking



$m_q=0$ : massless Goldstone bosons;  $m_q \neq 0$ : light physical pion

# Example: CSB

Charge Symmetry  $\equiv$  physics invariant under  $\text{up} \leftrightarrow \text{down}$

In Standard Model Charge Symmetry Breaking (CSB) via

→ different quark masses:

$$-m_u \bar{u}u - m_d \bar{d}d \implies \frac{\delta m_N^{\text{str}}}{2} N^\dagger \left( \tau_3 - \frac{1}{2F_\pi^2} \boldsymbol{\tau} \cdot \boldsymbol{\pi} \pi_3 \right) N$$

→ different quark charges:

$$+q_u \bar{u} \not{A} u + q_d \bar{d} \not{A} d \implies \frac{\delta m_N^{\text{em}}}{2} N^\dagger \left( \tau_3 - \frac{1}{2F_\pi^2} (\tau_3 \boldsymbol{\pi}^2 - \boldsymbol{\tau} \cdot \boldsymbol{\pi} \pi_3) \right) N$$

with  $\delta m_N^{\text{str}} + \delta m_N^{\text{em}} = M_p - M_n = 1.3 \text{ MeV}$

There are two recent measurements:

- Forward–backward asymmetry in  $pn \rightarrow d\pi^0$  ( $A_{fb}$ )  
to see this assume recall  $T_d = 0$  and  $T_\pi = 1$   
thus: initial state is  $T = 1 \Rightarrow$  should behave as  $pp$   
**forward and backward not defined**  
F–b asymmetry due to  $T = 0$  admixture in initial state
- $\sigma_{tot}(dd \rightarrow \alpha\pi^0)$   
to see this recall  $T_d = T_\alpha = 0$

Experimentally:

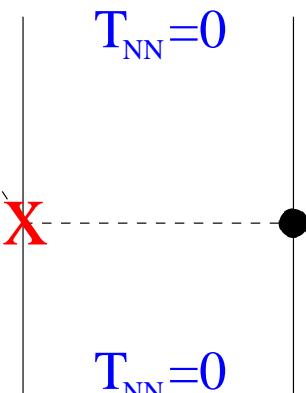
$$A_{fb}(pn \rightarrow d\pi^0) = 0.17 \pm 0.08(\text{st.}) \pm 0.06(\text{sys.}) \%$$
 (TRIUMF, 2003)

$$\sigma_{tot}(dd \rightarrow \alpha\pi^0) = \begin{cases} 13 \pm 2 \text{ pb} & \text{at } T_{lab} = 228.5 \text{ MeV} \\ 15 \pm 3 \text{ pb} & \text{at } T_{lab} = 231.8 \text{ MeV} \end{cases}$$
 (IUCF, 2003)

# Theoretical aspects

Relevant transition:

A. Filin et al. (2010); for  $dd$  reaction: CSB collaboration



$$A_{\text{fb}}^{\text{LO}} = (11.5 \pm 3.5) \times 10^{-4} \frac{\delta m_N^{\text{str}}}{\text{MeV}}$$

$$\implies \delta m_N^{\text{str}} = (1.5 \pm 0.8 \text{ (exp.)} \pm 0.5 \text{ (th.)}) \text{ MeV},$$

from lattice

Beane et al. (2008)

$$\implies \delta m_N^{\text{str}} = 2.26 \pm 0.57 \pm 0.42 \pm 0.10 \text{ MeV}$$

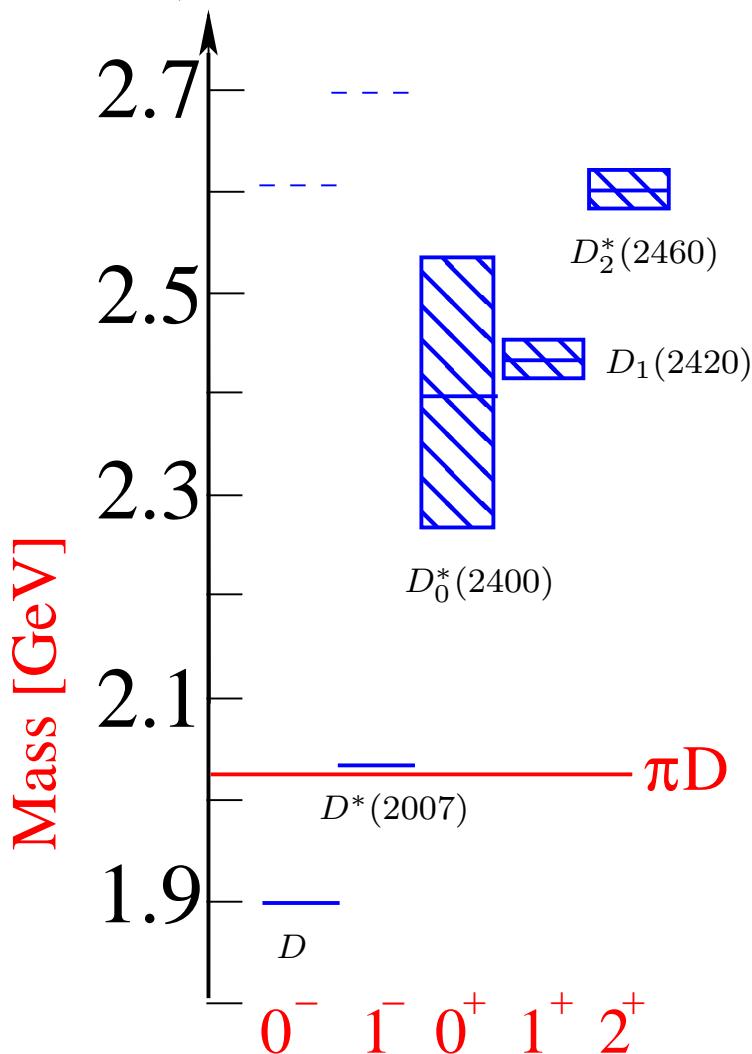
consistent with  $\delta m_N^{\text{em}} = -0.7 \pm 0.3 \text{ MeV}!!$

Gasser, Leutwyler (1982)

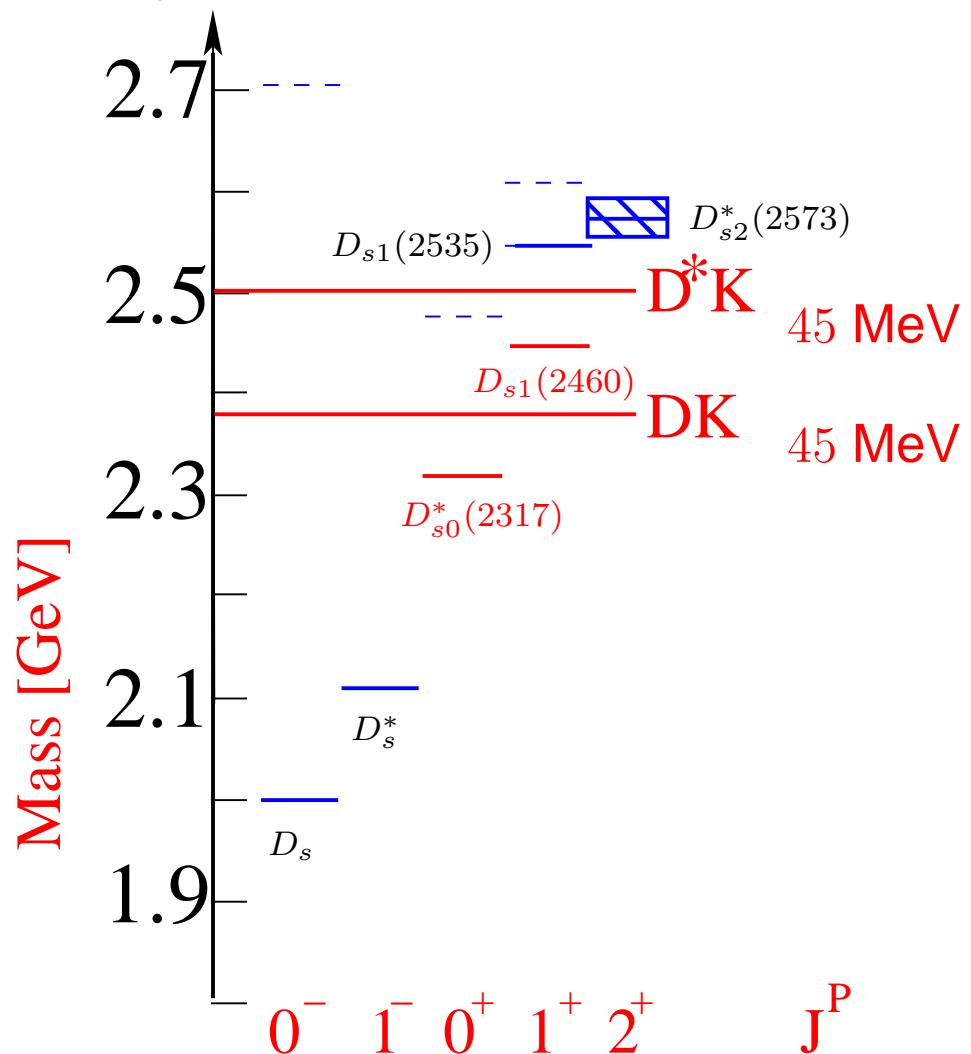
- gives **leading order effect** to CSB pion production
- (potentially) sizable loop/higher order corrections; **to be done**
- $dd \rightarrow \alpha\pi^0$  has **complicated few–body physics**
  - ▷ data for  $dd \rightarrow {}^3\text{He}N\pi \implies$  constrains  $dd$  dynamics
  - ▷ higher energy data for  $dd \rightarrow \alpha\pi^0 \implies$  **important test**

# Spectrum with open Charm

S=0, I=1/2



S=1, I=0



Quark Modell: Di Pierro, Eichten (2001)

# Heavy light Systems

- Effective field theorie ChPT and HQEFT control interactions
  - ▷ Systematic improvements + uncertainty estimates
  - ▷ consistent inclusion of isospin breaking
  - ▷ quark mass dependence
- unitarization  
generation of resonances and bound states

Kaiser, Weise/ Oller, Oset, Pelaez/ Lutz, Kolomeitsev, Soyeur/  
Guo, C.H., Krewald, Meißner (2008)/ Guo, C.H., Meißner (2009)

Here:  $\pi/K/\eta$ - $D/D_s$  Streuung

→  $D_{s0}^*(2317)$  as bound state mass fitted

- $D_{s1}(2460)$  + other other things predicted;
- uncertainty estimates through higher order Operators

Isospin breaking in QCD and EFT through **quark mass and charge differences**

The **same effective operators** lead to

→ **mass differences**, e.g.

$$\begin{aligned}\triangleright m_{D^+} - m_{D^0} &= \Delta m^{\text{strong}} + \Delta m^{\text{e.m.}} \\ &= ((2.5 \pm 0.2) + (2.3 \pm 0.6)) \text{ MeV}\end{aligned}$$

→  $\pi^0 - \eta$  mixing → **parameters fixed**

→ **Isospin breaking scattering amplitude**

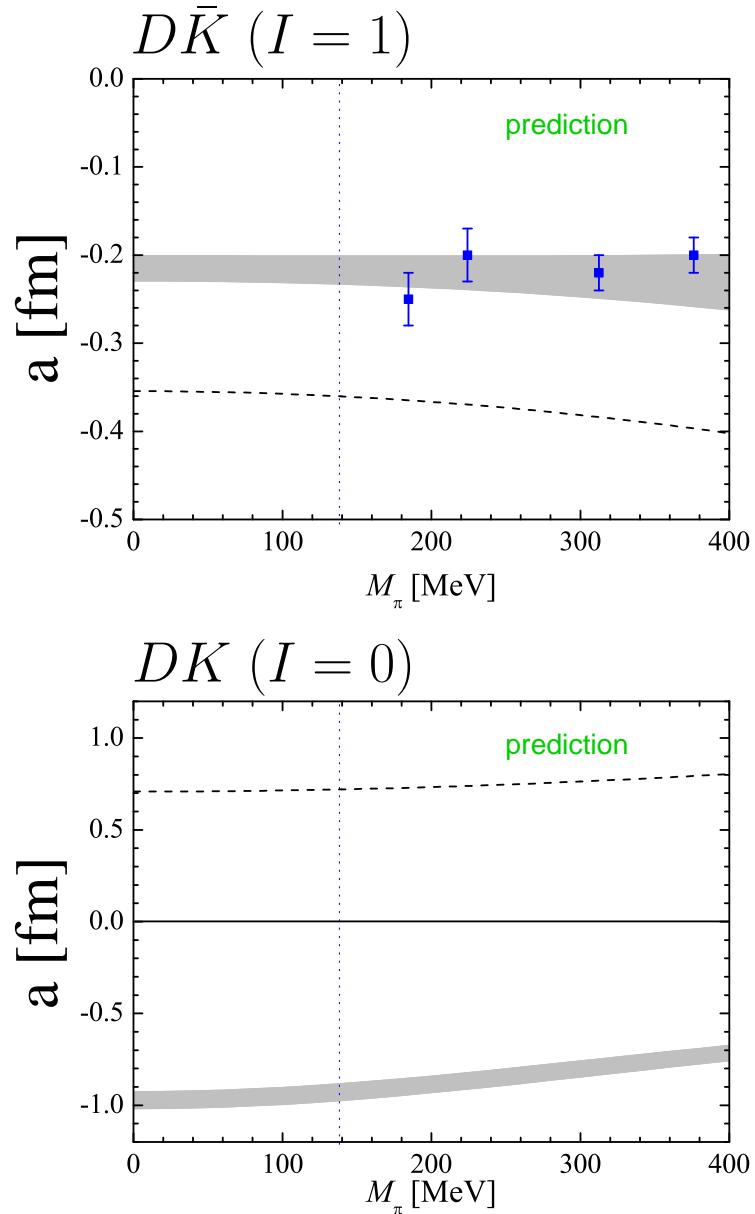
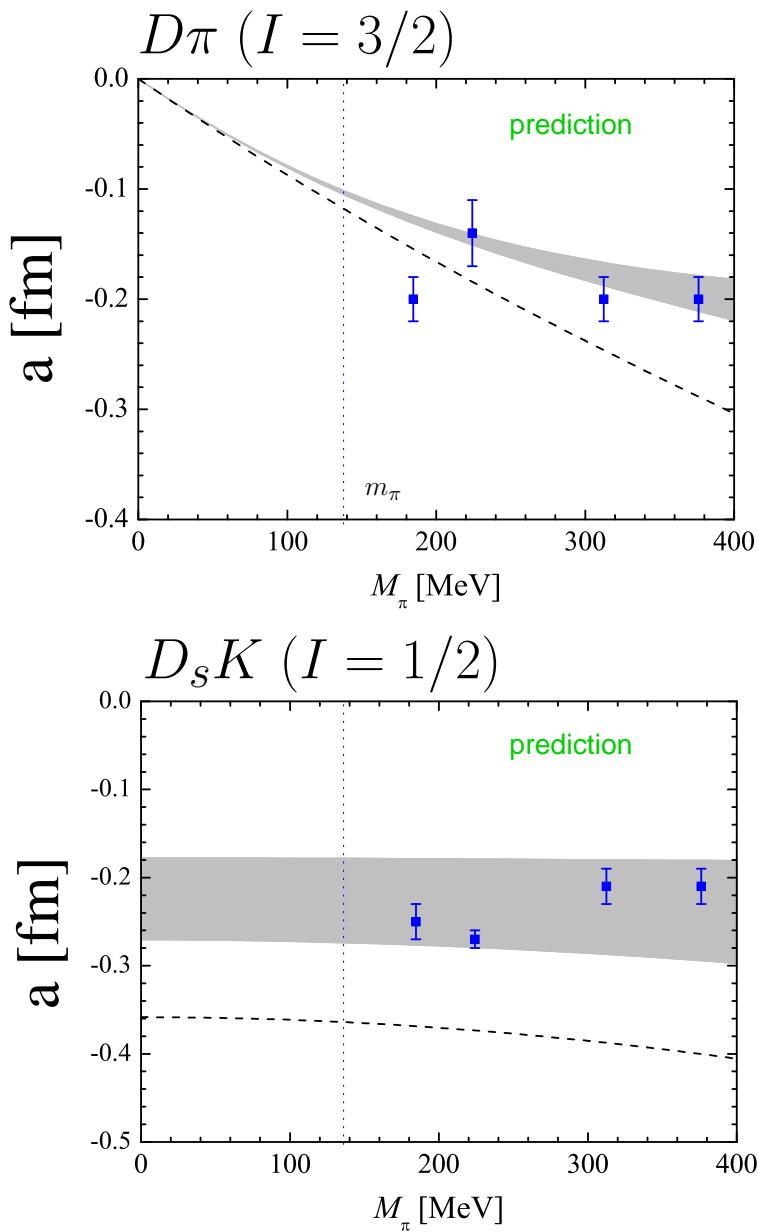
→ e.g.  $K D \rightarrow \pi^0 D_s$  **predicted**; with this

$$\begin{aligned}\Gamma(D_{s0}^*(2317) \rightarrow D_s \pi^0) &= (133 \pm 110) \text{ keV} \\ \Gamma(D_{s1}(2460) \rightarrow D_s \pi^0) &= (126 \pm 110) \text{ keV}\end{aligned}$$

Lutz, Soyeur (2007); complete to NLO+uncertainty estimate: Guo et al. (2008)

much smaller in quark model → **direct measurement necessary**

# Chiral extrapolation

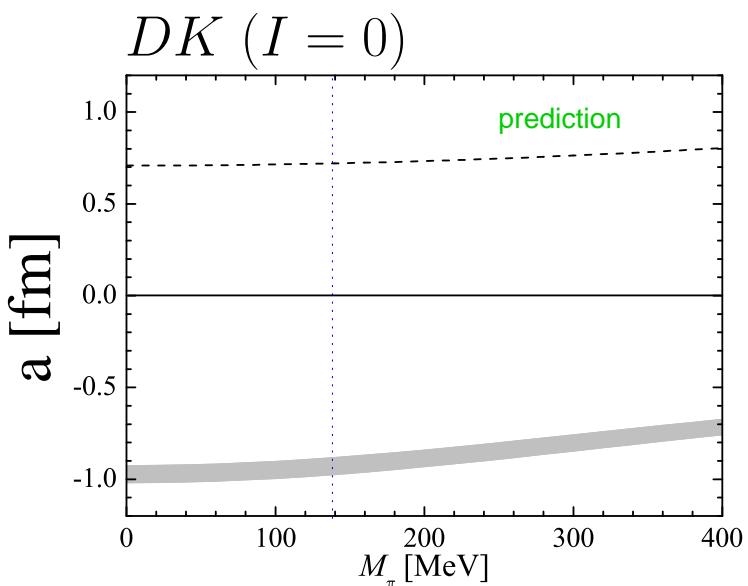
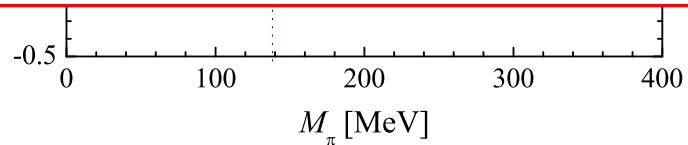
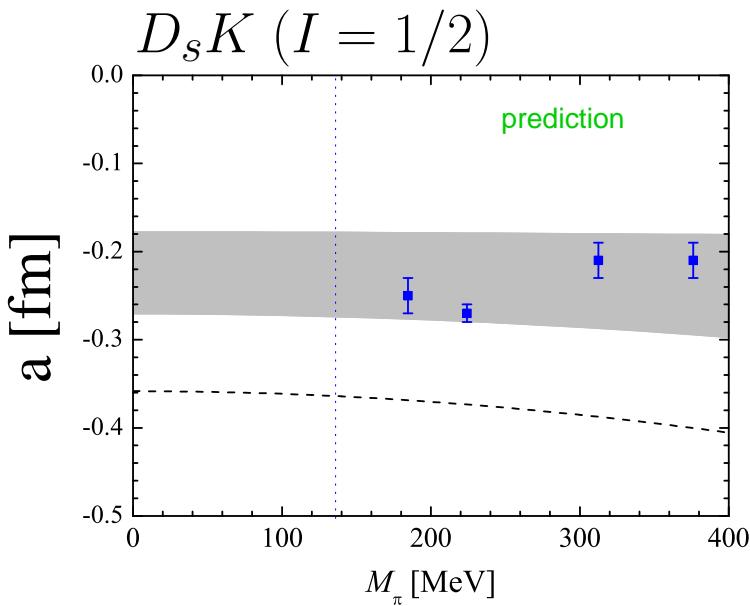
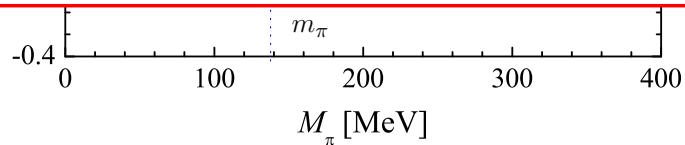


Lattice: Liu, Lin, Orginos (2008); UChPT: Guo et al. (2009)

# Chiral extrapolation

$$D_S(2317): a = g_{\text{eff}} + \mathcal{O}(1/\beta) \simeq \left( \frac{2\lambda^2}{1+\lambda^2} \right) \frac{-1}{\sqrt{2m_K\epsilon}}$$

$a = 1 \text{ fm}$  for molecule ( $\lambda^2 = 1$ ); smaller otherwise



Lattice: Liu, Lin, Orginos (2008); UChPT: Guo et al. (2009)

# Thus ...

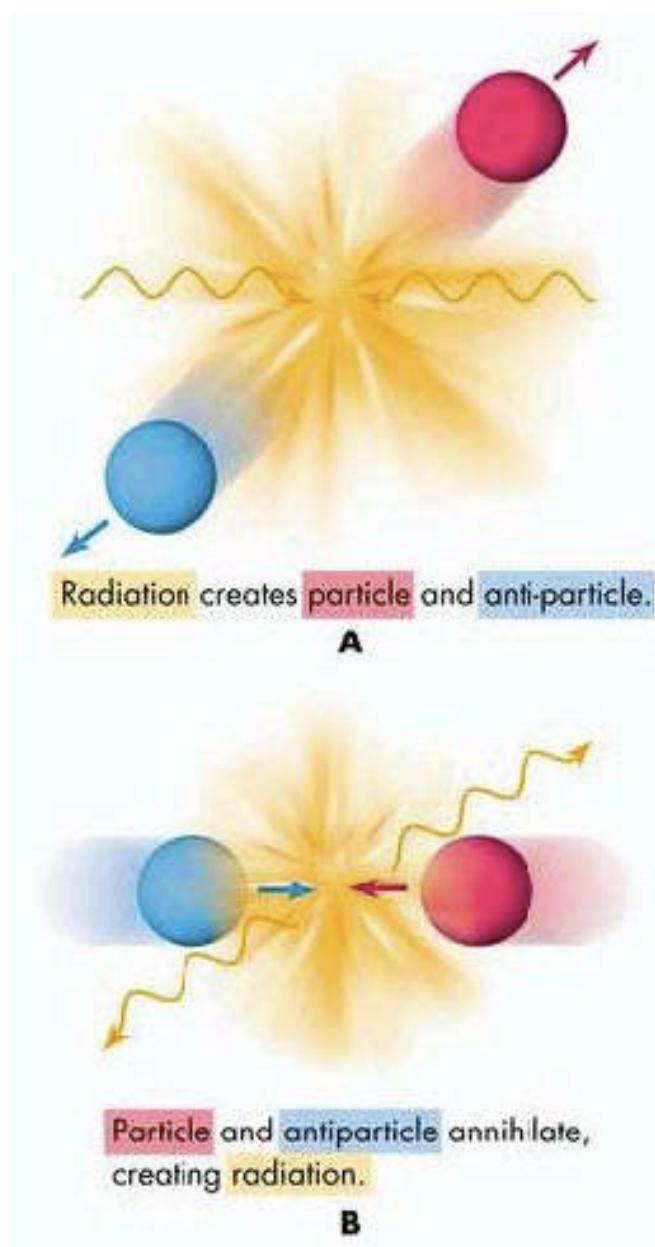
- Effective field theories hadron physics can be studied
  - ▷ to high precision
  - ▷ with controlled uncertainties
- Combination of EFT and resummation schemes very promising
- Symbiosis of lattice gauge theory and EFT on the horizon
  - we are on a good way to understand QCD

Standard Model tested in many aspects

→ it (nearly) always works

But: WE KNOW IT IS NOT COMPLETE!

# Matter Excess



In the early universe:

matter:antimatter as

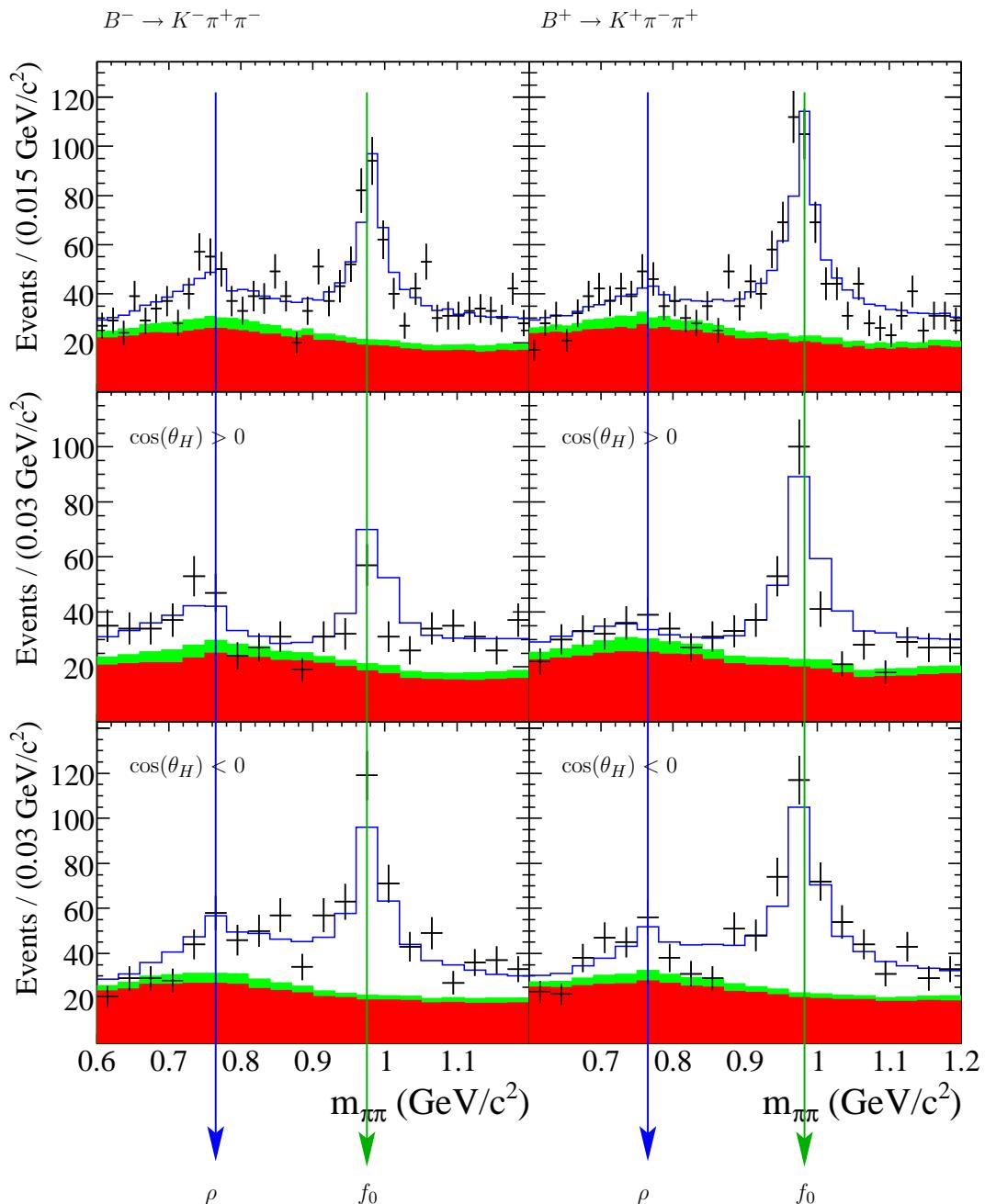
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**Problem:** The Standard Model falls short by several orders of magnitude!

→ What is missing?

Look for the needle in the haystack

# Sensitive Observables



Here:  $B^\pm \rightarrow K^\pm \pi^\mp \pi^\pm$

BABAR (2008); Belle (2006)

- $3.7\sigma$  in  $\rho K$
- $3.5\sigma$  in  $f_2(1270)K$
- $1.8\sigma$  in  $f_0(980)K$

Consistent with SM

Important next steps:

- Improve analysis and thus sensitivity  
⇒ Les NABIS
- Study  $D$  decays  
→ SM prediction tiny

# Summary

Various experiments (will be) looking for systems with charm

- BABAR (USA – data taking finished)
- Belle (Japan) (in operation) + Belle II (planned)
- BES-III (China) (in operation)
- LHCb (at LHC, CERN; just started)
- PANDA (FAIR, Germany; after 2014)

together with the expected developments of theoretical tools

There are exciting times to come!

THANKS FOR YOUR ATTENTION