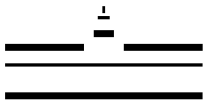


Investigation on the tensor analyzing power t_{20} in the reaction $\vec{d} + p \rightarrow {}^3\text{He} + \eta$

Michael Papenbrock

Westfälische Wilhelms-Universität Münster
for the ANKE collaboration

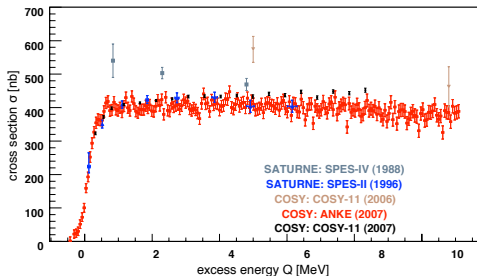


WESTFÄLISCHE
WILHELMS-UNIVERSITÄT
MÜNSTER

March 5th, 2013
DPG spring conference, Dresden

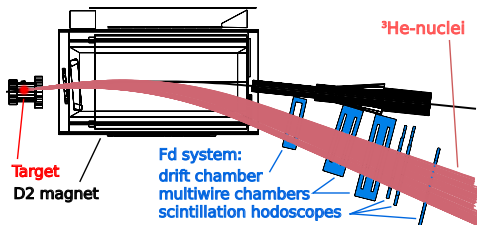
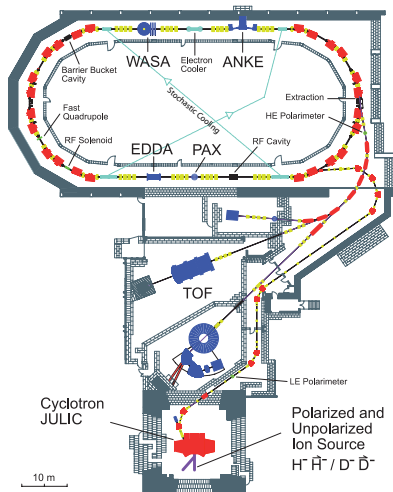
Motivation

- ▶ Previous unpolarized measurements on $\mathbf{dp} \rightarrow {}^3\mathbf{He}\eta$ provided detailed cross section data near the production threshold
- ▶ Excitation function has shown indications of an unexpected strong final state interaction (FSI)



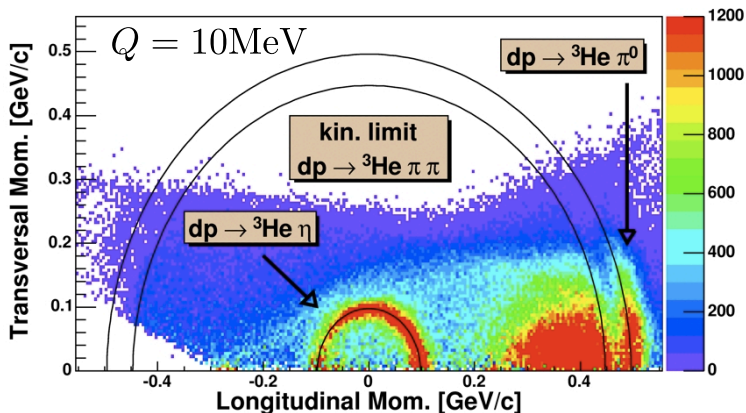
- ▶ Test for contributions to the cross section other than FSI possible by polarized measurements
- ▶ ANKE: Investigation of the reaction with a tensor/vector polarized deuteron beam

Measuring $dp \rightarrow {}^3\text{He} \eta$ at ANKE at COSY



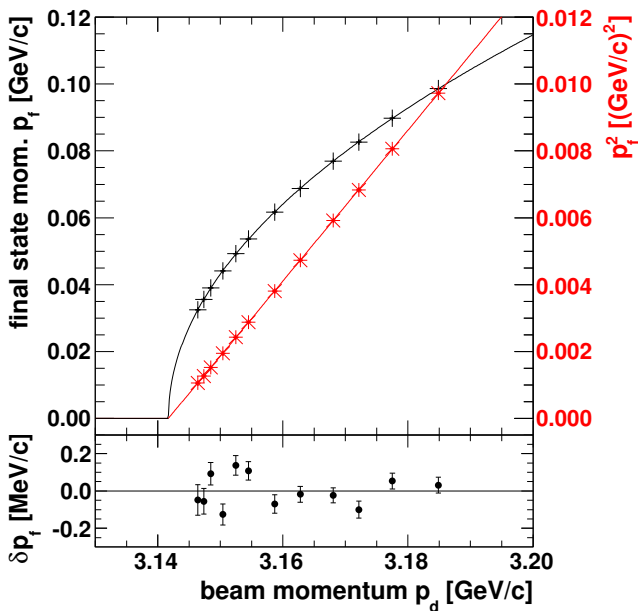
- ▶ Internal fixed target experiment with a cluster-jet target
- ▶ ${}^3\text{He}$ nuclei detected in the forward-system
- ▶ Full geometrical acceptance for $dp \rightarrow {}^3\text{He} \eta$ up to 20 MeV excess energy

Acceptance of the ANKE detector system



- ▶ Full acceptance for $dp \rightarrow {}^3\text{He} \eta$
- ▶ Continuously ramped beam from $Q = -5 \text{ MeV}$ to $Q = 11 \text{ MeV}$

High precision η mass determination

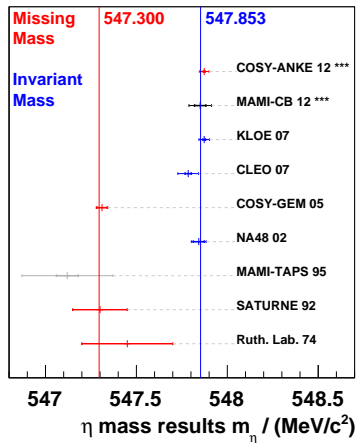


ANKE-COSY result of the η mass

$$m_{\eta} = (547.873 \pm 0.005_{\text{stat.}} \pm 0.023_{\text{sys.}}) \text{ MeV}/c^2$$

ANKE η meson mass

- ▶ Highest precision measurement yet
- ▶ In agreement with higher η meson mass measurements
- ▶ Published in Physical Review D85, 112011 (2012)



$^3\text{He}n\eta$ final state interaction

Description of the differential cross section

- ▶ Differential cross section of a two-body reaction:

$$\frac{p_i}{p_f} \cdot \frac{d\sigma}{d\Omega} = |f|^2 = |f_B \cdot FSI|^2 = |f_B|^2 \cdot |FSI|^2$$

- ▶ Effective range approximation:

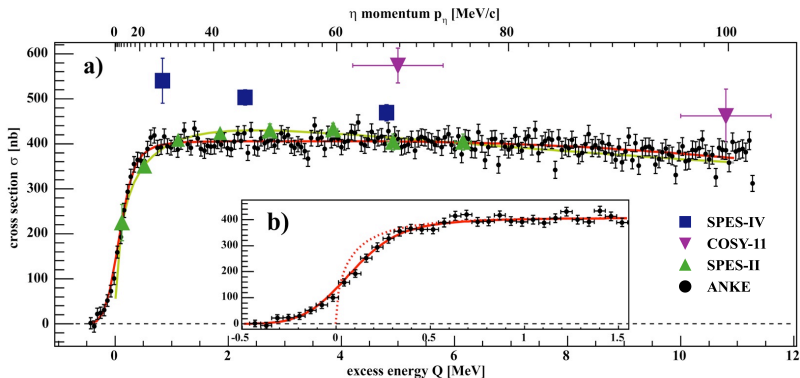
$$FSI = \frac{1}{1 + i \cdot a \cdot p_f + \frac{1}{2} a \cdot r_0 \cdot p_f^2}$$

- ▶ Alternative description with poles:

$$FSI = \frac{1}{\left(1 - \frac{p_f}{p_1}\right) \cdot \left(1 - \frac{p_f}{p_2}\right)}$$

with $a = -i \cdot \frac{p_1 + p_2}{p_1 \cdot p_2}$ and $r_0 = +\frac{2 \cdot i}{p_1 + p_2}$

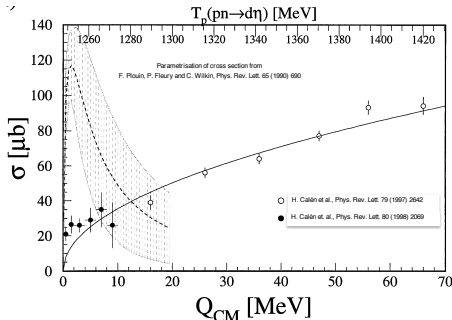
${}^3\text{He}\eta$ final state interaction: unpolarized data



- ▶ Very good description for the whole energy range
- ▶ Pole close to threshold ($|Q_0| \approx 0.4$ MeV) might be an indication for a quasi-bound state

Investigation of the reaction $pn \rightarrow d\eta$

- ▶ Investigation of A dependence of the η -nucleus FSI
→ further important information about η -mesic nuclei
- ▶ Determination of $d\eta$ scattering length
→ information about still unprecisely known ηN scattering length



- ▶ Measurement has been performed earlier this year

Why polarised measurements?

Additional information for FSI ansatz needed

- ▶ Test for further contributions besides FSI
- ▶ Investigate production amplitude close to threshold

Alternative description of the $\eta^3\text{He}$ s-wave production amplitude:

$$f = \bar{u}_{3\text{He}} \hat{p}_p \cdot (A \vec{\epsilon}_d + iB \vec{\epsilon}_d \times \vec{\sigma}) u_p$$

(Shown in J.-F. Germond & C. Wilkin, J. Phys. G 14, 181 (1988))

$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_f}{p_i} \left[|A|^2 + 2 \cdot |B|^2 \right]$$

with A, B as the two s-wave amplitudes of the $^3\text{He}\eta$ system

Why polarised measurements?

Assumption 1 Cross sections only differ from phase space because of FSI

⇒ A, B show same energy dependence

Assumption 2 Cross sections depend on spin dependent production amplitudes (no FSI effect)

⇒ A, B might show different energy dependence

How to separate $|A|^2$ and $|B|^2$?

⇒ Tensor analysing power t_{20} :

$$t_{20}(p_f) = \sqrt{2} \frac{|B|^2 - |A|^2}{|A|^2 + 2|B|^2}$$

Why polarised measurements?

Assumption: $|A|^2$ and $|B|^2$ only depend on FSI

$$|A|^2 = |A_0|^2 \cdot \text{FSI}(p_f)$$

$$|B|^2 = |B_0|^2 \cdot \text{FSI}(p_f)$$

$$\Rightarrow t_{20}(p_f) = \sqrt{2} \cdot \frac{|B_0|^2 - |A_0|^2}{|A_0|^2 + 2|B_0|^2} \cdot \frac{\text{FSI}(p_f)}{\text{FSI}(p_f)} = \text{const.}$$

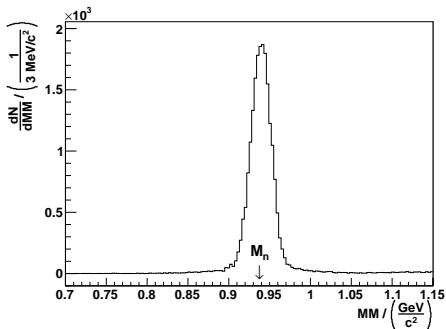
Measure t_{20} in collinear kinematics

$$t_{20}(p_f) = \frac{2\sqrt{2}}{p_{zz}} \cdot \left(\frac{\frac{d\sigma^0}{d\Omega}(\vartheta) - \frac{d\sigma^\uparrow}{d\Omega}(\vartheta)}{\frac{d\sigma^0}{d\Omega}(\vartheta)} \right)$$

Determining ρ_{zz}

Investigate reaction $d + p \rightarrow (pp) + n$ with known analyzing powers at lower energy ($T_d = 1.2$ GeV)

$$\frac{d\sigma^\uparrow}{dt}(q, \varphi) / \frac{d\sigma^0}{dt}(q) = 1 + \sqrt{3}\rho_{zz}t_{11}(\vartheta) \cos \varphi - \frac{1}{2\sqrt{2}}\rho_{zz}t_{20}(\vartheta) - \frac{\sqrt{3}}{2}\rho_{zz}t_{22}(\vartheta) \cos 2\varphi$$



Determining t_{20}

- ▶ Reminder:

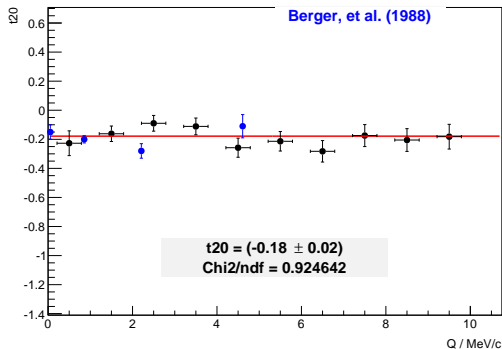
$$t_{20}(p_f) = \frac{2\sqrt{2}}{p_{zz}} \cdot \left(\frac{\frac{d\sigma^0}{d\Omega}(\vartheta) - \frac{d\sigma^\uparrow}{d\Omega}(\vartheta)}{\frac{d\sigma^0}{d\Omega}(\vartheta)} \right) \quad \vartheta = 0^\circ, 180^\circ$$

- ▶ Hence:

$$t_{20}(p_f) = \frac{2\sqrt{2}}{p_{zz}} \left(1 - \frac{N^\uparrow}{N^0} \right)$$

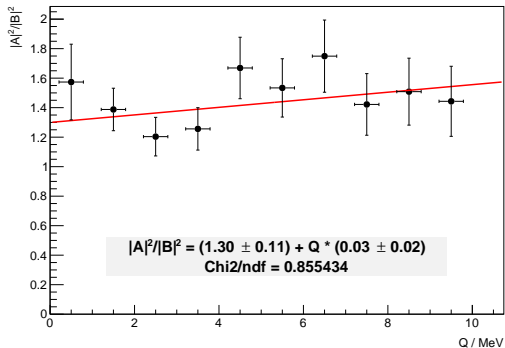
N^\uparrow, N^0 are normalized counts

Current status



- ▶ Data on t_{20} consistent with Berger, et al.
- ▶ No energy dependence visible within uncertainties

Current status



- ▶ Allow slope to put upper limit on spin dependent contributions

Summary

- ▶ Studied reaction $\vec{d} + p \rightarrow {}^3\text{He} + \eta$ in energy range $Q = -5 \text{ MeV}$ to 15 MeV relative to η -threshold
 - ▶ No energy dependence of t_{20} observed, t_{20} consistent with Berger, et al.
 - ▶ Amplitudes A, B show same energy dependence within uncertainties
 - ▶ Ratio of $|A|^2/|B|^2$ will allow to quantify spin dependent contributions
- Threshold behavior of $d + p \rightarrow {}^3\text{He} + \eta$ consistent with pure FSI ansatz

Final data available soon

Thank you for your attention



Additional Slides

Production amplitude

- ▶ Two independent η s-wave amplitudes (close to threshold)
- ▶ Possible dp spin combinations: $S = \frac{3}{2}, \frac{1}{2}$
- ▶ Couples with orbital angular momentum of dp: $L_{dp} = 1$
 $\Rightarrow J = \frac{1}{2}^-$
- ▶ Alternative production amplitude used for analogous π^0 production

$$f = \bar{u}_{3\text{He}} \hat{p}_p \cdot (A \vec{\varepsilon}_d + iB \vec{\varepsilon}_d \times \vec{\sigma}) u_p$$

$\bar{u}_{3\text{He}}$:Spinors of the particles

$\vec{\varepsilon}_d$:Polarisation vector of the deuteron

\hat{p}_p :Direction of momentum

$\vec{\sigma}$:Pauli matrix

$^3\text{He}n\eta$ final state interaction

- ▶ Fit to the data for $Q < 11\text{MeV}$

$$p_1 = \left[\left(-5 \pm 7_{-1}^{+2} \right) \pm i \cdot \left(19 \pm 2 \pm 1 \right) \right] \frac{\text{MeV}}{c}$$

$$p_2 = \left[\left(106 \pm 5 \right) \pm i \cdot \left(76 \pm 13_{-2}^{+1} \right) \right] \frac{\text{MeV}}{c}$$

- ▶ Pole of the production amplitude:

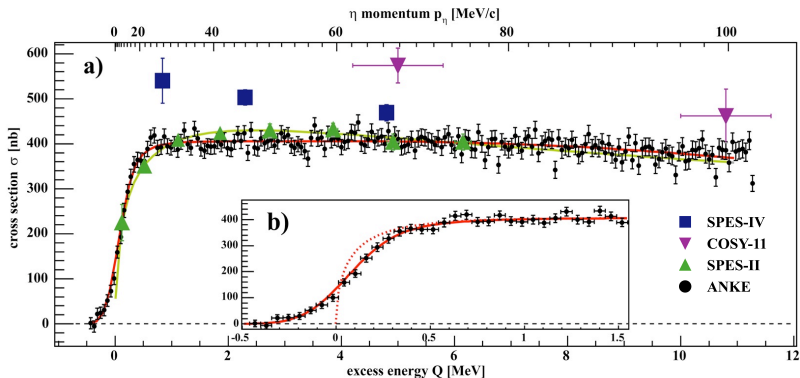
$$Q_0 = \frac{p_1^2}{2 \cdot m_{\text{red}}}$$

$$= \left[\left(-0.30 \pm 0.15 \pm 0.04 \right) \pm i \cdot \left(0.21 \pm 0.29 \pm 0.06 \right) \right] \text{MeV}$$

$$|Q_0| \approx 0.4$$

- ▶ Pole close to threshold expected from a quasi-bound state

$^3\text{He}\eta$ final state interaction



- ▶ Very good description for the whole energy range
- ▶ Unexpected large scattering length:

$$a_{^3\text{He}\eta} = \left[\pm \left(10.7 \pm 0.8_{-0.5}^{+0.1} \right) + i \cdot \left(1.5 \pm 2.6_{+1.0}^{-0.3} \right) \right] \text{ fm}$$