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Investigation of the ${}^3\text{He}-\eta$ system with polarized beams at ANKE

II International Symposium on Mesic Nuclei

September 22-25, 2013

wissen.leben
WWU Münster

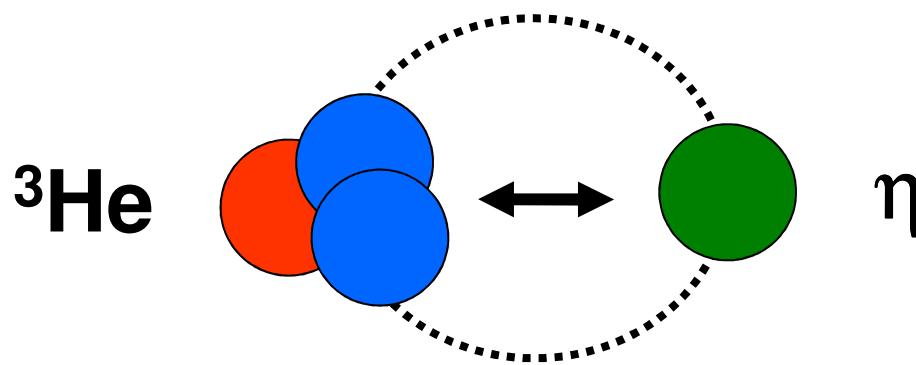
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Why η -Meson Production Close to Threshold?

- Do bound meson-nucleus systems exist?



- ANKE: $d+p \xrightarrow{\text{(→)}} ^3\text{He}+\eta$
- Excitation function close to threshold \rightarrow FSI
- Polarized beam \rightarrow Test of FSI hypothesis, role of spins



The COSY-Accelerator at Jülich



COSY (Cooler Synchrotron)

Energy range

- 0.045 – 2.8 GeV (p)
- 0.023 – 2.3 GeV (d)
(momentum 3.7 GeV/c)

Beam cooling

- Electron cooling
- Stochastic cooling

Polarisation

- p, d beams & targets

Beams

- internal, external

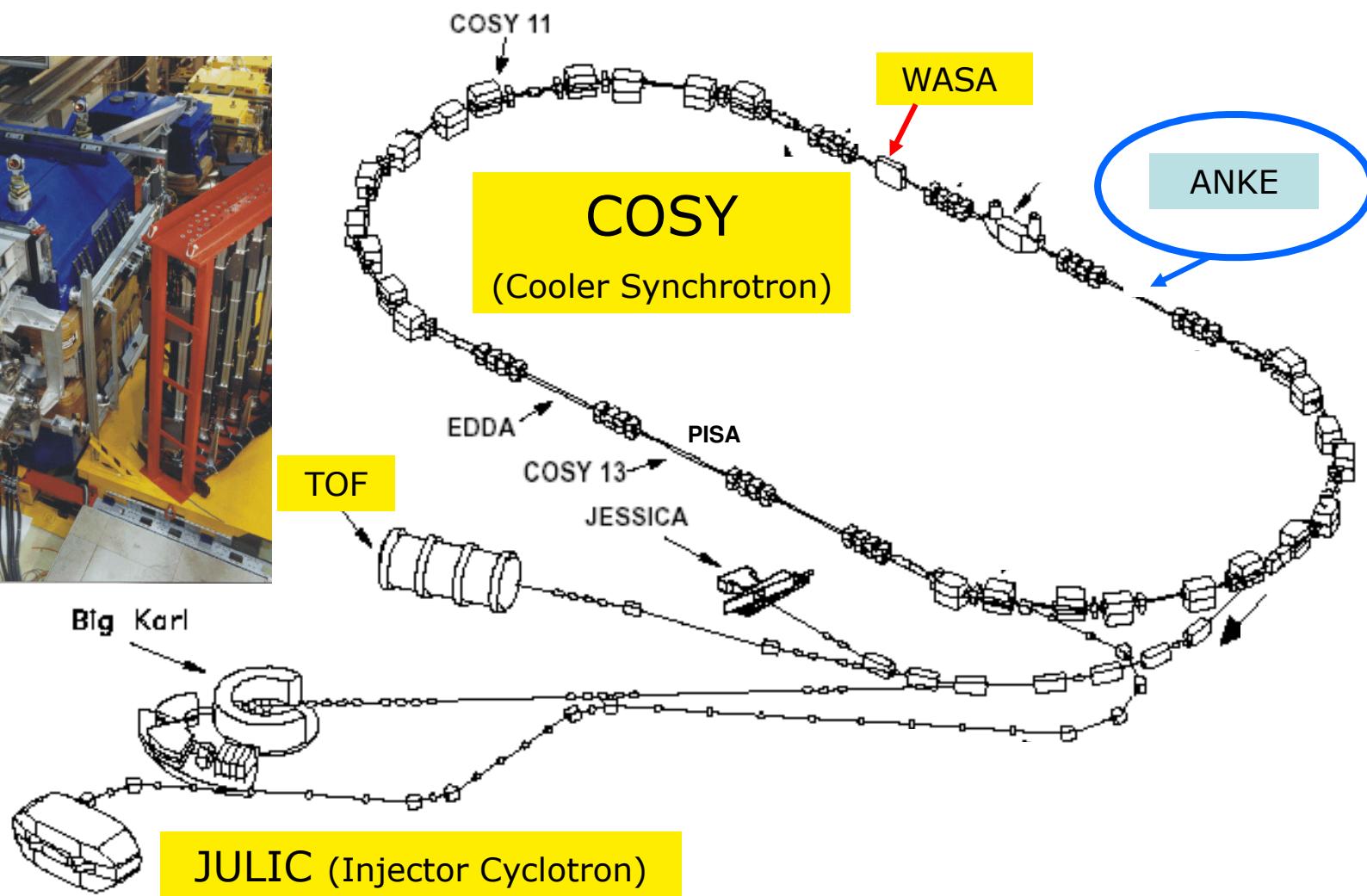
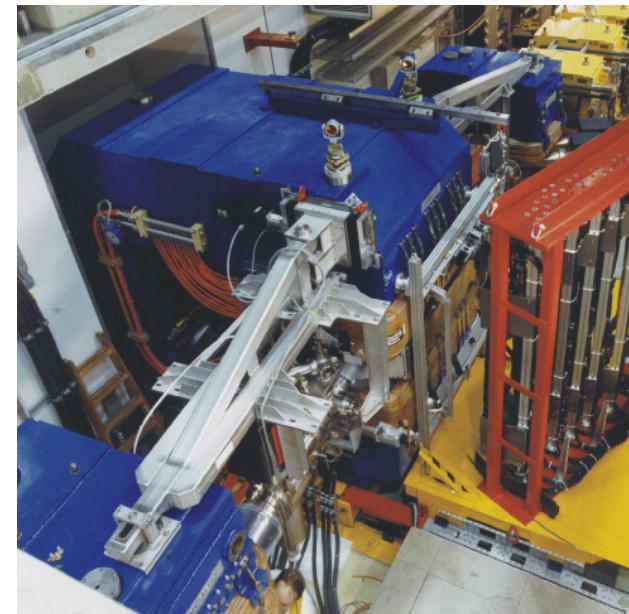
Experiments, Detectors

- ANKE, TOF, WASA, ...

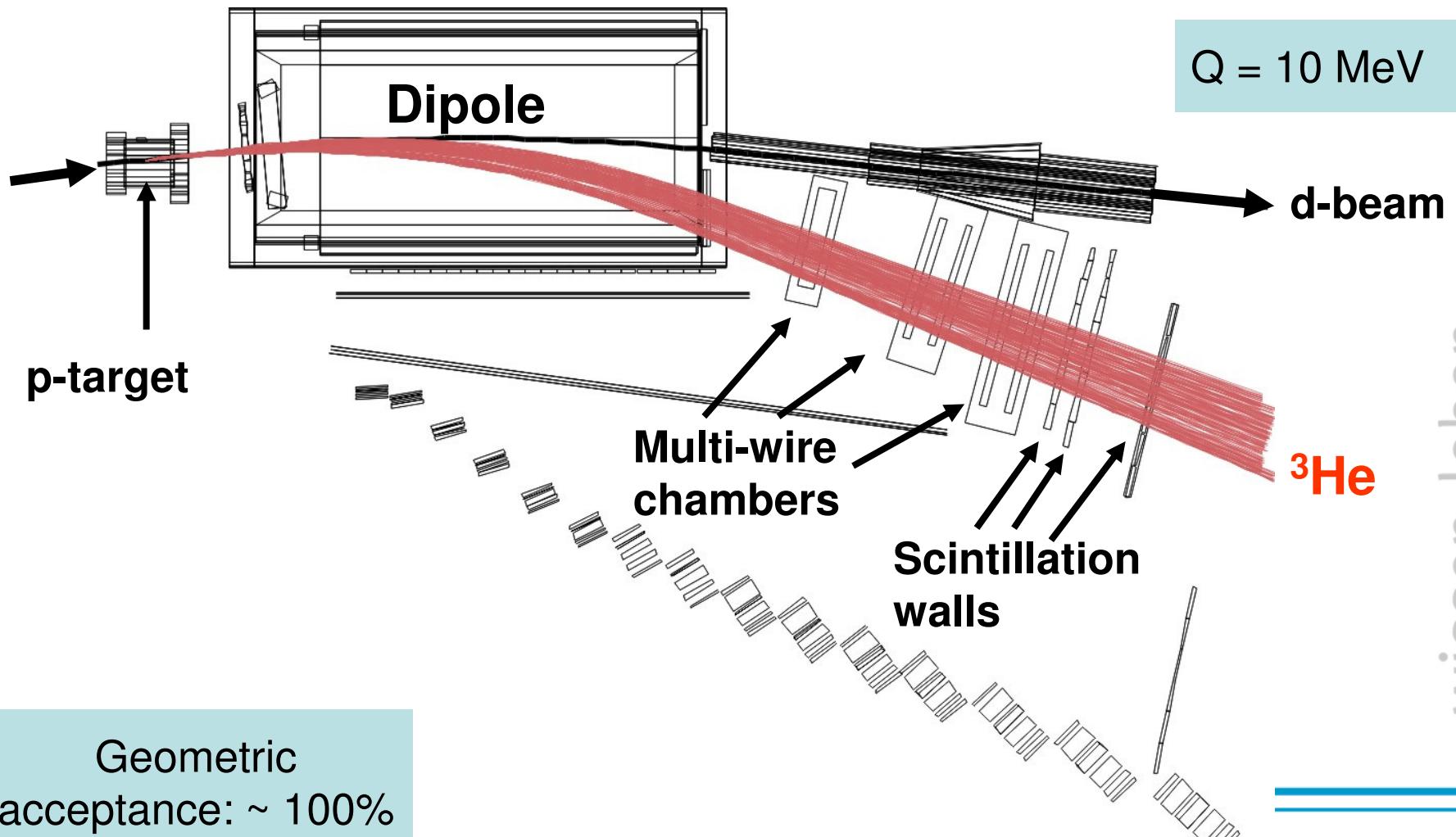


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The ANKE-Facility

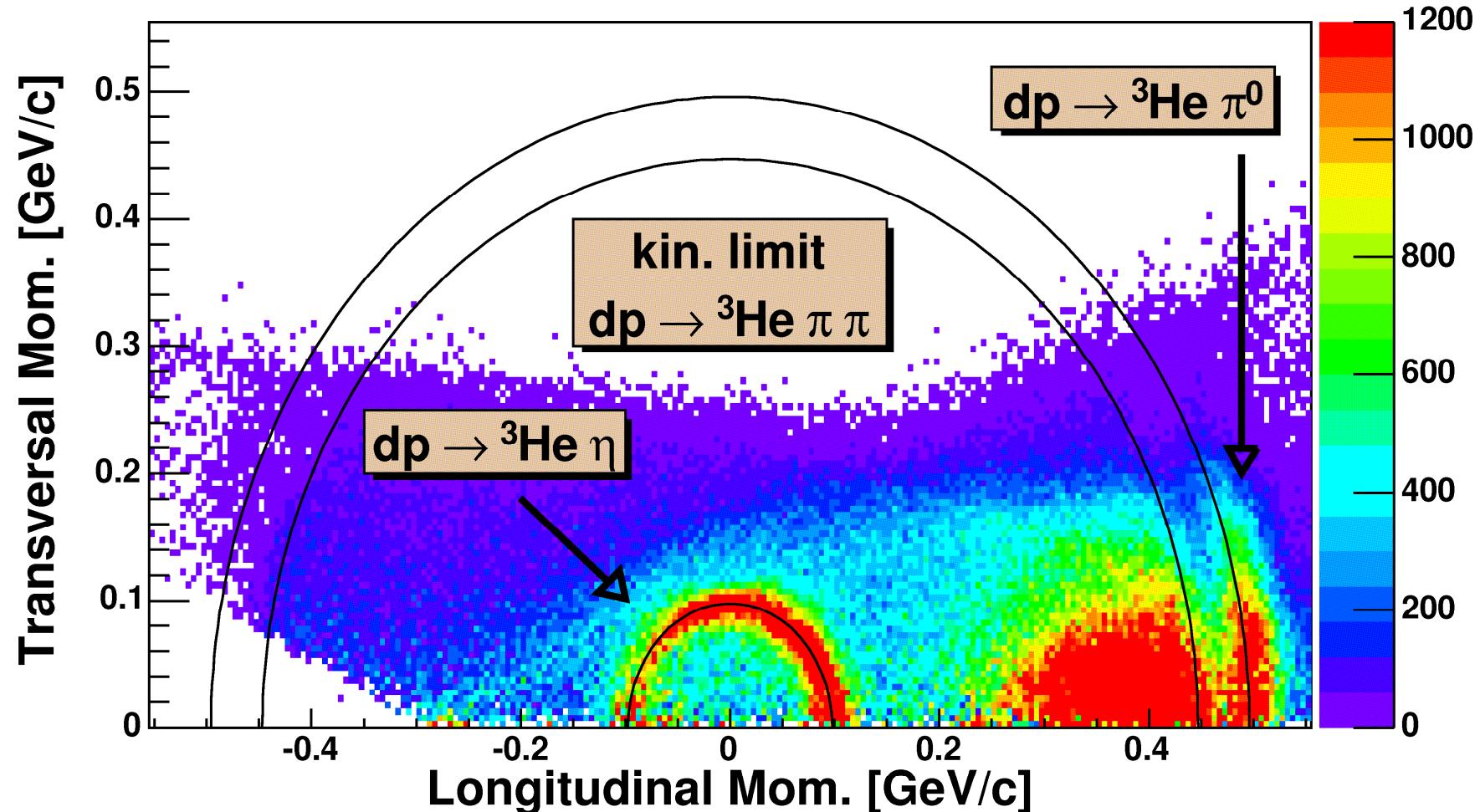


Identification of ${}^3\text{He}$ Nuclei at ANKE



Identification of the Reactions: $d+p \rightarrow {}^3\text{He}+X$

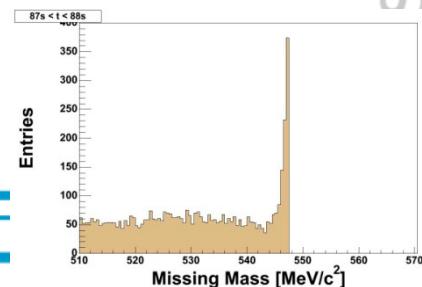
„Momentum rabbit“

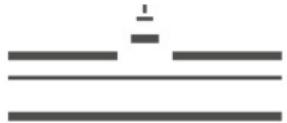




Identification of the Reactions: $d+p \rightarrow {}^3\text{He}+X$

- Energies and momenta of the incoming particles (d,p) known
 - Deuteron (mass = m_d):
energy + momentum: Adjustable by the accelerator
 - Proton (mass = m_p):
target particle at rest, momentum = 0
- Energy of the ${}^3\text{He}$ nucleus measurable by detectors
- η -meson: Not directly detectable at ANKE
 - Identification of the reaction via the missing mass analysis





Two-Particle Final State: Phase Space

Assumption:

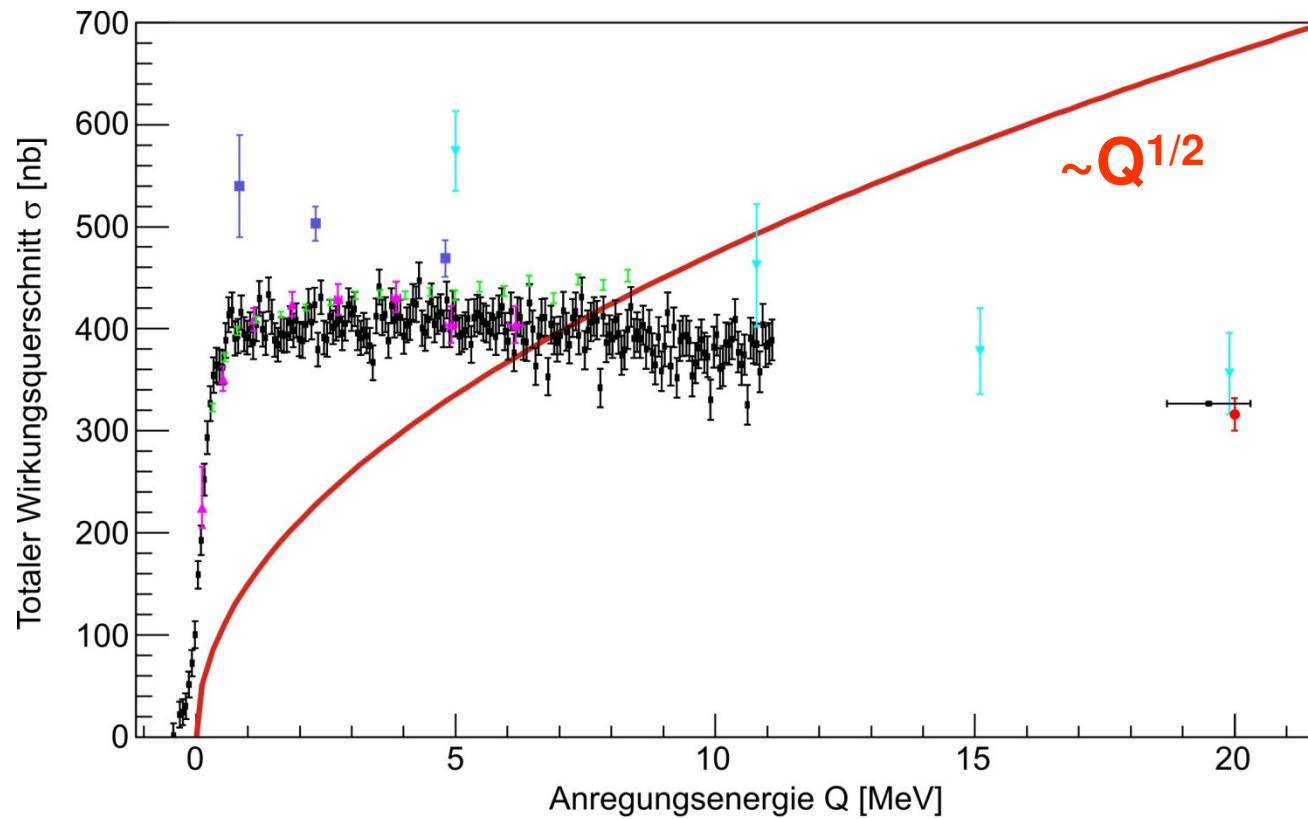
- Two-particle reaction $a+b \rightarrow c+d$ without initial and final state interactions („ISI“ and „FSI“):
- Scattering (and production) amplitude $f = \text{const.}$
→ Increase of the cross section according to phase space expectations

$$\frac{d\sigma(\vartheta)}{d\Omega} = \frac{p_f}{p_i} |f_s|^2 \propto p_f \propto \sqrt{Q}$$

p_i / p_f : Momenta of in- and outgoing particles in the CMS

Q: Q-value = Sum of kinetic energies im CMS

Results for the Reaction $d+p \rightarrow {}^3\text{He}+\eta$



But:

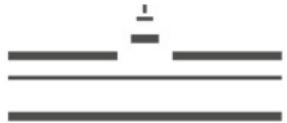
- Strong deviation from phase space expectation!
- Most probably not caused by higher partial waves



The Reaction $d+p \rightarrow {}^3\text{He}+\eta$

- Extreme increase of the total cross section close to the production threshold
- Increase of the cross sections within $\Delta Q < 1 \text{ MeV}$
 - strong energy dependence at threshold
- After that total cross sections remain almost constant
 - Additional effect beside pure phase space

Explanation: Strong final state interaction (FSI) between ${}^3\text{He}$ nucleus and η -meson



Scattering Theory and Final State Interaction

Description of the cross section including FSI:

$$\frac{d\sigma(\vartheta)}{d\Omega} = \frac{p_f}{p_i} |f_s|^2 = \frac{p_f}{p_i} \cdot \frac{|f_{\text{prod}}|^2}{\left|1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r_0 \cdot p_f^2\right|^2}$$

Assumption:

- Energy dependence of the production amplitude f_{Prod} is negligible close to threshold: $f_{\text{Prod}} \sim \text{const.}$
- Initial State Interaction (ISI) also: $\text{ISI} = \text{const.}$



Scattering Theory and Final State Interaction

- The scattering length can deliver informationen about possible bound states
- Conditions for bound $\eta^3\text{He}$ state:
 - Existence of a pole in the complex p_f plane

$$f_s = \frac{f_{\text{prod}}}{1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r \cdot p_f^2}$$

$$a \equiv a_r + ia_i$$

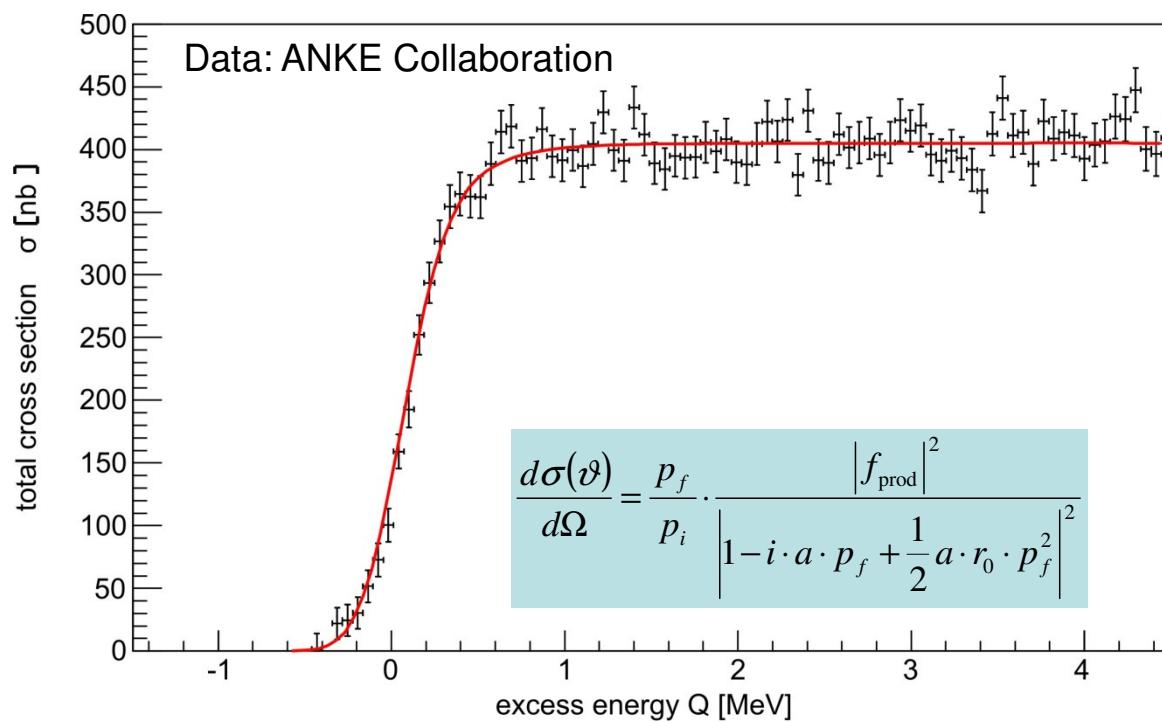
$$r \equiv r_r + ir_i$$

- As well as

$$a_r < 0, \quad a_i > 0, \quad R = \frac{|a_i|}{|a_r|} < 1$$

The Reaction $d+p \rightarrow {}^3\text{He}+\eta$

Fit to data very close to threshold: Only s-wave



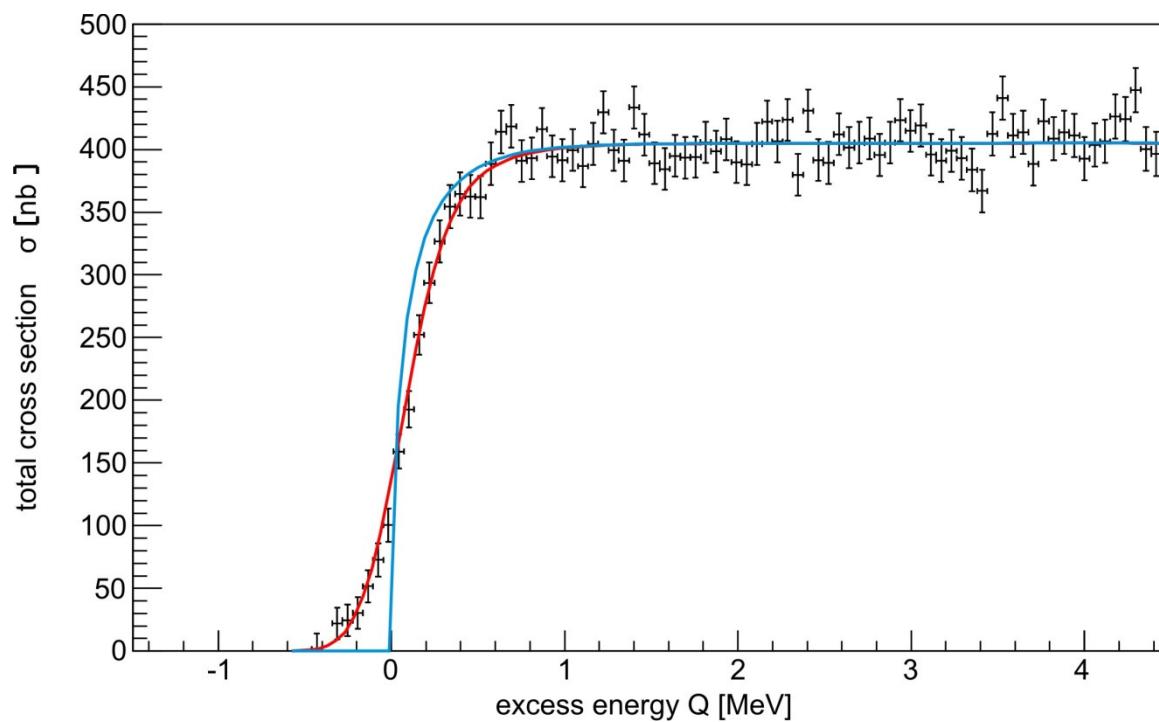
Fit parameter:

- Complex scattering length $a=a_r+ia_i$
- Complex effective range $r=r_r+ir_i$
- Finite momentum width δp_{beam} of the accelerator beam



The Reaction $d+p \rightarrow {}^3\text{He}+\eta$

Excitation function without accelerator beam smearing δp_{beam} :

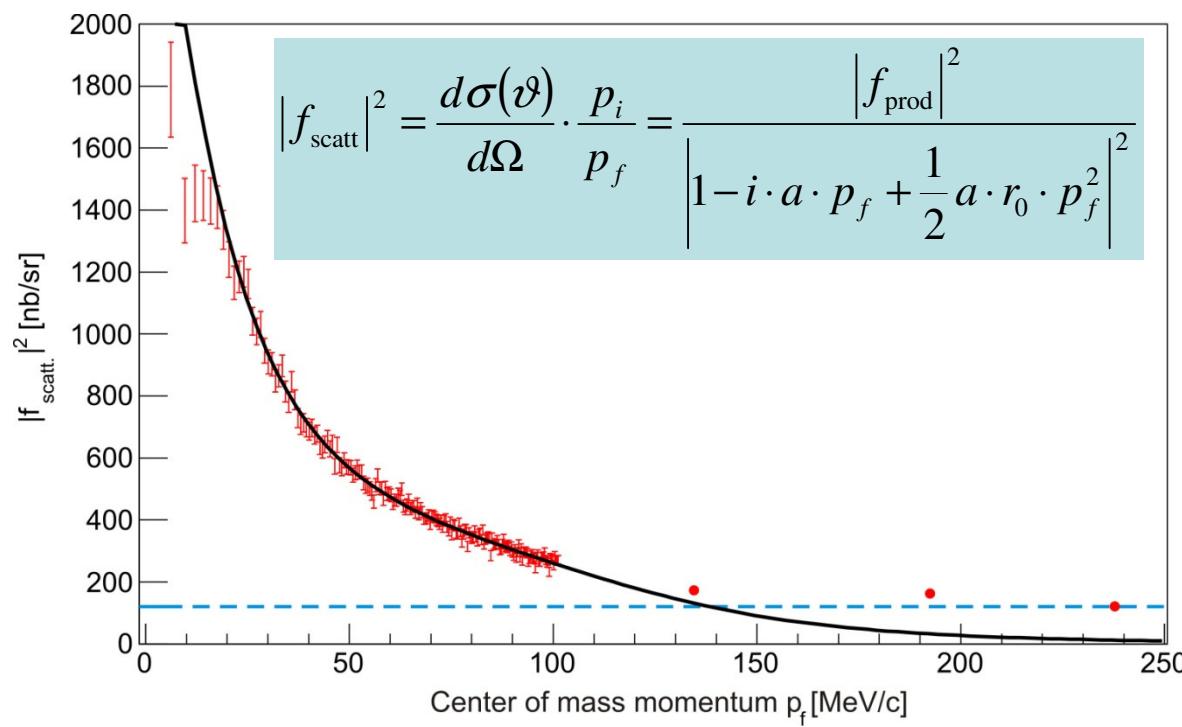


Blue line:

- Defolded shape, extracted from data (no accelerator beam smearing)
→
- Total cross section reaches maximum already $\Delta Q < 0.5$ MeV above threshold

The $d+p \rightarrow {}^3\text{He}+\eta$ Scattering Amplitude

Extracted scattering amplitude ($Q > 0$ MeV)



- Scattering amplitude decreases rapidly with increasing final state momentum p_f
- Scattering amplitude almost constant at high energies

→ strong FSI in $\eta {}^3\text{He}$ system



η - ${}^3\text{He}$ Scattering Length

Fit to data delivers information about the complex η - ${}^3\text{He}$ scattering length:

$$\left(\frac{d\sigma(\vartheta)}{d\Omega} \right) \cdot \frac{p_i}{p_f} = |f_{\text{scat}}|^2 = |f_{\text{prod}} \cdot FSI|^2 = |f_{\text{prod}}|^2 \cdot |FSI|^2$$



Result:

$$a = [\pm (10.7 \pm 0.8^{+0.1}_{-0.5}) + i(1.5 \pm 2.6^{+1.0}_{-0.9})] \text{ fm}$$

$$FSI = \frac{1}{1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r_0 \cdot p_f^2}$$

Notice: Determination of $|a_r|$!



η - ^3He -Interaction: Determination of Pols

$$\left(\frac{d\sigma(\vartheta)}{d\Omega} \right) \cdot \frac{p_i}{p_f} = |f_{\text{scatt}}|^2 = |f_{\text{prod}} \cdot FSI|^2 = |f_{\text{prod}}|^2 \cdot |FSI|^2$$

$$FSI = \frac{1}{1 - i \cdot a \cdot p_f + \frac{1}{2} a \cdot r_0 \cdot p_f^2}$$

$$FSI = \frac{1}{\left(1 - \frac{p_f}{p_1}\right) \cdot \left(1 - \frac{p_f}{p_2}\right)}$$

$$a = -i \cdot \frac{p_1 + p_2}{p_1 \cdot p_2} \quad r_0 = + \frac{2 \cdot i}{p_1 + p_2}$$

$$p_1 = \left[(-5 \pm 7^{+2}_{-1}) \pm i \cdot (19 \pm 2 \pm 1) \right] \text{MeV/c}$$
$$p_2 = \left[(106 \pm 5) \pm i \cdot (76 \pm 13^{+1}_{-2}) \right] \text{MeV/c}$$

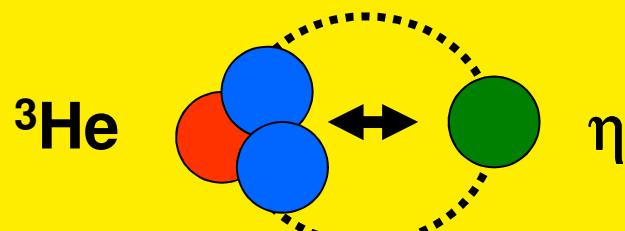
η - ^3He -Interaction: Determination of Pols

- Pole close to the reaction threshold

$$|Q_0| = \left| \frac{p_1^2}{2 \cdot m_{red}} \right| = 0.37 \text{ MeV}$$

- Position of the near-threshold pole (and scattering length) stable, i.e. nearly independent of fit range
- Large real part of scattering length and $|a_r| > a_i$

→ indication for the existence of a bound state
(strong interaction!)





Polarized Measurements

Production amplitude for $d\bar{p} \rightarrow {}^3\text{He} + \eta (\pi^0)$:

$$f_B = \bar{u}_\tau \vec{p}_p \cdot (A \vec{\epsilon}_d + iB \vec{\epsilon}_d \times \vec{\sigma}) u_p$$

see:
C. Kerboul et al.,
Phys. Lett. B 181, 28 (1986)

Determination of the energy dependence of the **amplitudes A** and **B** by measurement of:

$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_\eta}{p_p} \left[|A|^2 + 2|B|^2 \right]$$

$$|A|^2 = \frac{p_p}{p_\eta} (1 - \sqrt{2} T_{20}) \frac{d\sigma}{d\Omega}$$

$$T_{20} = \sqrt{2} \left[\frac{|B|^2 - |A|^2}{|A|^2 + 2|B|^2} \right]$$

$$|B|^2 = \frac{p_p}{p_\eta} (1 + \frac{1}{\sqrt{2}} T_{20}) \frac{d\sigma}{d\Omega}$$

$$T_{20} = \frac{2 \cdot \sqrt{2}}{p_{zz}} \cdot \frac{d\sigma_0/d\Omega(\vartheta) - d\sigma_\uparrow/d\Omega(\vartheta)}{d\sigma_0/d\Omega(\vartheta)} \quad \vartheta = 0^\circ \text{ or } 180^\circ$$



Polarized Measurements

Assumption: $\vec{d}p \rightarrow {}^3\text{He} + \eta$

- Negligible effect of ISI
- Energy dependence of $|f|^2$ only given by FSI
 - Shape of excitation function independent of spins
 - Same energy dependence of amplitudes $|A|^2$ and $|B|^2$

$$|A|^2 = |A_0|^2 \cdot FSI(p_\eta)$$

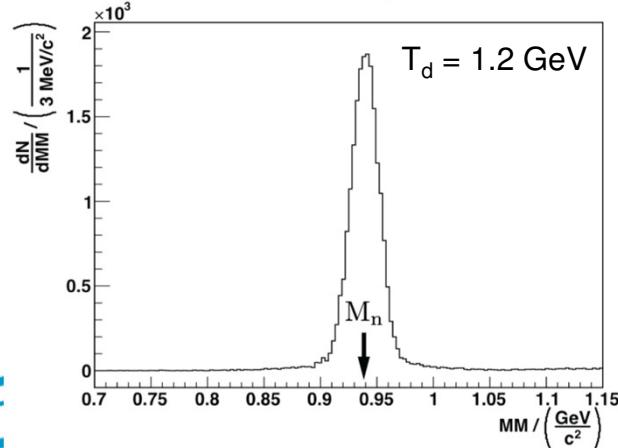
$$|B|^2 = |B_0|^2 \cdot FSI(p_\eta)$$

$$\Rightarrow T_{20} = \sqrt{2} \left[\frac{|B_0|^2 - |A_0|^2}{|A_0|^2 + 2|B_0|^2} \right] \cdot \frac{FSI(p_\eta)}{FSI(p_\eta)} = \text{const.}$$

- Measure T_{20} as function of the excess energy

The Reaction $d+p \rightarrow {}^3\text{He}+\eta$ at ANKE

- Alternating injection of unpolarized and tensor polarized deuterons in COSY
- Ramped COSY beam: $Q = -5 \text{ MeV} \dots +10 \text{ MeV}$ (300 s)
- Full geometrical acceptance of ANKE for $d+p \rightarrow {}^3\text{He}+\eta$
- Determination of p_{zz} by, e.g., $d+p \rightarrow (\text{pp})+n$ (analyzing powers known)

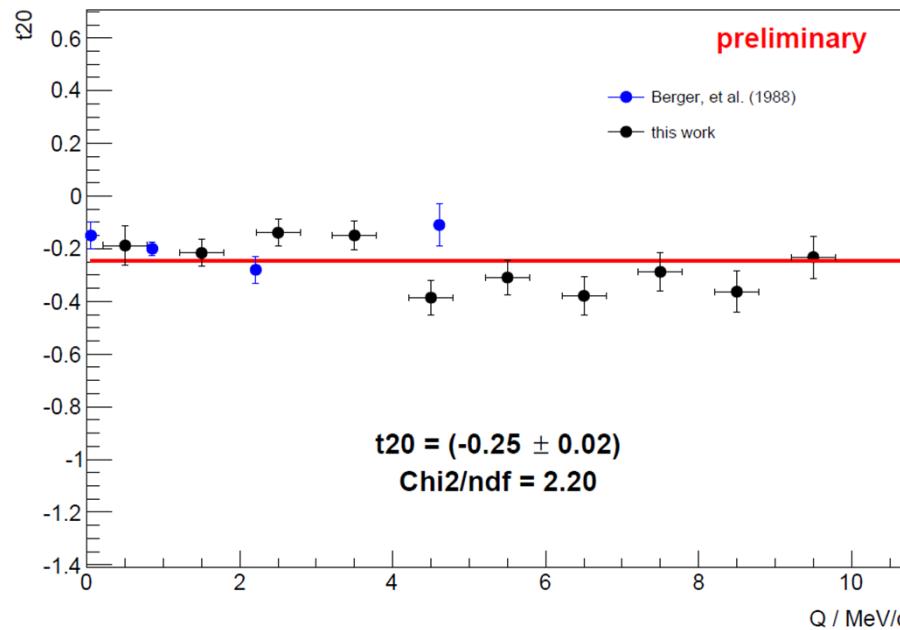


$$\frac{d\sigma_\uparrow}{dt}(q, \varphi) / \frac{d\sigma_0}{dt}(q, \varphi) =$$

$$1 + \sqrt{3} p_z i t_{11}(\vartheta) \cos(\varphi) - \frac{1}{2\sqrt{2}} p_{zz} t_{20}(\vartheta)$$

$$- \frac{\sqrt{3}}{2} p_{zz} t_{22}(\vartheta) \cos(2\varphi)$$

Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$



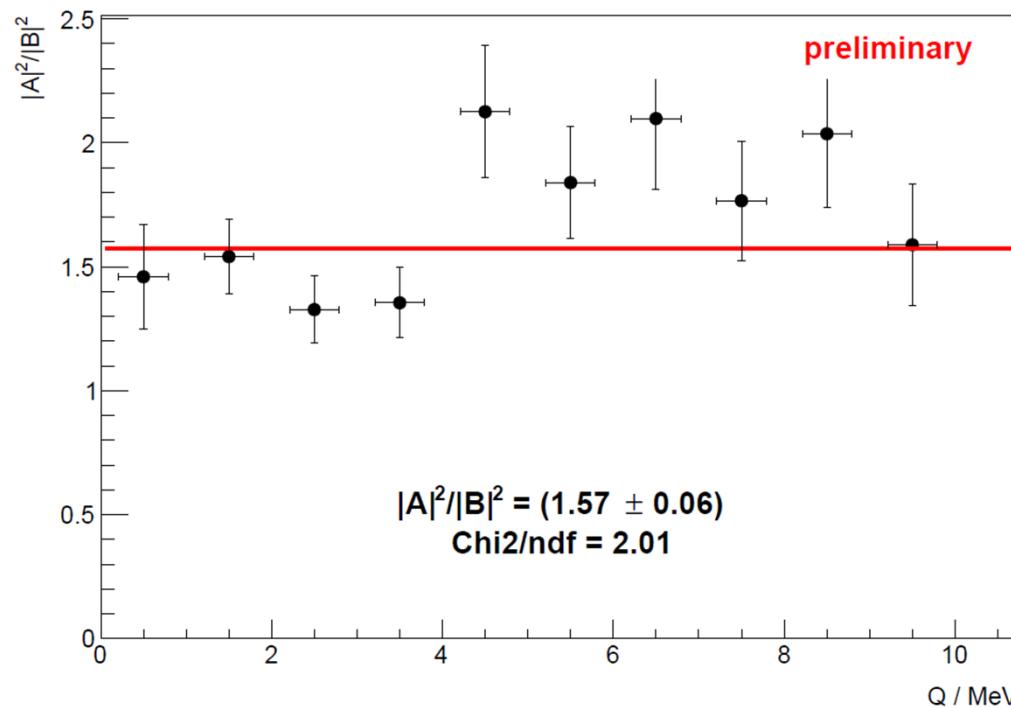
$$T_{20} = \frac{2 \cdot \sqrt{2}}{p_{zz}} \cdot \frac{d\sigma_0/d\Omega(\vartheta) - d\sigma_\uparrow/d\Omega(\vartheta)}{d\sigma_0/d\Omega(\vartheta)}$$

- Data are consistent with $T_{20} = \text{const.}$ close to threshold
- Extraction of $|A|^2/|B|^2$:

$$T_{20} = \sqrt{2} \left[\frac{|B|^2 - |A|^2}{|A|^2 + 2|B|^2} \right] \rightarrow \frac{|A|^2}{|B|^2} = \frac{1 - \sqrt{2} \cdot T_{20}}{1 + T_{20}/\sqrt{2}}$$

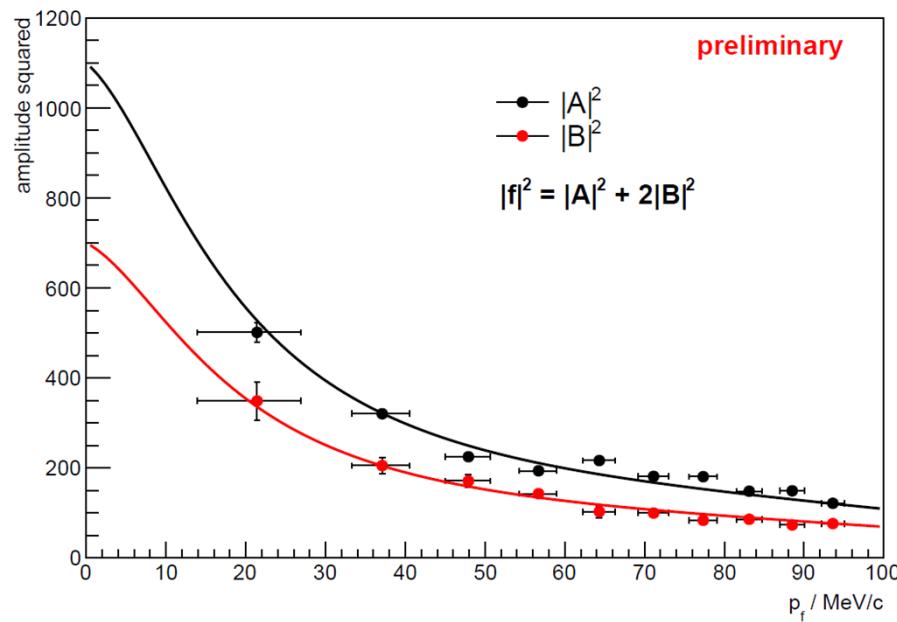
Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

- Assumption: $T_{20} = \text{const.} \rightarrow |A|^2/|B|^2 = \text{const.}$



Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

- Energy dependence of $|f|^2$ known from „old“ unpolarized measurements
 $\rightarrow |A|^2(p_f)$ and $|B|^2(p_f)$ can be calculated



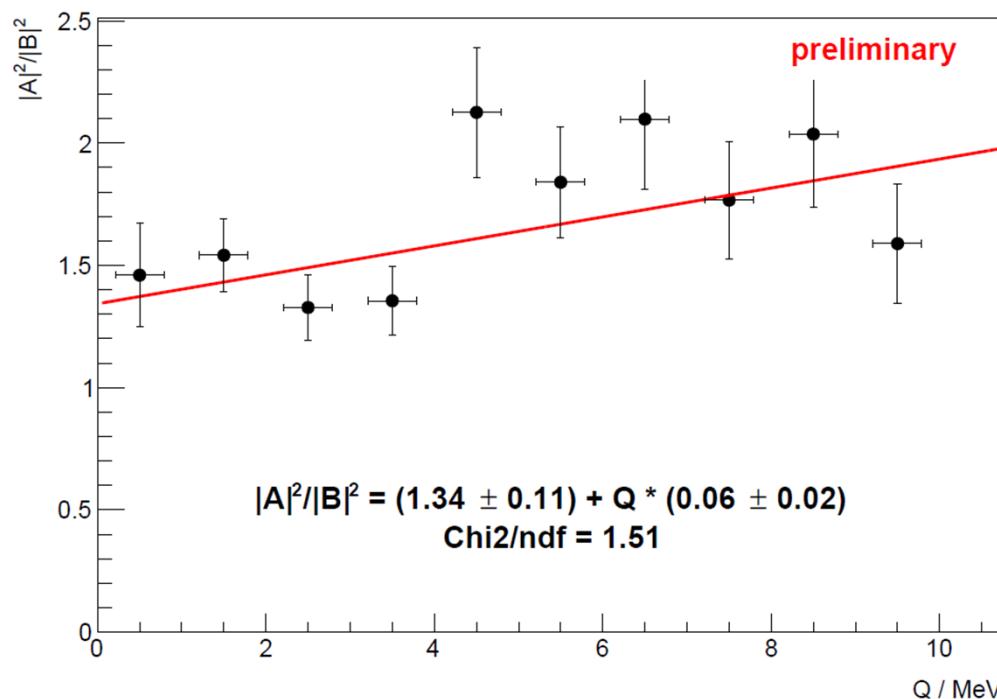
$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_\eta}{p_p} \left[|A|^2 + 2|B|^2 \right]$$

$$|A|^2 = \frac{p_p}{p_\eta} (1 - \sqrt{2} T_{20}) \frac{d\sigma}{d\Omega}$$

$$|B|^2 = \frac{p_p}{p_\eta} \left(1 + \frac{1}{\sqrt{2}} T_{20} \right) \frac{d\sigma}{d\Omega}$$

Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

- Allow for an energy dependence of $|A|^2/|B|^2$:
 → Test: Different energy dependence of $|A|^2(p_f)$ and $|B|^2(p_f)$?

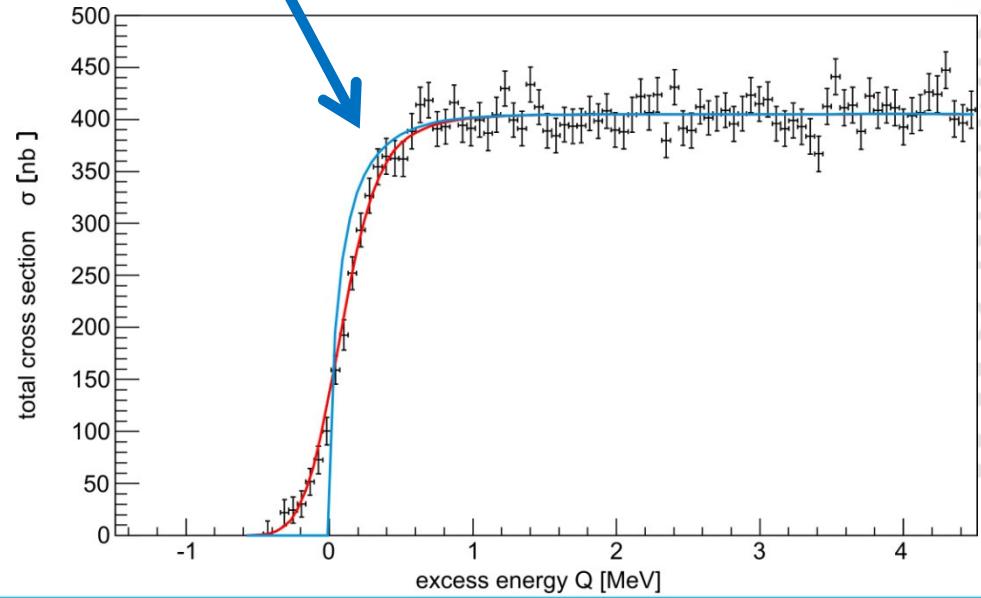
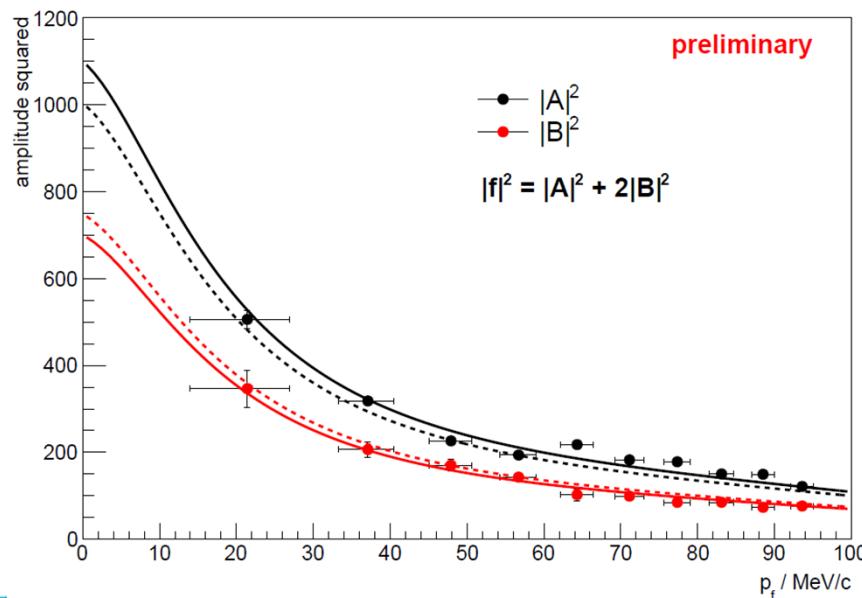


$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_\eta}{p_p} \left[|A|^2 + 2|B|^2 \right]$$

$$\frac{|A|^2}{|B|^2} = m \cdot Q + n$$

Preliminary Results: $d+p \rightarrow {}^3\text{He}+\eta$

- No significant different energy dependence of $|A|^2$ and $|B|^2$
- Remarkable excitation function of $d+p \rightarrow {}^3\text{He}+\eta$ still an indication for very strong FSI effect





Next Steps:

- Quantification of upper limits for non-FSI effect
- Evaluation of effect on pole position or scattering length

In parallel:

- Analysis of new data on $p+n \rightarrow d+\eta$ via $p+d \rightarrow d+\eta+p_{\text{spec}}$
- Comparison of results from
 - $p+n \rightarrow d+\eta$
 - $d+p \rightarrow {}^3\text{He}+\eta$
 - $d+d \rightarrow {}^4\text{He}+\eta$





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Summary

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