The η meson mass determination with ANKE at COSY*

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on behalf of the ANKE-Collaboration

Recent measurements on the η meson mass performed at different experimental facilities (i.e. CERN-NA48, COSY-GEM, CESR-CLEO, DAΦNE-KLOE, MAMI-Crystall Ball) resulted in very precise data, but differ by up to more than eight standard deviations, i.e. $0{,}5\,\mathrm{MeV/c^2}$ [1]. In order to clarify this situation a high precision measurement using the ANKE spectrometer at the COoler SYnchrotron COSY has been realized.

Using the two-body reaction $\mathrm{d\,p} \to {}^3\mathrm{He}\,\eta$ at low excess energies the η mass can be determined only from pure kinematics by the determination of the production threshold. Therefore, twelve data points at fixed excess energies in the range of $Q = 1-12\,\mathrm{MeV}$ were investigated. The final state momentum p_f of the ^{3}He-particles in the Center of Mass (CM) frame

$$p_f\,(s) = \frac{\sqrt{\left[s-(m_{^3\mathrm{He}}+m_\eta)^2\right]\cdot\left[s-(m_{^3\mathrm{He}}-m_\eta)^2\right]}}{2\cdot\sqrt{s}}\,, \quad (1)$$

measured with the ANKE spectrometer, is very sensitive on the η mass and the total energy $\sqrt{s}$, where the latter one is completely defined in a fixed target experiment by the masses of the initial particles and the momentum of the deuteron beam p_d:

$$\sqrt{s} = \sqrt{2m_p\sqrt{m_d^2+p_d^2}+m_d^2+m_p^2}\,. \quad (2)$$

For a precise determination of the production threshold both quantities, the final state momenta of the ^{3}He-particles and the corresponding deuteron beam momenta have to be measured with highest accuracy. Fitting the dependence of the final state momentum p_f on the beam momentum p_d and the η mass, $p_f = p_f(p_d, m_\eta)$, for the twelve points the mass can be extracted as fitting parameter (see Fig. 3).

The beam momentum for each fixed excess energy was determined using a method developed at VEPP-2M at Novosibirsk at an electron-positron machine [2] using the spin dynamics of a polarized beam. Here the spin precession frequency of a relativistic particle is disturbed by an artificial spin resonance induced by a horizontal rf-magnetic field of a solenoid leading to a depolarization of the polarized accelerator beam. The depolarizing resonance frequency f_r depends on the kinematical γ-factor (i.e. the beam momentum $p = m\sqrt{\gamma^2-1}$) and the beam revolution frequency f_0 via the resonance condition:

$$f_r = (k+\gamma G)\,f_0\,, \quad (3)$$

where k is an integer and G the gyromagnetic anomaly of the beam particle. By measuring these two frequencies the beam momentum of a polarized beam can be determined with a precision below $\Delta p/p < 10^{-4}$. For the first time this method was used at COSY with a vector polarized deuteron beam and the momenta in the threshold range of 3.1 - 3.2 GeV/c were determined with an accuracy of $\Delta p/p < 6\cdot10^{-5}$, i.e. with 170 keV/c [3].

The correct final state momenta for the twelve different energies of the ^{3}He-nuclei of the reaction $\mathrm{d\,p} \to {}^3\mathrm{He}\,\eta$ can only be extracted fulfilling two conditions: a clear identification of the reaction of interest and a precise detector calibration. The feature that the ANKE facility has full geometrical acceptance for the reaction $\mathrm{d\,p} \to {}^3\mathrm{He}\,\eta$ near threshold up to $Q = 15\,\mathrm{MeV}$ allows to verify and improve the detector calibration by studying the kinematics of this two-body reaction. According to Eq. 1 for a fixed CM energy $\sqrt{s}$ the final state momenta in the CM frame are distributed on a momentum sphere with constant radius $p_f = (p_x^2+p_y^2+p_z^2)^{1/2}$, which can be visualized by plotting the transversal versus the longitudinal reconstructed momentum, as shown in Fig 1a. One expects a centered momentum locus with a fixed radius p_f, indicated as dashed line. The main idea to verify the cali-

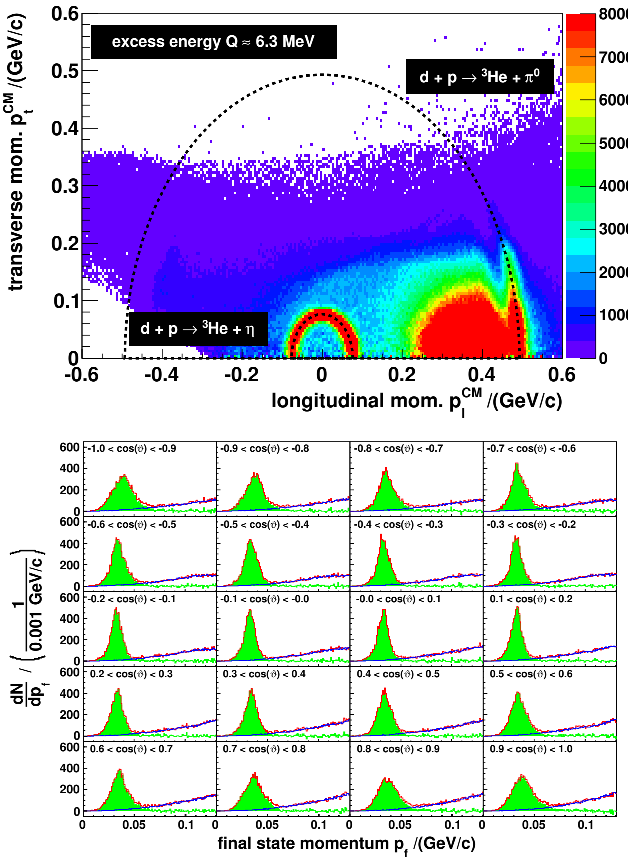

Fig. 1: **a)** The momentum loci for the ^{3}He η and ^{3}He π channels. For the ^{3}He η channel ANKE covers near threshold the full solid angle, while for the ^{3}He π channel only forward scattered ^{3}He-nuclei are detected. **b)** Final state momentum $p_f = p_f(\cos\vartheta)$ (red) at an excess energy of $Q \approx 1\,\mathrm{MeV}$, the background description (blue) and the extracted ^{3}He η signal (green).

bration is that the momentum sphere has to be completely symmetric in p_x, p_y and p_z (or ϑ and ϕ) so that the final state momentum p_f should be constant in all directions. By a careful investigation of the momentum dependence on the cosine of the polar angle ϑ and the azimuthal angle ϕ

$$p_f = p_f(\cos\vartheta) \quad (4)$$
$$p_f = p_f(\phi) \quad (5)$$

the shape of the momentum sphere or locus can be studied in more detail. Deviations from this symmetric shape will indicate the need for an improvements of the calibration. Therefore, the ^{3}Heη signal has to be extracted background free. At ANKE, the ^{3}He-nuclei produced can be identified using the energy loss and time of flight information. By this, the background consisting mainly of protons and deuterons of the dp elastic scattering and the deuteron break-up, can be suppressed effectively. The remaining background, originating mainly from the multi pion production can be subtracted by data taken below the η threshold at an excess energy of $Q \approx -5$ MeV, but analyzed as if they were taken above. In [4] the successful applicability of this approach on missing mass spectra is shown, but it is also applicable to final state momentum spectra as shown in Fig. 1b for different $\cos\vartheta$ bins for the data point closest to threshold, i.e. Q $\approx$ 1 MeV. Similarly the final state momentum dependency on ϕ can be studied. The background free ^{3}Heη distributions in $\cos\vartheta$ and ϕ allow one to extract the mean final state momenta for individual $\cos\vartheta$ and ϕ bins as shown in Fig. 2. In contrast to Monte

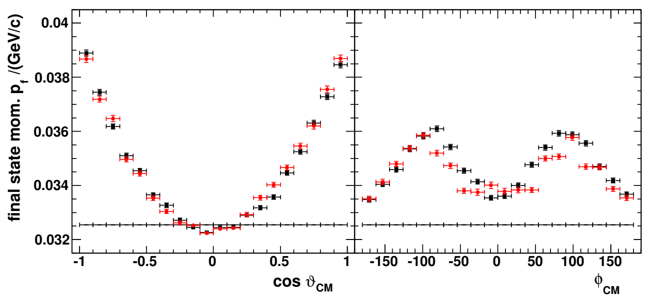

Fig. 2: Final state momentum dependence on $\cos\vartheta$ and ϕ: $p_f = p_f(\cos\vartheta)$ and $p_f(\phi)$ for data (red circles), Monte Carlo simulations without momentum smearing (black line) and with momentum smearing (black squares).

Carlo simulations on $p_f(\cos\vartheta)$ and $p_f(\phi)$ without momentum smearing (black line) the data points obtained, indicated as red crosses show a dependence of the final state momentum on $\cos\vartheta$ and ϕ. The shape of the momentum sphere is stretched to values $\cos\vartheta \to \pm 1$ and shows an oscillation in ϕ. This behavior is caused by a kinematic effect due to different momentum resolutions of the ANKE detector for p_x, p_y and p_z. Assuming that the p_x, p_y and p_z distributions are gaussian distributed with different widths σ, it is possible to reproduce the final state momentum dependence on $\cos\vartheta$ and ϕ with Monte Carlo simulations shown as black crosses.
For the final state momentum determination the plots shown in Fig. 2 are of high importance for the following three reasons.

1. Improvement of the calibration. Asymmetric shapes in $\cos\vartheta$ and ϕ originated by a "unpolished" calibration can be corrected by minor changes of the ANKE calibration parameters. These changes are so small that they have no impact on typical calibration quantities like missing masses of different reactions. That means that by studying the $\cos\vartheta$ and ϕ dependences the sensitivity for the different calibration parameter can be increased.
2. Extraction of the correct momentum resolution in (p_x, p_y, p_z). Assuming gaussian distributions with different widths we found for ANKE $(\sigma_x, \sigma_y, \sigma_z) = (3.2, 7.8, 16.4)$ MeV/c.
3. Correction of the reconstructed final state momentum. Because of the finite momentum resolution, the extracted average over all $\cos\vartheta$ and ϕ bins (see Fig. 2 red circles) is larger compared to the original one (black line). This has to be considered in the determination of the twelve final state momenta.

Currently the final state momentum analysis is still in progress but already now the twelve momenta in the range of $30-100\,\mathrm{MeV/c}$ can be determined with a precision better than $400\,\mathrm{keV/c}$. Fitting the $p_f = p_f(p_d, m_\eta)$ dependence, as shown in Fig. 3 the η mass is preliminarily determined to be

$$m_\eta = (547.869 \pm 0.007 \pm 0.040)\,\mathrm{MeV/c^2}. \quad (6)$$

The accuracy, which will be achieved at ANKE, will be comparable and competitive to the precision achieved at other recent experiments.

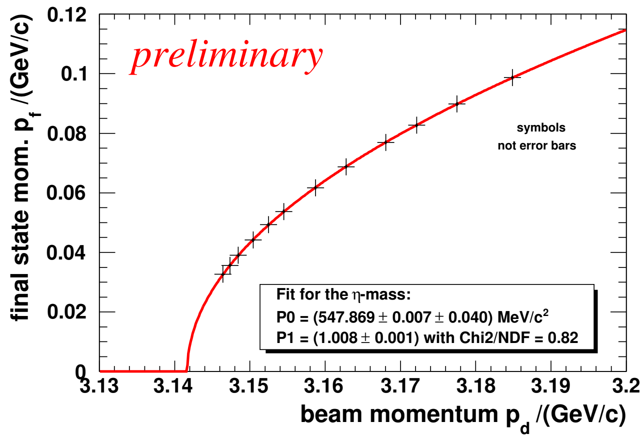

Fig. 3: Preliminary determination of the η mass by fitting the dependence $p_f = p_f(p_d, m_\eta)$. The η mass is extracted as the fitting parameter to be $m_\eta = (547.869 \pm 0.007 \pm 0.040)\,\mathrm{MeV/c^2}$.

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