

## Investigation of the ${}^3\text{He}\eta$ final state in polarized dp-collisions at ANKE

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Measurements at ANKE on the reaction  $d + p \rightarrow {}^3\text{He} + \eta$  close to the production threshold have yielded a rich and precise data set on total and differential cross sections [1]. The interpretation of these data indicated a very strong final state interaction (FSI) and the existence of a possible bound state of the  ${}^3\text{He}\eta$  system. With only unpolarised data, however, it is not possible to quantify possible contributions from the initial state such as angular momenta to the shape of the total cross section. In order to further investigate this situation, measurements with a polarized deuteron beam have been performed at ANKE covering the same excess energy range of  $Q = 0$  MeV to 11 MeV.

Close to threshold there are two independent s-wave amplitudes for the  $\vec{d} + p \rightarrow {}^3\text{He} + \eta$  reaction which can be thought of as spin- $\frac{3}{2}$  or spin- $\frac{1}{2}$  combinations of the  $dp$  spins coupled with  $L_{dp} = 1$  to give the overall  $J^P = \frac{1}{2}^-$  state. An alternative decomposition was given by Germond & Wilkin [3]

$$f = \bar{u}_3 \hat{p}_p \cdot (A \vec{\epsilon}_d + iB \vec{\epsilon}_d \times \vec{\sigma}) u_p. \quad (1)$$

With these two amplitudes  $A$  and  $B$  the differential cross section can be expressed by:

$$\frac{d\sigma}{d\Omega} = \frac{1}{3} \frac{p_f}{p_i} [|A|^2 + 2 \cdot |B|^2] \quad (2)$$

From Eq. 2 it is obvious that, if the energy dependence of  $A$  and  $B$  differs from each other, which would indicate further effects than only final state interactions, the shape of the total cross section would be affected as well. In that case, the pure FSI ansatz presented in [1] would have to be expanded to account for the contributions from the different initial states. The moduli-squared of the two amplitudes  $A$  and  $B$  can be separated by the determination of the spherical tensor analysing power  $t_{20}$ :

$$t_{20} = \sqrt{2} \frac{|B|^2 - |A|^2}{|A|^2 + 2|B|^2} \quad (3)$$

It can be shown that if  $A$  and  $B$  have the same energy dependence, which would be the case if the initial state does not influence the energy dependence of the excitation function, then  $t_{20}$  is expected to be independent of  $Q$ .

In order to determine  $t_{20}$  in collinear kinematics,

$$t_{20} = \frac{2\sqrt{2}}{p_{zz}} \cdot \left( \frac{\frac{d\sigma^0}{d\Omega}(\theta) - \frac{d\sigma^\uparrow}{d\Omega}(\theta)}{\frac{d\sigma^0}{d\Omega}(\theta)} \right) \quad \theta = 0^\circ, 180^\circ, \quad (4)$$

the tensor polarization  $p_{zz}$  has to be known. For this purpose the *polarization export* was used where additional data below the  $\eta$  production threshold have been taken at a deuteron momentum  $p_d = 2.435$  GeV/c. By the investigation of different reactions with known analysing powers it is possible to determine  $p_{zz}$ . One suitable reaction for the determination of  $p_{zz}$  is  $d + p \rightarrow {}^3\text{He} + \pi^0$  where the differential cross section for forward scattered  ${}^3\text{He}$  particles depends only on  $p_{zz}$ :

$$\frac{d\sigma^\uparrow}{d\Omega}(\theta, \phi) / \frac{d\sigma^0}{d\Omega}(\phi) = 1 + \frac{1}{2} p_{zz} A_{yy} \quad \theta = 0^\circ \quad (5)$$

The diproton production  $d + p \rightarrow (pp) + n$  is influenced by  $p_z$  and  $p_{zz}$ , though it was shown that the  $p_z$  part vanishes for small excitation energies of the diproton ( $E_{pp} < 3$  MeV) [4].

$$\frac{d\sigma^\uparrow}{dt}(q, \phi) / \frac{d\sigma^0}{dt}(q) = 1 + \sqrt{3} p_z i t_{11}(\theta) \cos \phi - \frac{1}{2\sqrt{2}} p_{zz} t_{20}(\theta) - \frac{\sqrt{3}}{2} p_{zz} t_{22}(\theta) \cos 2\phi \quad (6)$$

After the determination of  $p_{zz}$ , missing mass spectra of  $d + p \rightarrow {}^3\text{He} + \eta$  are analysed to determine the ratio of differential cross section and by this  $t_{20}$ . Using the same background description as [1] yields clean signals of the  $\eta$  production (Fig. 1).

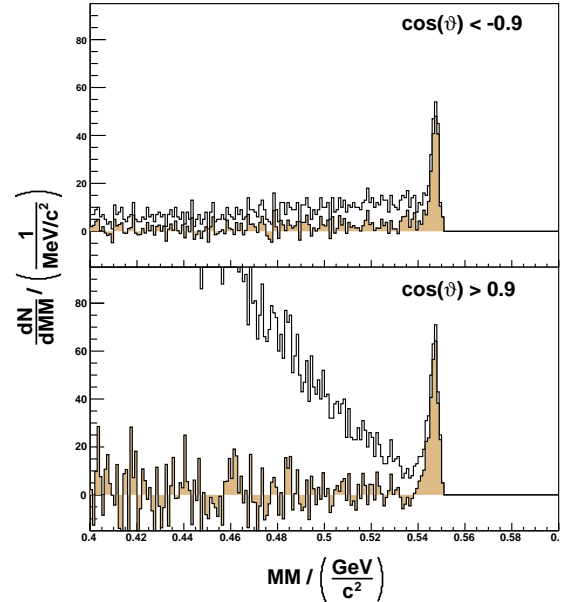


Fig. 1: Missing mass spectra for  $d + p \rightarrow {}^3\text{He} + \eta$  with the  ${}^3\text{He}$  particles scattered in backward (upper frame) and in forward direction (lower frame).

The data on the  ${}^3\text{He}\eta$  production, which were obtained with a ramped COSY beam, are currently analysed for ten energy bins close to threshold with a bin width of 1 MeV each. To extract the tensor analysing power  $t_{20}$  for each of these energy bins, the ratio of the differential cross sections as well as the individual tensor polarizations  $p_{zz}$  are currently determined. Although still preliminary the  $t_{20}$  are found to be compatible with the data presented in [5]. Final results are expected to be available soon and will provide further information on the energy dependence of  $t_{20}$  in a broader excess energy range.

### References:

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