

The K^- - ${}^3\text{He}$ scattering length and the reaction $pd \rightarrow {}^3\text{He} K^+ K^-$ near threshold

V. Yu. Grishina^a, M. Büscher, L.A. Kondratyuk^b and C. Wilkin^c

In a recent paper [1], we presented an estimation of the s -wave $K^- \alpha$ scattering length $A(K^- \alpha)$ obtained within the Multiple Scattering Approach. We further discussed the possibility of determining $A(K^- \alpha)$ from the $K^- \alpha$ invariant mass distribution measured in the reaction $dd \rightarrow \alpha K^+ K^-$ near threshold. Here we show the results of similar calculations for $A(K^- {}^3\text{He})$ (see also Ref. [2].)

In Table 1 the predictions of $A(K^- {}^3\text{He})$ are given for different $\bar{K}N$ inputs. These are taken from a K -matrix fit (Set 1) [3], a constant scattering-length fit (Set 2) [4], and a chiral unitary approach (Set 3) [5].

Table 1: The K^- - ${}^3\text{He}$ scattering length for different $\bar{K}N$ scattering lengths $a_I(\bar{K}N)$ ($I = 0, 1$).

Ref.	$a_0(\bar{K}N)$ [fm]	$a_1(\bar{K}N)$ [fm]	$A(K^- {}^3\text{He})$ [fm]
[3]	$-1.59 + i0.76$	$0.26 + i0.57$	$-1.50 + i0.83$
[4]	$-1.03 + i0.95$	$0.94 + i0.72$	$-1.52 + i1.80$
[5]	$-1.31 + i1.24$	$0.26 + i0.66$	$-1.66 + i1.10$
[6]	$2.88 + i1.12$	$0.43 + i0.30$	$-3.93 + i4.03$

We note that the $\bar{K}N$ scattering lengths described by Sets 1-3 correspond to their vacuum values. In order to demonstrate the sensitivity of our results to possible modifications of the $\bar{K}N$ scattering amplitudes in the presence of nuclear matter, we consider as Set 4 the strongly attractive in-medium solution found in Ref. [6] and given by

$$a_0^{\text{eff}} = (2.9 + i1.1) \text{ fm}, \quad a_1^{\text{eff}} = (0.43 + i0.30) \text{ fm}.$$

The results of our calculations are listed in the last column of Table 1. The values are very similar for Sets 1 and 3, being in the range $A(K^- {}^3\text{He}) = -(1.5 \div 1.7) + i(0.83 \div 1.1)$ fm. Set 2 gives a much larger imaginary part, with $A(K^- {}^3\text{He}) = -1.52 + i1.80$ fm. The “exotic” Set 4 leads to a very large $K^- {}^3\text{He}$ scattering length, with a real part of -4 fm.

We now consider the $K^- {}^3\text{He}$ FSI effect in the reaction $pd \rightarrow K^+ K^- {}^3\text{He}$ near threshold. In Fig. 1 we present calculations of the $K^- {}^3\text{He}$ relative energy spectrum at an excess energy of 40 MeV. The solid line shows the calculations for pure phase space, *i.e.* for a constant production amplitude and neglecting FSI. All other lines in the figure represent the results obtained using the different sets of $K^- N$ parameters given in Table 1. The predictions are normalized to the total $pd \rightarrow K^+ K^- {}^3\text{He}$ cross section of 9.3 nb. It is clear that the FSI can change significantly the $K^- {}^3\text{He}$ mass spectrum. The blue line corresponding to Set 4, which demonstrates a very pronounced deformation of the $K^- {}^3\text{He}$ invariant mass spectrum in the region of small invariant masses, is in clear contradiction to the data.

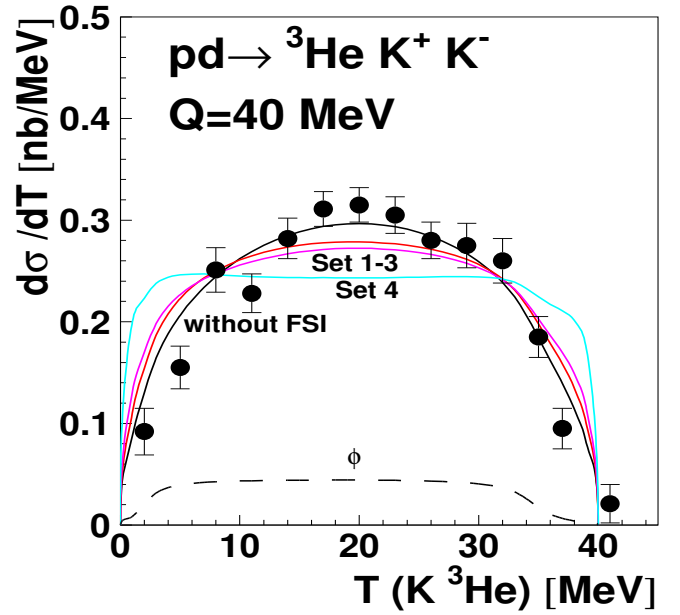


Fig. 1: Distribution of $K^- {}^3\text{He}$ relative energies from the $pd \rightarrow {}^3\text{He} K^+ K^-$ reaction at 40 MeV. The MOMO experimental data are taken from Refs. [7, 8]. The model predictions (lines) are symmetric around $T = 20$ MeV since the MOMO experiment was insensitive to the signs of the kaon charges. The $\phi(1020)$ resonance, whose effects are shown by the dashed line, contributes about 15% to the total rate.

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^aINR, Moscow, Russia

^bITEP, Moscow, Russia

^cPhysics & Astronomy Dept., UCL, London, UK

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